



Numerical simulations of mixing conditions and aerosol dynamics in the CERN CLOUD chamber

J. Voigtlander¹, J. Duplissy^{2,3}, L. Rondo⁴, A. Kürten⁴, and F. Stratmann¹

¹Leibniz Institute for Tropospheric Research, Leipzig, Germany

²CERN, PH Department, Geneva, Switzerland

³Department of Physics, University of Helsinki, P.O. Box 64, 00014 Helsinki, Finland

⁴Institute for Atmospheric and Environmental Science, Goethe-University, Frankfurt/Main, Germany

Correspondence to: J. Voigtlander (jensv@tropos.de)

Received: 24 June 2011 – Published in Atmos. Chem. Phys. Discuss.: 14 July 2011

Revised: 26 January 2012 – Accepted: 16 February 2012 – Published: 29 February 2012

Abstract. To study the effect of galactic cosmic rays on aerosols and clouds, the Cosmics Leaving OUtdoor Droplets (CLOUD) project was established. Experiments are carried out at a 26.1 m³ tank at CERN (Switzerland). In the experiments, the effect of ionizing radiation on H₂SO₄ particle formation and growth is investigated. To evaluate the experimental configuration, the experiment was simulated using a coupled multidimensional computational fluid dynamics (CFD) – particle model. In the model the coupled fields of gas/vapor species, temperature, flow velocity and particle properties were computed to investigate mixing state and mixing times of the CLOUD tank's contents. Simulation results show that a 1-fan configuration, as used in first experiments, may not be sufficient to ensure a homogeneously mixed chamber. To mix the tank properly, two fans and sufficiently high fan speeds are necessary. The 1/e response times for instantaneous changes of wall temperature and saturation ratio were found to be in the order of few minutes. Particle nucleation and growth was also simulated and particle number size distribution properties of the freshly nucleated particles (particle number, mean size, standard deviation of the assumed log-normal distribution) were found to be distributed over the tank's volume similar to the gas species.

1 Introduction

Atmospheric aerosols are complicated multiphase systems, influencing Earth' climate directly via absorption and scattering of solar radiation and indirectly via cloud processes. Key parameters for physical and chemical behavior are microphysical properties, i.e., particle/droplet number size distribution and chemical composition of the particles. However, the processes, which control these properties are not well understood. Largest uncertainties in understanding the current climate change are attributed to aerosols and clouds (IPCC2007). These uncertainties partly result from solar-related contributions and require further research. For example, still under discussion are galactic cosmic ray ionization effects on aerosols and clouds (e.g., Carslaw et al., 2002; Enghoff and Svensmark, 2008; Kirkby, 2007; Kulmala et al., 2010; Svensmark and Friis-Christensen, 1997).

To investigate quantitatively both, particle nucleation and the effects of ionisation on particle nucleation, the Cosmics Leaving OUtdoor Droplets (CLOUD) project was established. Within this project, experiments are carried out at a large volume cloud chamber (26.1 m³) located at CERN (Switzerland). The chamber has been carefully designed for carrying out experiments under very clean and thermodynamically stable conditions (Kirkby et al., 2011) and can be exposed to a particle beam provided by the CERN Proton Synchrotron (PS) particle accelerator. The particle beam is applied to create ions and to study their effect on aerosol particle formation and on cloud condensation nuclei and ice

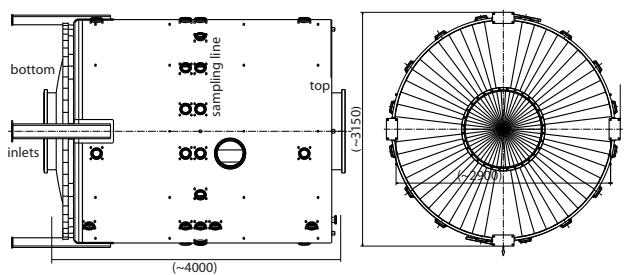


Fig. 1. Schematic diagram of the CLOUD chamber, a stainless steel cloud tank with 26.1 m^3 . According to the figures shown in the following, the front view of the tank was rotated by 90° (clockwise). At the mid-height, the holes for the sampling probes (sampling line) are visible.

nuclei activation. The chamber is equipped with a large number of different instruments to study aerosol cloud cosmic rays micro physics under well defined conditions.

A big issue in large volume cloud chambers like the CLOUD chamber at CERN is to attain special homogeneity regarding the prevailing thermodynamic conditions, gas concentrations and particle properties. Achieving homogeneity inside the tank becomes even more challenging, if several parameters are changed during the experiments (e.g., via UV-illumination system, particle nucleation, trace gas input, wall cooling). To ensure homogeneity, mixing fans are usually applied in such experiments. To check the mixing state, measurements are made at several selected points of the chamber. However, it is not possible to check all parameters continuously, and therefore numerical simulations are helpful and necessary.

A theoretical study concerning cloud droplet formation in a similar (shape, aspect ratio), but smaller (12 m^3) cloud tank was already given by Schütze and Stratmann (2008). It was stated that a 2-fan configuration gives, with respect to homogeneity in the chamber, much better results than a 1-fan configuration. However, simulations presented in Schütze and Stratmann (2008) were not performed for the actual CLOUD geometry and not evaluated with experimental cloud chamber data.

This paper presents numerical simulation results for a cloud tank with the geometry of the CLOUD chamber concerning the mixing state with respect to both, gaseous species and particles. Simulations were carried out using a coupled computational fluid dynamics (CFD) – particle model called CLOUD-FPM. After a very brief description of the experimental set up, fundamental aspects of the simulations are given, followed by several model results in comparison to experimental data.

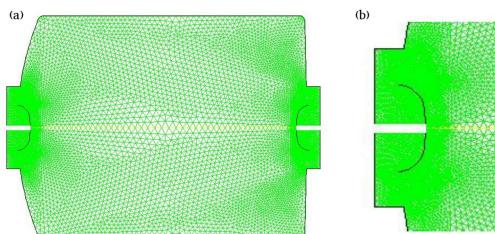


Fig. 2. (a) Numerical grid used for the CLOUD-FPM simulations. The grid has about 20 000 grid cells and was generated according to the requirements of the turbulence model applied for the simulations. The thin lines next to the bottom and the top are representing the mixing fans and the small, not meshed areas in between are the fan hubs. The mixing fan layer is also shown in detail in (b).

2 The CLOUD chamber

The cylindrical CLOUD chamber, located at CERN, is an electro-polished stainless steel tank with a diameter of 3.0 m , a height of approx. 4.0 m and a corresponding volume of 26.1 m^3 . The tank was designed, after a pilot experiment (Duplissy et al., 2010), with respect to achieve highest standards of cleanliness and temperature stability (Kirkby et al., 2011). A schematic diagram of the chamber is shown in Fig. 1. Different inlets and outlets at the chamber wall can be used to connect sampling probes, to introduce trace gases into the chamber, and to evacuate the chamber. To continuously mix the tank's contents, two fans can be installed next to the flanges at the top and the bottom (thin black lines in Fig. 2). At the top of the tank, there is also an UV-illumination system (Kupc et al., 2011), illustrated in Fig. 3), which is used to trigger the OH production via ozone photolysis. The OH radicals then react with SO_2 to form sulfuric acid (H_2SO_4). Dependent on H_2SO_4 precursor concentrations, thermodynamic conditions and beam intensity, H_2SO_4 particle nucleation will occur (Kirkby et al., 2011).

3 Numerical model

3.1 General remarks

An important requirement for the experiments at the CLOUD chamber are well positioned sampling points and a tank's mixing state being as homogeneous as possible. In case of a non well-mixed chamber, the samples may not be representative for the whole tank. To evaluate the experimental configuration, the CLOUD chamber was simulated using the commercially available computational fluid dynamics (CFD) code FLUENT (ANSYS Inc., Canonsburg, PA, USA) together with the Fine Particle Model FPM, Particle Dynamics GmbH, Leipzig, Germany, Wilck et al., 2002). The FLUENT model allows the simulation of a wide range of small scale fluid problems, while the FPM was developed

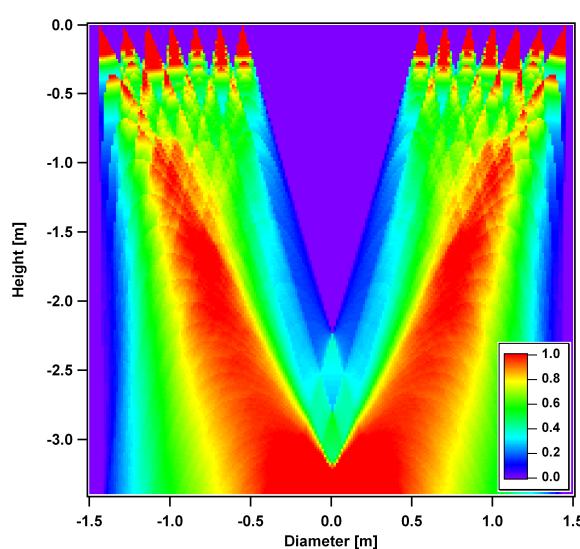


Fig. 3. UV illumination below 317 nm inside the CLOUD chamber (Kupc et al., 2011). In the numerical simulations, H_2SO_4 production was assumed to take place only in the parts of high UV intensity (red and yellow areas).

as a general purpose particle dynamics model. Both models together form the so-called CLOUD-FPM, a model being capable of handling the coupled fluid and particle dynamical processes taking place inside the CLOUD chamber. In CLOUD-FPM, all relevant properties like velocity, temperature, pressure, turbulence parameters, composition of gas/particle phase, and nucleation/growth of ultrafine aerosol particles are treated explicitly. The coupling of the models also includes the consideration of mass and heat transfer between gas and particle phase due to particle nucleation and condensational growth.

For the simulations, the geometry of the CLOUD chamber must be discretized on a numerical grid. Subsequently, the fluid and particle dynamics equations are solved on this grid. Because of the cylindrical geometry of the tank, a 2-dimensional (2-D) axis-symmetric grid can be used to efficiently evaluate the mixing state. Compared to a 3-D treatment, such a grid reduces the computational costs significantly, although it is obvious that 2-D simulations do not allow a one to one description of the experimental set up. For example, individual holes (inlets/outlets) at the chamber bottom/side cannot be simulated using an axis-symmetric grid (bottom: if outside of the tank axis).

The flow field inside the chamber is of turbulent nature, even for small flow velocities in the order of few centimeters per second. Therefore, a $k-\epsilon$ -turbulence-model (Jones and Launder, 1972; Launder and Spalding, 1973) with enhanced wall functions was applied. Utilizing such a near wall approach means that the laminar sublayer with a thickness of about 10^0 mm has to be resolved by the numerical grid. Ac-

cordingly, a grid with about 20 000 grid cells (Fig. 2) was generated and applied for the simulations shown here.

3.2 Mixing fans

One of the key parameters in the numerical simulations is a proper description of the mixing fans. As simulations on a 2-D grid do not allow a consideration of the rotating fan blades, the fans are represented by zero thickness pressure jump layers (e.g. Fluent User's guide). The pressure jump δ_p across the fan plane is described by a polynomial function dependent on flow speed. To provide an accurate representation of the flow field, pressure jump and shape of the fan planes have to be adjusted to experimental flow field data. For this study this was done by a comparison with a radial velocity profile determined 50 cm above the bottom mixing fan in the CLOUD chamber.

The measured velocity profile compared to the simulated one is shown in Fig. 4. Experimental data for a 1-fan configuration (no hood, CLOUD-1) are represented by the black dots (Fig. 4a), data for the current 2-fan configuration (CLOUD-5) are given by the blue curves (Fig. 4b). Figure 4a shows that 2-D simulations with the simplest approach of a flat disc shaped fan layer are not suitable to reproduce the experimental data. The simulated jet above the fan (Fig. 4a, green line) was not observed in the experiments. In fact, the measured profile suggests a much more divergent velocity field. To match the measured velocity profile, the shape of the fan had to be changed to arc (as illustrated in Fig. 2, velocity profile shown in Fig. 4a, red line). To improve the efficiency of the mixing fans hoods were installed around each of them (current CLOUD-5 set up). However, measured velocity profiles above the mixing fans are quite similar compared to the 1-fan configuration without a hood (CLOUD-1, Fig. 4b), indicating that the effect of the hoods is only small.

In Fig. 5 cross sectional profiles of the velocity magnitude are presented. The data visualize the jet above the flat fan (Fig. 5a) and the much more divergent flow field of the adjusted arc-fan simulation (Fig. 5b). Furthermore, it is shown that the upper half of the tank is almost not influenced by the arc shaped fan, as the velocity is almost zero in this part of the chamber. The turbulent intensity around the fan is much larger than for the flat fan approach, but turbulent mixing is limited to the region next to the fan (not shown here). It can be concluded that to mix the tank properly under realistic conditions (arc fan) there is a need for a second fan. In agreement, the simulated flow field for a 2-fan configuration indicates that the whole tank is mixed by such a set up (Fig. 5c). Finally, results in Fig. 5d illustrate the simulated flow field for a 2-fan configuration adjusted to the improved CLOUD-5 set up (with hoods), showing that the mixing fans still produce a divergent flow field. The low efficiency of the hoods is caused by the overall low flow velocities in the chamber.

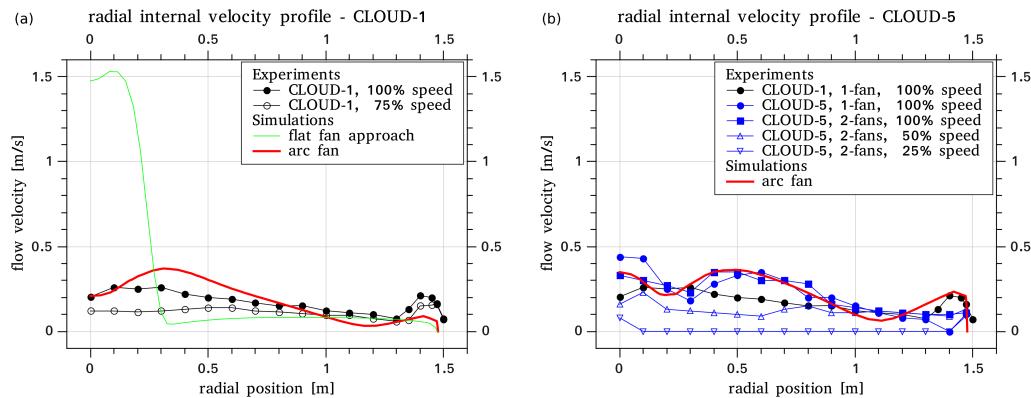


Fig. 4. Measured internal air velocity 50 cm above the fan (black and blue dots) compared to simulation results. The Fig. shows results for a 1-fan configuration (CLOUD-1, **(a)**) and a 2-fan configuration with hoods around the fans (CLOUD-5, **(b)**). The flow direction was not measured, but probably has large azimuthal and radial components, especially in the outer region. Different shapes of the zero thickness pressure jump layer (representing the mixing fan) result in different velocity profiles. A flat fan configuration results in a velocity profile significantly different from experimental data (green line). To match the experimental data, arc shaped layers are necessary (red line).

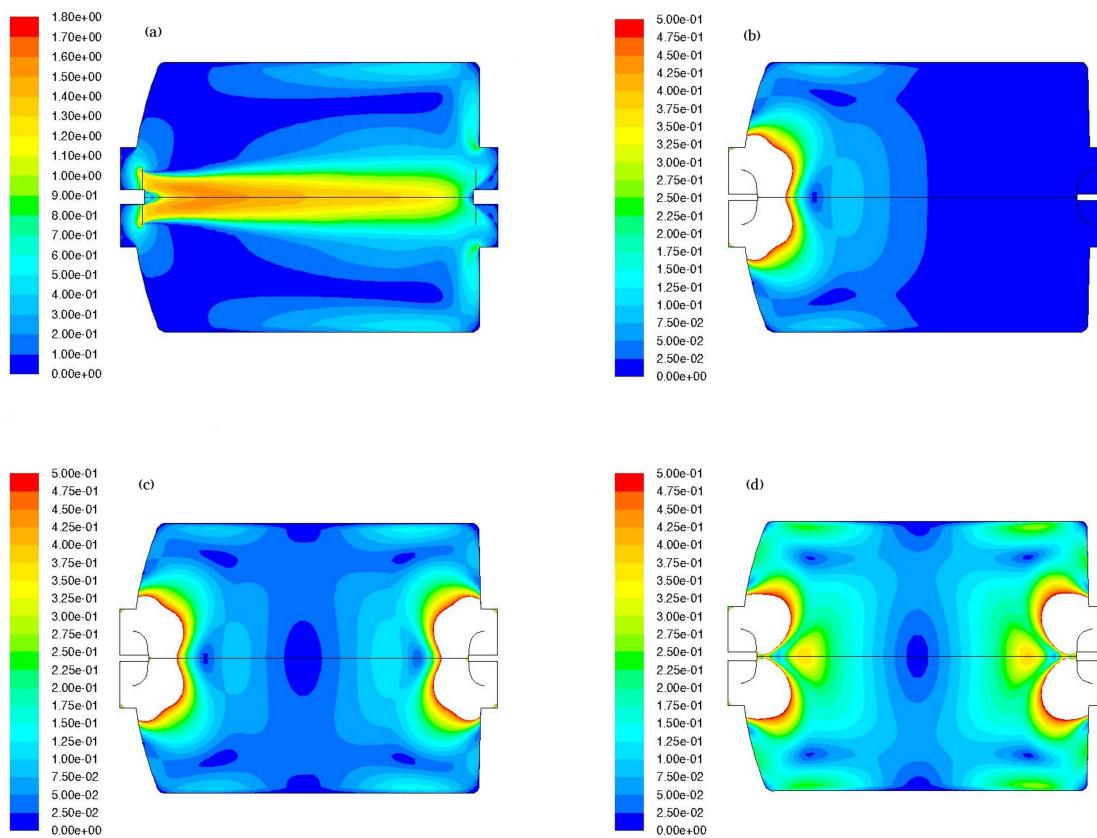


Fig. 5. Calculated cross sectional profiles of the velocity magnitude. **(a)** 1-fan configuration with flat fan approach, **(b)** 1-fan configuration with modified arc shaped pressure jump layer adjusted to the measured velocity profile 50 cm above the mixing fan (CLOUD-1), **(c)** 2-fan configuration with fan settings as used in **(b)** and **(d)** 2-fan configuration adjusted to velocity profile of the CLOUD-5 set up (with hoods around the fans).

4 Results and discussion

4.1 H₂SO₄ lifetime experiments

4.1.1 Description and experimental data

Due to molecular and turbulent vapor mass diffusion, H₂SO₄ is continuously lost to the wall. If there is no source, successive H₂SO₄ measurements allow to determine the wall losses of H₂SO₄, which depend on H₂SO₄ concentration, thermodynamic conditions and fan speed. For a proper representation of the flow field in the model (as discussed above), and suitable vapor diffusion coefficients, the numerical simulations should match the experimental data. The experiments were therefore used to validate the numerical description of the CLOUD chamber in the CLOUD-FPM model.

For the H₂SO₄ lifetime experiments, the chamber was continuously fed with the H₂SO₄ precursor gases (ozone, SO₂ and water vapor). H₂SO₄ then was produced via ozone photolysis (OH production) and reaction with SO₂. After a certain period of time, H₂SO₄ production and loss to the wall were in equilibrium, resulting in a quasi constant H₂SO₄ concentration. After switching off the UV-illumination system, the H₂SO₄ production was stopped. The subsequent decrease of H₂SO₄ due to the transport to the wall was recorded outside of the boundary layer (thickness about 10 to 20 cm) at one sampling line of the tank. H₂SO₄ was measured by a chemical ionization mass spectrometer acquired from THS instruments. The basic instrument design and the method for measuring sulfuric acid has been described in detail elsewhere (Eisele and Tanner, 1993; Berresheim et al., 2000).

The temporal decrease of the H₂SO₄ concentration for such an experiment using a 1-fan configuration is shown in Fig. 6a (black line). The x-axis gives the running time and the y-axis the H₂SO₄ concentration. Thereby, the time axis was shifted so that the H₂SO₄ decrease starts at the zero line. Figure 6a shows that, at the sampling point, the H₂SO₄ concentration is reduced by more than 90 percent after 15 min.

The experimental uncertainties of gaseous H₂SO₄ concentration measurements are about a factor of 2. On the other hand, observed short term fluctuations of the H₂SO₄ concentrations, which represent the combination of instrumental noise and local fluctuations in the small sampling volume, were much smaller (less than 20 percent, see Fig. 6). It can be concluded that the measurement uncertainties might influence the (initial) average H₂SO₄ concentration of the experiment, but did not affect the temporal characteristics of the H₂SO₄ concentrations at the sampling point, as well as the comparison with the modeling data.

4.1.2 Simulation results

For the CLOUD-FPM simulations, suitable initial and boundary conditions must be chosen. All thermodynamic conditions (e.g., T , RH) were adjusted to values equal

to those in the experiments ($T = 291.65\text{ K}$, $\text{RH} = 38\%$). The H₂SO₄ concentration was prescribed as initial value. Thereby, a homogeneous distribution was assumed.

Typical concentrations of H₂SO₄ in the CLOUD experiments are in the order of 10^6 cm^{-3} to 10^8 cm^{-3} , corresponding to a H₂SO₄ mass in the order of 10^{-12} kg up to 10^{-10} kg for the whole tank. Due to the small total amount of H₂SO₄, it is a suitable assumption to consider the tank's wall as an infinite sink for H₂SO₄. In the model this was done by defining a H₂SO₄ mass fraction of zero at the wall. Sources and additional sinks (particle nucleation) for H₂SO₄ were excluded in the simulation, because particle nucleation rates are negligible at H₂SO₄ concentrations used in these experiments (Kirkby et al., 2011).

Furthermore, binary diffusion coefficients of $0.09\text{ cm}^2\text{ s}^{-1}$ (H₂SO₄ in air) and $0.06\text{ cm}^2\text{ s}^{-1}$ (H₂SO₄ in H₂O) were applied in the simulations (Marti et al., 1997; Hanson and Eisele, 2000; Brus et al., 2010). These values are comparable to values determined using the methods of Fuller (FSG, Fuller et al., 1966), FSG-LaBas (Lyman, 1990) or Wilke and Lee (WL, Wilke and Lee, 1955), which yield $0.11\text{ m}^2\text{ s}^{-1}$, $0.093\text{ m}^2\text{ s}^{-1}$, respectively $0.1\text{ m}^2\text{ s}^{-1}$.

To evaluate the mixing state, volume average values and standard deviations of the H₂SO₄ concentrations were computed and compared to the experimental data (Fig. 6). The standard deviations represent a measure for the inhomogeneity in the CLOUD tank and were calculated as:

$$\sigma_{\text{H}_2\text{SO}_4} = \sqrt{\frac{\sum_{\text{cell}=1}^N V_{\text{cell}} (\rho_{\text{H}_2\text{SO}_4,\text{cell}} - \rho_{\text{H}_2\text{SO}_4,\text{mean}})^2}{\sum_{\text{cell}=1}^N V_{\text{cell}}}} \quad (1)$$

where $\rho_{\text{H}_2\text{SO}_4,\text{cell}}$ is the H₂SO₄ concentration in the actual grid cell and $\rho_{\text{H}_2\text{SO}_4,\text{mean}}$ is the volume average H₂SO₄ concentration.

Figure 6a shows that for the 1-fan configuration the calculated concentration at the sampling volume (red dotted line) is in good agreement with the experimental data. Calculated volume average concentrations are clearly above this values (red solid line). However, in the range of the calculated standard deviations (light gray area), the mean values also agree with the measurements. Since the concentration at the sampling volume is not representative for the whole tank, the results again indicate that the tank is not well mixed by a 1-fan configuration. Considering a 2-fan configuration, calculated standard deviations of the H₂SO₄ concentration are much smaller (Fig. 6b). On the other hand, the wall losses are larger compared to the 1-fan set up. But also for the 2-fan set up, the agreement between experimental data and simulation results is good.

In simulations shown above, the H₂SO₄ concentration was initially defined. In additional simulations, the production of H₂SO₄ in the tank was considered. Because of the quasi constant concentrations of the precursor gases during the time scale of a typical experiment (max. several hours), the production rate was assumed to be constant with respect to time.

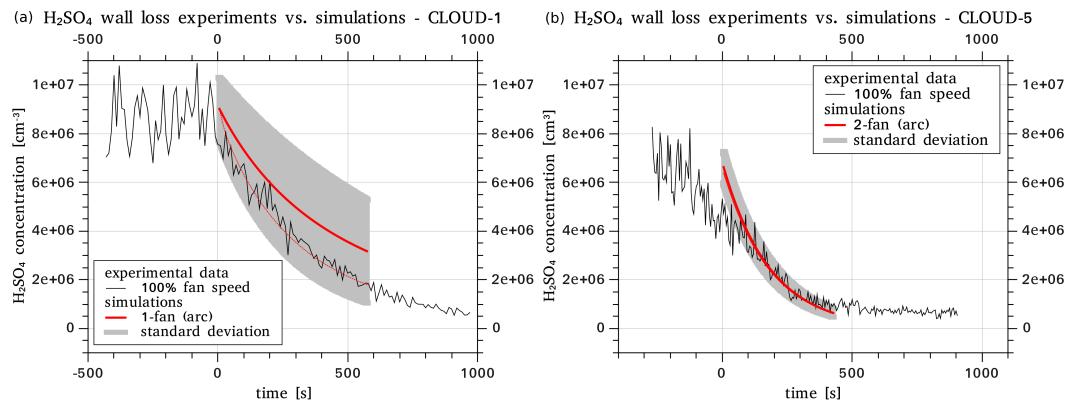


Fig. 6. H₂SO₄ lifetime experiments compared to numerical simulations for a 1-fan configuration (CLOUD-1, **(a)**) and a 2-fan configuration with hoods around the fans (CLOUD-5, **(b)**). The blue line for a 2-fan simulation (also arc shaped). The dotted lines are the results at the assumed sampling spot, the solid lines are calculated volume average values. The gray areas standard deviations, calculated for the average values.

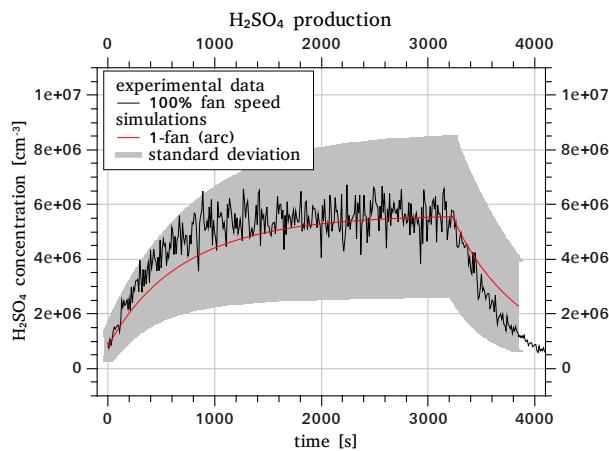


Fig. 7. Calculated H₂SO₄ concentration with constant H₂SO₄ production rate in comparison to experimental CLOUD-1 data (1-fan configuration). The source rate was $1.0 \times 10^4 \text{ cm}^{-3} \text{s}^{-1}$.

Also for that reason and to reduce the computational costs, the calculation of the chemical reactions was left out of the simulations. It was further assumed that the H₂SO₄ production only takes place in the (with high intensity) UV-light illuminated part of the tank (Fig. 3, red and yellow colored areas).

The H₂SO₄ source rates were estimated from the CLOUD experiments and span about two orders of magnitude in the range between about $10^4 \text{ cm}^{-3} \text{s}^{-1}$ and $10^6 \text{ cm}^{-3} \text{s}^{-1}$. Experimental data compared to simulation results assuming a constant source rate of $1 \times 10^4 \text{ cm}^{-3} \text{s}^{-1}$ are exemplarily shown in Fig. 7. Within the calculated standard deviation modeled and measured H₂SO₄ concentrations are in agreement, indicating that the simple approach of a constant

Table 1. 1/e volume-to-surface exchange times for an instantaneous change of wall temperature or water mass fraction for different fan configurations.

Configuration	T_{wall}		$x_{\text{H}_2\text{O},\text{wall}}$
	$t_{1/\text{e}}[\text{s}]$	$\Delta T_{\max}[\text{K}]$	
1 × fan (flat)	184	1.88	45
2 × fan (flat)	151	1.61	8
1 × fan (arc)	222	3.03	282
2 × fan (arc)	100	2.00	21
			287
			118

H₂SO₄ production rate gives proper results. Again, the large standard deviations indicate a not well mixed tank for the 1-fan configuration investigated here.

4.2 Mixing of the CLOUD tank's contents

Time resolved simulations were carried out to estimate time scales for mixing the tank's contents. To investigate volume-to-surface exchange, the response of the system to an instantaneous change of (a) the wall temperature by 20 K (291.65 to 271.65 K), and (b) the water mass fraction at the wall by 0.015 (from 0.05 to 0.20) was simulated for both, the 1-fan and the 2-fan configuration. Fan shape and pressure jump settings were the same as for the H₂SO₄ lifetime experiments described above. Simulation results of the temperature jump are shown in Fig. 8a and b. The corresponding wall exchange times to reduce the difference between wall and tank's average value to 1/e of the initial value, called 1/e time in the following, are about 100 s for the 1-fan and 220 s for the 2-fan set up (see also Table 1). This means, the second fan reduces the mixing time by a factor of more than two.

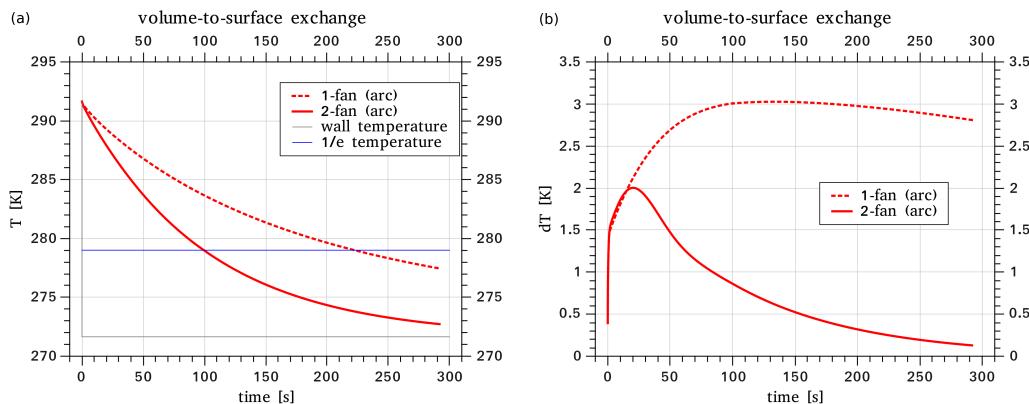


Fig. 8. Results of the simulations investigating volume-to-surface exchange in the CLOUD chamber. Shown are the temporal development of T (**a**) and ΔT (**b**) for an instantaneous drop of wall temperature by 20 K.

Similar as given for the H_2SO_4 concentration, a measure for the inhomogeneity in the tank is the volume-averaged deviation of T_{mean} given by:

$$\Delta \bar{T} = \sqrt{\frac{\sum_{\text{cell}=1}^N V_{\text{cell}} (T_{\text{cell}} - T_{\text{mean}})^2}{\sum_{\text{cell}=1}^N V_{\text{cell}}}} \quad (2)$$

where cell is the cell index, V_{cell} is cell volume and T_{cell} is cell temperature.

The volume-averaged deviation of T_{mean} is shown in Fig. 8b. The value is zero in the beginning, because the simulation started with the assumption of a homogeneous tank. Afterwards, it increases rapidly due to the temperature change of the wall. After reaching a maximum value ΔT_{mean} decreases back to zero for long time scales. The inhomogeneity of the tank is significantly reduced, if a second fan is installed (ΔT_{mean} reduced). Again it is obvious that the usage of only 1-fan (arc shaped) is not suitable to provide a well mixed tank, because ΔT_{mean} values are very high and decrease very slowly compared to the 2-fan configuration.

Volume-to-surface exchange time scales for heat and mass transport processes were found to be very similar. The 1/e-times for an instantaneous jump of the water mass fraction at the wall are almost identical to the 1/e-times for a temperature jump and are therefore not plotted again.

4.3 Simulation of particle nucleation and growth

Time dependent CLOUD-FPM calculations were also carried out to simulate nucleation and growth of $\text{H}_2\text{SO}_4\text{-H}_2\text{O}$ particles in the CLOUD tank. The focus was on the investigation of the mixing state and not a quantitative theoretical description of the experiments. For that reason, aspects of ion induced nucleation or additional condensing gases (beside H_2SO_4), as supposed in Duplissy et al. (2010) (according to Nieminen et al., 2010) and also discussed in Kirkby et al., 2011, were not considered.

For the simulations a parametrization of the H_2SO_4 vapor dependent nucleation rate and the subsequent particle growth had to be included into the model. With respect to the computational effort, classical nucleation theory is very expensive. Simplified parametrizations of nucleation rate J are based on the equation (e.g., Kulmala et al., 2004; McMurry, 1980) :

$$J = K \cdot [\text{conc}]^A \quad (3)$$

with the concentration of the considered nucleating vapor $[\text{conc}]$, and the fitting parameters K (kinetic coefficient) and A (exponential term), derived from experiments (e.g., Berndt et al., 2006; Kuang et al., 2008; Sihto et al., 2006; Riipinen et al., 2007). For simulations shown in this study, A and K were adjusted to data published in Kirkby et al. (2011) (neutral case) resulting in values $A = 3$ and $K = 0.8 \times 10^{-26} \text{ cm}^3 \text{ s}^{-1}$.

Particle growth of the freshly nucleated particles was described by a simple growth law given by Seinfeld and Pandis (1997). Applying this growth law, there is no kinetic description of the particle growth with respect to water. This means, concerning water the particles are always assumed to be in thermodynamical equilibrium. The growth law is given by:

$$\frac{dD_p}{dt} = \frac{M_s \bar{c} \alpha (C_{\text{vap}} - C_{\text{eq}})}{2\rho} WR \quad (4)$$

where D_p is particle diameter, ρ is particle density, M_s is molar weight of H_2SO_4 , \bar{c} is the mean molecular velocity of H_2SO_4 (assumed as 333 m s^{-1}), α is the mass transfer accommodation coefficient (assumed to be 1), C_{vap} is the H_2SO_4 concentration, C_{eq} is the equilibrium concentration of H_2SO_4 (assumed to be zero here), and WR is the ratio of wet to dry particle diameter (dry: only H_2SO_4). Neglecting the Kelvin term, this ratio only depends on RH (Herrmann et al., 2010). It was calculated by a linear fitting equation according to vapor pressure values given in Tabazadeh et al. (1997). The particle number size distribution was described by a single mode log-normal distribution.

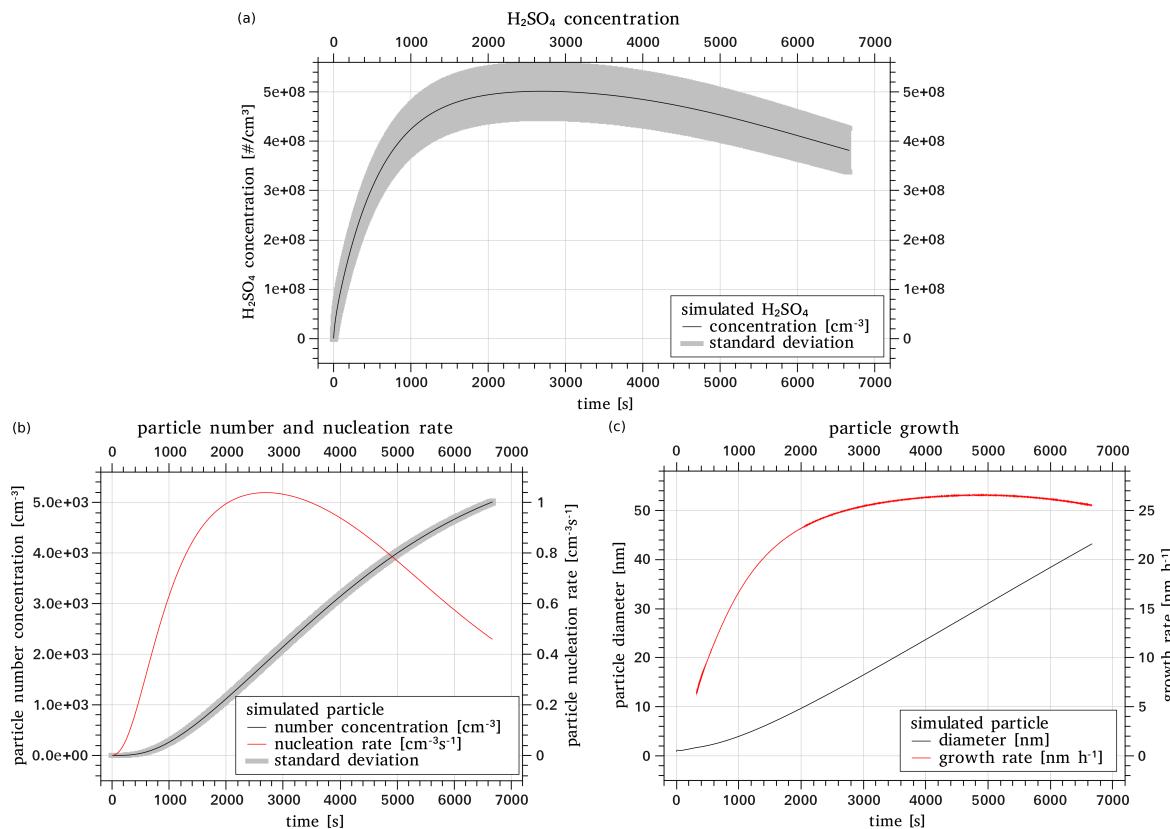


Fig. 9. Results of a simulation including particle nucleation and growth. Shown are calculated mean H₂SO₄ concentration (a), particle number concentration and nucleation rates (b), and mean particle size and growth rates (c). Details are given in the text.

The simulation shown here was done for a 2-fan configuration with arc shaped pressure jump layers as determined above. Assuming a H₂SO₄ production rate of $2.5 \times 10^6 \text{ cm}^{-3} \text{ s}^{-1}$, the resulting maximum H₂SO₄ concentration was about $5 \times 10^8 \text{ cm}^{-3}$ (see Fig. 9a). After reaching its maximum, the average H₂SO₄ concentration decreases again due to the increasing amount of H₂SO₄ transferred to the particle phase by particle nucleation and condensational growth. According to Kirkby et al. (2011), particle nucleation rates were up to about $1 \text{ cm}^{-3} \text{ s}^{-1}$ (Fig. 9b). Volume average particle diameters plotted in Fig. 9c show that particle growth rates up to $25\text{--}30 \text{ nm h}^{-1}$ were calculated.

As shown by the included (small) standard deviation of the particle number (Fig. 9b), the particles were also found to be quite homogeneously distributed over the tank's volume for this fan configuration. Furthermore, mean particle size and sigma of the log-normal distribution are also almost constant in the whole tank. Thus, the tank can be considered quite well mixed also for the freshly nucleated H₂SO₄ particles, if a suitable fan configuration, as given by the 2-fan set up, is applied.

5 Conclusions

The CLOUD-FPM model was applied to conduct numerical simulations of the CLOUD tank (26.1 m^3) established at CERN (Switzerland). In the simulations the coupled fields of temperature, saturation ratio, flow velocity, vapor species and particle size distribution were calculated.

In the model, the CLOUD chamber was described by a 2-D axis-symmetric grid. The description of the mixing fans were realized via polynomial pressure jump settings at zero thickness layers. For a suitable prediction of the flow field, the model was compared with measured velocity profiles. With adjusted fan settings, it was found that the mixing fans produce a divergent flow field with only low volume average velocity magnitudes for fan speeds used in the experiments. Considering a 1-fan configuration the simulations show that the second half of the chamber is not influenced by the fan. Consequently, a 1-fan configuration as used in the first experiments does not provide a homogeneously mixed tank. Furthermore, the simulations suggest that the efficiency of the fan-hoods is only moderate due to the low average velocity magnitude.

The adjusted model was applied to simulate the H₂SO₄ cycle in the CLOUD chamber including H₂SO₄ production and loss to the wall of the chamber. Thereby, H₂SO₄ production was calculated assuming constant H₂SO₄ production rates in the UV-illuminated part of the tank. Production rates were taken from experimental values. Simulated H₂SO₄ concentrations were compared to experimental data. Calculated and measured H₂SO₄ concentrations were in agreement within calculated standard deviations so that it can be concluded that the model provides an appropriate tool to evaluate the experimental data. Again, the simulations show that a 1-fan configuration does not provide a well mixed tank. On the other hand, the data suggest well mixed conditions for the 2-fan configuration.

The mixing conditions in the CLOUD chamber were further investigated by calculating 1/e volume-to-surface exchange times for the system response to an instantaneous change of the wall temperature or water saturation ratio. These times were found in the range of few minutes. Again, a second fan significantly reduces the standard deviations and it also may decrease the volume-to-surface exchange time by a factor of 2–3 (2 min compared to 5 min).

Particle nucleation and growth was also investigated in additional simulations. It was found that the mixing state of the particle number size distribution properties is quite similar to the mixing state of the gaseous components. In other words, if the tank is homogeneously mixed with respect to H₂SO₄, it can be also considered well mixed with respect to the freshly nucleated particles. In summary, it was found that an at least 2-fan configuration and sufficiently high fan speeds should be chosen to provide well mixed conditions inside of the CLOUD chamber.

Acknowledgements. We thank the German Federal Ministry of Education and Research (BMWF) for founding the modeling research presented here (project no. 01LK0902B). We further thank CERN for supporting CLOUD with important technical and financial resources, and for providing a particle beam from the CERN Proton Synchrotron. The CLOUD project further received funding from the EC Seventh Framework Programme (Marie Curie Initial Training Network "CLOUD-ITN" grant no. 215072, and ERC-Advanced "ATMNUCLE" grant no. 227463), the Swiss National Science Foundation (project nos. 206621_125025 and 206620_130527), the Academy of Finland Center of Excellence program (project no. 1118615), the Austrian Science Fund (FWF; project no. P19546 and L593), the Portuguese Foundation for Science and Technology (project no. CERN/FP/116387/2010), and the Russian Foundation for Basic Research (grant N08-02-91006-CERN). We also thank J. Kirkby, K. Carslaw and J. Curtius for contributions to this manuscript.

Edited by: V.-M. Kerminen

References

- Berndt, T., O. Böge, and F. Stratmann: Formation of atmospheric H₂SO₄/H₂O particles in the absence of organics: A laboratory study, *Geophys. Res. Lett.*, 33, L15817, doi:10.1029/2006GL026660, 2006.
- Berresheim, H. and T. Elste and C. Plass-Dümler and F. L. Eisele and D. J. Tanner: Chemical ionization mass spectrometer for long-term measurements of atmospheric OH and H₂SO₄, *Int. J. Mass Spectrom.*, 202, 91–109, 2000.
- Brus, D., Hyvärinen, A.-P., Viisanen, Y., Kulmala, M., and Lihavainen, H.: Homogeneous nucleation of sulfuric acid and water mixture: experimental setup and first results, *Atmos. Chem. Phys.*, 10, 2631–2641, doi:10.5194/acp-10-2631-2010, 2010.
- Carslaw, K., R. Harrison, and J. Kirkby: Cosmic rays, clouds, and climate, *Science*, 298, 1732–1737, 2002.
- Duplissy, J., Enghoff, M. B., Aplin, K. L., Arnold, F., Aufmhoff, H., Avnegaard, M., Baltensperger, U., Bondo, T., Bingham, R., Carslaw, K., Curtius, J., David, A., Fastrup, B., Gagné, S., Hahn, F., Harrison, R. G., Kellett, B., Kirkby, J., Kulmala, M., Laakso, L., Laaksonen, A., Lillestol, E., Lockwood, M., Mäkelä, J., Makhmutov, V., Marsh, N. D., Nieminen, T., Onnela, A., Pedersen, E., Pedersen, J. O. P., Polny, J., Reichl, U., Seinfeld, J. H., Sipilä, M., Stozhkov, Y., Stratmann, F., Svensmark, H., Svensmark, J., Veenhof, R., Verheggen, B., Viisanen, Y., Wagner, P. E., Wehrle, G., Weingartner, E., Wex, H., Wilhelmsson, M., and Winkler, P. M.: Results from the CERN pilot CLOUD experiment, *Atmos. Chem. Phys.*, 10, 1635–1647, doi:10.5194/acp-10-1635-2010, 2010.
- Eisele, F. L. and D. J. Tanner: Measurement of the Gas-Phase Concentration of H₂SO₄ and Methane Sulfonic-Acid and Estimates of H₂SO₄ Production and Loss in the Atmosphere, *J. Geophys. Res.*, 98, 9001–9010, 1993.
- Enghoff, M. B. and Svensmark, H.: The role of atmospheric ions in aerosol nucleation – a review, *Atmos. Chem. Phys.*, 8, 4911–4923, doi:10.5194/acp-8-4911-2008, 2008.
- Hanson, D., and Eisele, F.: Diffusion of H₂SO₄ in humidified nitrogen: Hydrated H₂SO₄, *J. Phys. Chem. A*, 104, 1715–1719, 2000.
- Fuller, E. N., Schettler, P. D., and Giddings, J. C.: A new method for prediction of binary gas-phase diffusion coefficients, *Ind. Eng. Chem.*, 58, 18–27, 1966.
- Herrmann, E., Brus, D., Hyvärinen, A.-P., Stratmann, F., Wilck, M., Lihavainen, H., and Kulmala, M.: A computational fluid dynamics approach to nucleation in the water-sulfuric acid-system, *J. Phys. Chem.*, 114, 8033–8042, 2010.
- IPCC2007: Climate change 2007, Tech. rep., Cambridge University Press, 2007.
- Jones, W. and Launder, B.: The prediction of laminarization with a two-equation model of turbulence, *Int. J. Heat Mass Transfer*, 15, 301–314, 1972.
- Kirkby, J.: Cosmic rays and climate, *Surv. Geophys.*, 28, 333–375, 2007.
- Kirkby, J., Curtius, J., Almeida, J., Dunne, E., Duplissy, J., Ehrhart, S., Franchin, A., Gagne, S., Ickes, L., Kürten, A., Kupc, A., Metzger, A., Riccobono, F., Rondo, L., Schobesberger, S., Tsagkogeorgas, G., Wimmer, D., Amorim, A., Bianchi, F., Breitenlechner, M., David, A., Dommen, J., Downard, A., Ehn, M., Flagan, R. C., Haider, S., Hansel, A., Hauser, D., Jud, W., Junninen, H., Kreissl, F., Kvashin, A., Laaksonen,

- A., Lehtipalo, K., Lima, J., Lovejoy, E. R., Makhmutov, V., Mathot, S., Mikkila, J., Minginette, P., Mogo, S., Nieminen, T., Onnela, A., Pereira, P., Petaja, T., Schnitzhofer, R., Seinfeld, J. H., Sipila, M., Stozhkov, Y., Stratmann, F., Tome, A., Vanhanen, J., Viisanen, Y., Vrtala, A., Wagner, P. E., Walther, H., Weingartner, E., Wex, H., Winkler, P. M., Carslaw, K. S., Worsnop, D. R., Baltensperger, U., and Kulmala, M.: Role of sulphuric acid, ammonia and galactic cosmic rays in atmospheric aerosol nucleation., *Nature*, 476, 429–433, doi:10.1038/nature10343, 2011.
- Kuang, C., McMurry, P., McCormick, A., and Eisele, F.: Dependence of nucleation rated on sulfuric acid vapor concentration in diverse atmospheric locations, *J. Geophys. Res.*, 113, D10209, doi:10.1029/2007JD009253, 2008.
- Kulmala, M., Vehkamki, H., Petj, T., DalMaso, M., Lauri, A., Kerminen, V.-M., Birmili, W., and McMurry, P.: Formation and growth rates of ultrafine atmospheric particles: a review of observations, *J. Aerosol Sci.*, 35, 143–176, 2004.
- Kulmala, M., Riipinen, I., Nieminen, T., Hulkkoenen, M., Sogacheva, L., Manninen, H. E., Paasonen, P., Petäjä, T., Dal Maso, M., Aalto, P. P., Viljanen, A., Usoskin, I., Vainio, R., Mirme, S., Mirme, A., Minikin, A., Petzold, A., Hrrak, U., Plaß-Dülmer, C., Birmili, W., and Kerminen, V.-M.: Atmospheric data over a solar cycle: no connection between galactic cosmic rays and new particle formation, *Atmos. Chem. Phys.*, 10, 1885–1898, doi:10.5194/acp-10-1885-2010, 2010.
- Kupc, A., Amorim, A., Curtius, J., Danielczok, A., Duplissy, J., Ehrhart, S., Walther, H., Ickes, L., Kirkby, J., Kürten, A., Lima, J. M., Mathot, S., Minginette, P., Onnela, A., Rondo, L., and Wagner, P. E.: A fibre-optic UV system for H_2SO_4 production in aerosol chambers causing minimal thermal effects, *J. Aerosol Sci.*, 42, 532–543, 2011.
- Launder, B. and Spalding, D.: The numerical computation of turbulent flows, *Comput. Methods Appl. Mech. Eng.*, 3, 269–289, 1973.
- Lyman, W. J., Reehl, W. F., and Rosenblatt, D. H.: Handbook of chemical property estimation methods, American Chemical Society: Washington, DC, USA, 1990.
- Marti, J., Jefferson, A., PingCai, X., Richert, C., McMurry, P., and Eisele, F.: H_2SO_4 vapor pressure of sulfuric acid and ammonium sulfate solutions, *J. Geophys. Res.*, 102, 3725–3735, 1997.
- McMurry, P.: Photochemical aerosol formation from SO_2 : A theoretical analysis of smog chamber data, *J. Colloid Interface Sci.*, 78, 513–527, 1980.
- Nieminen, T., Lehtinen, K. E. J., and Kulmala, M.: Sub-10 nm particle growth by vapor condensation effects of vapor molecule size and particle thermal speed, *Atmos. Chem. Phys.*, 10, 9773–9779, doi:10.5194/acp-10-9773-2010, 2010.
- Riipinen, I., Sihto, S.-L., Kulmala, M., Arnold, F., Dal Maso, M., Birmili, W., Saarnio, K., Teinilä, K., Kerminen, V.-M., Laaksonen, A., and Lehtinen, K. E. J.: Connections between atmospheric sulphuric acid and new particle formation during QUEST III-IV campaigns in Heidelberg and Hytilä, *Atmos. Chem. Phys.*, 7, 1899–1914, doi:10.5194/acp-7-1899-2007, 2007.
- Schütze, M. and Stratmann, F.: Numerical simulation of cloud droplet formation in a tank, *Comput. Geosci.*, 34, 1034–1043, 2008.
- Seinfeld, J. and Pandis, S.: *Atmospheric Chemistry and Physics: From air pollution to climate change*, Wiley, New York, USA, 1997.
- Sihto, S.-L., Kulmala, M., Kerminen, V.-M., Dal Maso, M., Petäjä, T., Riipinen, I., Korhonen, H., Arnold, F., Janson, R., Boy, M., Laaksonen, A., and Lehtinen, K. E. J.: Atmospheric sulphuric acid and aerosol formation: implications from atmospheric measurements for nucleation and early growth mechanisms, *Atmos. Chem. Phys.*, 6, 4079–4091, doi:10.5194/acp-6-4079-2006, 2006.
- Svensmark, H. and Friis-Christensen, E.: Variation of cosmic ray flux and global cloud coverage - a missing link in solar-climate relationships, *J. Atm. Sol. Terr. Phys.*, 59, 1225–1232, 1997.
- Tabazadeh, A., Toon, O., Clegg, S., and Hamill, P.: A new parameterization of H_2SO_4/H_2O aerosol composition: Atmospheric implications, *Geophys. Res. Lett.*, 24, 1931–1934, 1997.
- Wilck, M., Stratmann, F., and Whitby, E.: A fine particle model for fluent: Description and application, in Proc. Sixth Int. Aerosol Conf., pp. 1269–1270, Chinese Association for Aerosol Research in Taiwan/International Aerosol Research Assembly, Taipei, Taiwan, 2002.
- Wilke, C. R. and Lee, C. Y.: Estimation of Diffusion Coefficients for Gases and Vapors, *Ind. Eng. Chem.*, 47, 1253–1257, 1955.