

Study of the isospin character of 1^- states using hadronic probes at intermediate energies

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Abstract. The complementary (γ, γ') and $(\alpha, \alpha'\gamma)$ reactions were used to study the isospin properties of low-lying $E1$ excitations in the doubly-magic nucleus ^{48}Ca . In contrast to heavier nuclei, a state-to-state change in isospin character was revealed in ^{48}Ca and a dominant isoscalar excitation was found which is interpreted as an isoscalar oscillation. Recently, protons at 80 MeV were used as an additional hadronic probe in a p - γ coincidence experiment on ^{140}Ce for the first time. Results of the experiments on ^{48}Ca and first results of the ^{140}Ce will be presented in this contribution.

1 Introduction

Low-lying electric dipole excitations, in particular the electric Pygmy Dipole Resonance (PDR), were investigated in neutron-rich nuclei using various experimental methods [1]. Aiming at isospin characters of these 1^- states, it has been demonstrated that the complementary studies of real-photon scattering and high-resolution $(\alpha, \alpha'\gamma)$ coincidence experiments at 136 MeV allow separating isovector and isoscalar dipole response which is important for a deeper understanding of the underlying mechanisms generating the dipole strength [2]. Systematic studies using these complementary probes revealed a splitting into low-lying isospin-mixed $E1$ excitations and higher-lying dominantly isovector $E1$ excitations in neutron-magic, proton-magic, and non-magic nuclei (namely ^{140}Ce , ^{138}Ba , ^{124}Sn , and ^{94}Mo) [2–5]. This experimentally observed common feature of the low-lying $E1$ response is reproduced by several theoretical calculations (see, e.g., Refs. [6–9]) which suggest a distinction between a low-lying neutron-skin oscillation mode and a transitional mode towards the well-known isovector Giant Dipole Resonance (IVGDR) [10]. Recently, the crossover between the neutron-skin mode and the higher-lying proton-neutron oscillation mode was investigated in detail by means of random-phase approximation (RPA) calculations [11] and a decomposition method was introduced which might be important for a robust comparison with experimental data. In addition to the aforementioned α -particles, the hadronic probe of ^{17}O at 20 MeV/u was recently used in a high-resolution experiment on ^{208}Pb at Legnaro National Laboratory [12].

In the following, we present results of two experiments using the hadronic probes of α particles at 136 MeV and protons at 80 MeV which we performed to achieve a better understanding of the nature of low-lying $E1$ excitations in ^{48}Ca and ^{140}Ce . The $^{48}\text{Ca}(\alpha, \alpha'\gamma)$ and $^{140}\text{Ce}(p, p'\gamma)$ experiments go beyond the systematic study of low-lying dipole excitations in $(\alpha, \alpha'\gamma)$ and (γ, γ') experiments in two ways: On the one hand, by studying a much lighter neutron-rich nucleus and, on the other hand, by using a complementary hadronic probe with a higher energy per nucleon, namely protons at 80 MeV which penetrate more deeply into the nucleus.

2 Experimental setup and results

Both experiments were performed at the KVI Groningen, The Netherlands, exploiting the Big-Bite Spectrometer in combination with an array of HPGe detectors for γ spectroscopy. The scattered particles and the de-exciting γ -rays were acquired in hardware coincidence. In the data-analysis process, particle- γ coincidence matrices are constructed which are then used to generate energy spectra with specific energy conditions on the excitation energy (E_x) and the γ -ray energy (E_γ). Gates on ground-state transitions ($E_x = E_\gamma$) highly increase the sensitivity to dipole transitions which dominantly decay via this channel. Details on the experimental setup and data analysis tools can be found in Ref. [13]. Results of the two experiments on ^{48}Ca and ^{140}Ce are presented in the following subsections. Main experimental parameters are given in Table 2.

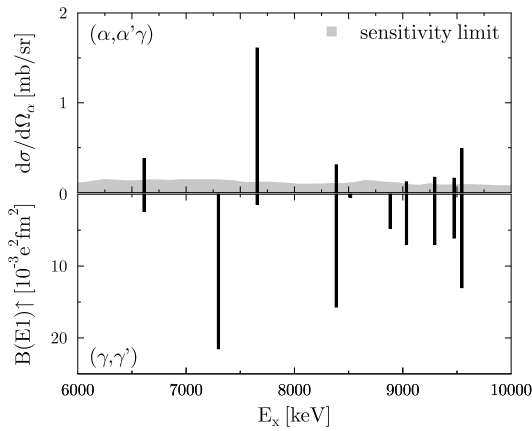
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Table 1. Experimental parameters for the $^{48}\text{Ca}(\alpha, \alpha'\gamma)$ and $^{140}\text{Ce}(p, p'\gamma)$ experiments.

	$^{48}\text{Ca}(\alpha, \alpha'\gamma)$	$^{140}\text{Ce}(p, p'\gamma)$
target thickness	1.7 mg/cm ²	20 mg/cm ²
isotopic enrichment	99%	99.72%
beam energy	136 MeV	80 MeV
average beam current	0.9 pA	0.6 pA
BBS central angle	5.8°	5.6°
collected charge	794 μC	408 μC

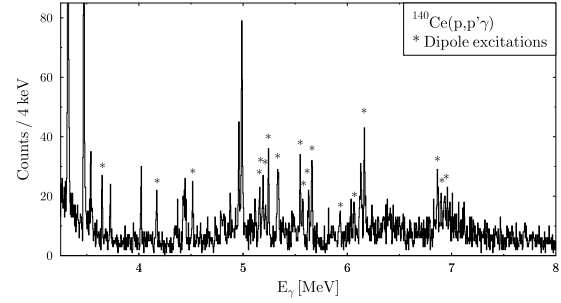
2.1 Results of the $^{48}\text{Ca}(\alpha, \alpha'\gamma)$ experiment

Low-lying $E1$ excitations and their $B(E1)$ strength distributions in stable Ca isotopes, including the doubly-magic ^{48}Ca , have been studied in several (γ, γ) experiments [14, 15]. Hence, the results of the $^{48}\text{Ca}(\alpha, \alpha'\gamma)$ experiment of this work, allow a comparison of singles α -scattering cross sections and reduced $B(E1)\uparrow$ transition strengths [14] (presented in Fig. 1) which gives access to the isospin character of low-lying 1^- states in ^{48}Ca . It reveals a state-to-state dependent isospin character of the 1^- states where isoscalar, isospin-mixed, and isovector dipole excitations are in close vicinity. An isospin splitting as observed in the heavier nuclei is not present in ^{48}Ca . Remarkable is also a strong dominantly isoscalar state at 7.6 MeV close to a strong isovector state at 7.3 MeV. The strong isoscalar state is interpreted as a pure isoscalar oscillation on basis of RPA calculations [16] under consideration of the corresponding velocity fields and transition densities. Furthermore, the well-separated dominant isovector and isoscalar states at 7.3 MeV and 7.6 MeV, respectively, are an optimal test case for the investigation of isospin mixing in a two-state mixing approach. A value of 0.061(6) was determined for the squared mixing amplitude. This results in an isospin-mixing matrix element of 85(3) keV. Further results of this experiment were presented in more detail in Ref. [17].

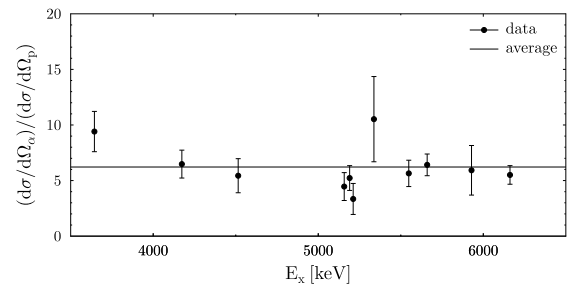
**Figure 1.** Cross sections obtained in the $^{48}\text{Ca}(\alpha, \alpha'\gamma)$ experiment (upper panel) in comparison with results of a real-photon scattering experiment [14] (lower panel).

2.2 Results of the $^{140}\text{Ce}(p, p'\gamma)$ experiment

The neutron-magic ^{140}Ce was the first nucleus for which an isospin splitting of low-lying $E1$ excitations was observed and it is the first nucleus for which the studies were extended by a $(p, p'\gamma)$ experiment at intermediate energy. The summed γ -ray spectrum of all HPGe detectors with a gate on ground-state transitions obtained in the $^{140}\text{Ce}(p, p'\gamma)$ experiment is shown in Fig. 2.

**Figure 2.** Summed γ -ray spectrum of all HPGe detectors with gate on $E_x = E_\gamma$ measured in the $^{140}\text{Ce}(p, p'\gamma)$ experiment. Dipole transitions are marked with stars.

Dipole transitions were identified via their ground-state transition energy in comparison with (γ, γ) data [18, 19]. Singles proton-scattering cross sections were determined on basis of this spectrum. Compared to the results from the inelastic α -scattering experiment, the proton-scattering cross sections are almost one order of magnitude smaller. The ratio of the cross sections is shown in Fig. 2. It is nearly constant with an average value of $(d\sigma/d\Omega_\alpha)/(d\sigma/d\Omega_p) = 6.2(7)$. The central part of the nucleon-nucleon interaction is strongly energy-dependent [20] and decreases with increasing energy from 34 MeV/u to 80 MeV/u. This also leads to a higher degree of transparency for the proton probe and additional cancellation effects for dipole transitions. The experimentally determined ratio for the cross sections is reproduced by Distorted-Wave Born Approximation (DWBA) calculations.

**Figure 3.** Ratios of singles α -scattering cross sections [2] and proton-scattering cross sections (this work) for the dipole excitations observed in both experiments. The horizontal line indicates the average ratio.

3 Summary and outlook

Inelastic scattering experiments using hadronic probes at intermediate energies, namely an $(\alpha, \alpha'\gamma)$ experiment at $E_\alpha = 136$ MeV on the doubly-magic ^{48}Ca and a $(p, p'\gamma)$ experiment at $E_p = 80$ MeV on the neutron-magic ^{140}Ce were performed at KVI Groningen to further study the nature of low-lying $E1$ strength. These experiments complement previous real-photon scattering experiments on both nuclei [14, 18] and an additional $^{140}\text{Ce}(\alpha, \alpha'\gamma)$ experiment [2]. In contrast to the observations in heavier nuclei, low-lying $E1$ excitations in the lighter nucleus ^{48}Ca show a state-to-state dependent isospin character. Furthermore, a strong isoscalar dipole state was revealed which supports theoretical predictions of a strong isoscalar oscillation [16].

For the dipole excitations in ^{140}Ce , the singles proton-scattering cross sections are considerably smaller compared to the α -scattering cross sections. Nevertheless, the general excitation pattern seems to be very similar which is supported by a nearly constant ratio of these cross sections.

In the future, different experimental techniques will be combined to achieve a deeper understanding of co-existing electric dipole modes throughout the nuclear landscape. Essential observables testing the underlying structures more thoroughly include γ -decay branchings, isospin characters, as well as single-particle contents. In particular, the isospin character of 1^- states will be addressed in particle- γ coincidence experiments with hadronic probes at intermediate energies at iThemba LABS in Somerset West, South Africa, and within an experimental PDR campaign with the CAGRA Clover-detector array at RCNP in Osaka, Japan.

Acknowledgment

This work was supported by the DFG (ZI 510/4-2), by the European Commission within the Sixth Framework Programme through I3-EURONS (contract No. RII3-CT-2004-506065), by the Alliance Program of the Helmholtz Association (HA216/EMMI), and by the Helmholtz International Center for FAIR (HIC for FAIR).

References

- [1] D. Savran, T. Aumann, A. Zilges, Prog. Part. Nucl. Phys. **70**, 210 (2013)
- [2] D. Savran, M. Babilon, A.M. van den Berg, M.N. Harakeh, J. Hasper, A. Matic, H.J. Wörtche, A. Zilges, Phys. Rev. Lett. **97**, 172502 (2006)
- [3] J. Endres, D. Savran, A.M. van den Berg, P. Dendooven, M. Fritzsche, M.N. Harakeh, J. Hasper, H.J. Wörtche, A. Zilges, Phys. Rev. C **80**, 034302 (2009)
- [4] J. Endres, E. Litvinova, D. Savran, P.A. Butler, M.N. Harakeh, S. Harissopoulos, R.D. Herzberg, R. Krücken, A. Lagoyannis, N. Pietralla et al., Phys. Rev. Lett. **105**, 212503 (2010)
- [5] V. Derya, J. Endres, M. Elvers, M.N. Harakeh, N. Pietralla, C. Romig, D. Savran, M. Scheck, F. Siebenhühner, V.I. Stoica et al., Nucl. Phys. A **906**, 94 (2013)
- [6] N. Tsoneva, H. Lenske, Phys. Rev. C **77**, 024321 (2008)
- [7] N. Paar, Y.F. Niu, D. Vretenar, J. Meng, Phys. Rev. Lett. **103**, 032502 (2009)
- [8] E.G. Lanza, A. Vitturi, E. Litvinova, D. Savran, Phys. Rev. C **89**, 041601 (2014)
- [9] E. Litvinova, P. Ring, V. Tselyaev, K. Langanke, Phys. Rev. C **79**, 054312 (2009)
- [10] M.N. Harakeh, A. van der Woude, *Giant Resonances* (Oxford University Press, New York, 2001)
- [11] H. Nakada, T. Inakura, H. Sawai, Phys. Rev. C **87**, 034302 (2013)
- [12] F.C.L. Crespi, A. Bracco, R. Nicolini, D. Mengoni, L. Pellegrini, E.G. Lanza, S. Leoni, A. Maj, M. Kmiecik, R. Avigo et al., Phys. Rev. Lett. **113**, 012501 (2014)
- [13] D. Savran, A.M. van den Berg, M.N. Harakeh, K. Ramspeck, H.J. Wörtche, A. Zilges, Nucl. Instr. and Meth. A **564**, 267 (2006)
- [14] T. Hartmann, J. Enders, P. Mohr, K. Vogt, S. Volz, A. Zilges, Phys. Rev. C **65**, 034301 (2002)
- [15] J. Isaak, D. Savran, M. Fritzsche, D. Galaviz, T. Hartmann, S. Kamedzhiev, J.H. Kelley, E. Kwan, N. Pietralla, C. Romig et al., Phys. Rev. C **83**, 034304 (2011)
- [16] P. Papakonstantinou, H. Hergert, V.Yu. Ponomarev, R. Roth, Phys. Lett. B **709**, 270 (2012)
- [17] V. Derya, D. Savran, J. Endres, M.N. Harakeh, H. Hergert, J.H. Kelley, P. Papakonstantinou, N. Pietralla, V.Yu. Ponomarev, R. Roth et al., Phys. Lett. B **730**, 288 (2014)
- [18] S. Volz, N. Tsoneva, M. Babilon, M. Elvers, J. Hasper, R.D. Herzberg, H. Lenske, K. Lindenberg, D. Savran, A. Zilges, Nucl. Phys. A **779**, 1 (2006)
- [19] B. Löher (2014), private communication
- [20] W.G. Love, M.A. Franey, Phys. Rev. C **24**, 1073 (1981)

