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Investigations of anisotropic flow using multi-particle azimuthal correlations in pp, p–Pb, Xe–Xe, and Pb–Pb collisions at the LHC

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Abstract

Measurements of anisotropic flow coefficients (v_n) and their cross-correlations using two- and multi-particle cumulant methods are reported in collisions of pp at $\sqrt{s} = 13$ TeV, p–Pb at $\sqrt{s_{NN}} = 5.02$ TeV, Xe–Xe at $\sqrt{s_{NN}} = 5.44$ TeV, and Pb–Pb at $\sqrt{s_{NN}} = 5.02$ TeV recorded with the ALICE detector. These measurements are performed as a function of multiplicity in the mid-rapidity region $|\eta| < 0.8$ for the transverse momentum range $0.2 < p_T < 3.0$ GeV/ c . An ordering of the coefficients $v_2 > v_3 > v_4$ is found in pp and p–Pb collisions, similar to that seen in large collision systems, while a weak v_2 multiplicity dependence is observed relative to nucleus–nucleus collisions in the same multiplicity range. Using the novel subevent method, v_2 measured in pp and p–Pb collisions with four-particle cumulants is found to be compatible with that from six-particle cumulants. The symmetric cumulants $SC(m, n)$ calculated with the subevent method which evaluate the correlation strength between v_n^2 and v_m^2 are also presented. The presented data, which add further support to the existence of long-range multi-particle azimuthal correlations in high multiplicity pp and p–Pb collisions, can neither be described by PYTHIA8 nor by IP-Glasma+MUSIC+UrQMD model calculations, and hence provide new insights into the understanding of collective effects in small collision systems.

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Experiments investigating ultra-relativistic collisions of heavy ions intend to recreate and explore the state of matter known as Quark–Gluon Plasma (QGP), a deconfined state of quarks and gluons, which prevailed in the early universe. Investigations of azimuthal correlations of final state particles over a wide range in pseudorapidity relative to the collision symmetry plane Ψ_n ($n \geq 1$), whose magnitude is usually quantified by flow coefficients v_n , provide important insights into the nature of the hot and dense matter created in these collisions [1–3]. Extensive measurements of v_n for inclusive [4–9] and identified hadrons [10] were performed at the Large Hadron Collider (LHC). These studies, together with quantitative predictions by hydrodynamic calculations, have enabled an extraction of the properties of the created matter [11], revealing in particular that the QGP behaves as a nearly-perfect fluid with a shear viscosity over entropy density ratio η/s close to the universal lower limit $1/4\pi$ from AdS/CFT [12]. Over the past few years, significant progress was achieved in measuring v_n distributions, as well as correlations between different flow coefficients and symmetry planes [6, 7, 13–18]. In particular, the correlation strength between different orders of flow coefficients (m and n), quantified by symmetric cumulants $SC(m, n)$ [19], was found to be sensitive to the temperature dependence of η/s and the details of the initial conditions [14]. The experimental measurements of $SC(m, n)$, together with v_n , thus provide tighter constraints on theoretical models than the individual flow coefficients alone [14, 17].

Striking similarities between numerous observables, thought to indicate the emergence of a QGP, were revealed across the pp, p–Pb and Pb–Pb collision systems at LHC energies, when compared at similar multiplicity of produced particles within a specific phase space [20]. The so-called “ridge” structure measured using two-particle correlations as a function of the pseudorapidity difference $\Delta\eta$ and the azimuthal angle difference $\Delta\phi$, which in heavy-ion collisions results from anisotropic flow, was also observed in high multiplicity p–Pb and pp collisions [21]. In addition, reported measurements of azimuthal correlations using multi-particle cumulants revealed signatures believed to be associated with the presence of collective effects in heavy-ion collisions, such as a negative four-particle cumulant $c_2\{4\}$ [22–26]. These results have indicated the presence of collective behaviour also in pp and p–Pb collisions.

Whether the observed similarities between small (pp and p–Pb) and large (Xe–Xe and Pb–Pb) collision systems arise from the same physics mechanism is under intense debate. Besides hydrodynamic descriptions [27–31], calculations from transport models [32–34], hadronic rescattering [35, 36], a string rope and shoving mechanism [37], as well as initial stage effects [38–40], can also reproduce qualitatively or even quantitatively the multiplicity dependence observed in data. No consensus on a single scenario which consistently explains all measurements across collision systems was reached so far. A better understanding of the initial conditions in small collision systems is needed, in particular in pp collisions where sub-nucleon degrees of freedom are relevant [29, 41].

In this letter, we report measurements of v_n and $SC(m, n)$ as a function of produced particle multiplicity across small and large collision systems. These measurements are expected to provide information on collective effects for all systems, which can be revealed via long-range multi-particle correlations. They rely on a new technique of performing multi-particle correlations in ranges of pseudorapidity named the *subevent method* [42, 43], where the produced particles are grouped into two or more subevents, separated by a pseudorapidity gap, to achieve a long range separation between particles that are correlated. The subevent approach minimises biases from few particle correlations such as resonances and jets, usually called non-flow, which are not associated with a collision symmetry plane.

The analyzed data are from collisions of pp at $\sqrt{s} = 13$ TeV, p–Pb at $\sqrt{s_{\text{NN}}} = 5.02$ TeV, Xe–Xe at $\sqrt{s_{\text{NN}}} = 5.44$ TeV and Pb–Pb at $\sqrt{s_{\text{NN}}} = 5.02$ TeV. They were recorded with the ALICE detector [44, 45] in Run 2 of the LHC during the years 2015, 2016 and 2017. Minimum bias events were triggered using a coincidence signal in the two scintillator arrays of the V0 detector, V0A and V0C, which cover the pseudorapidity ranges $2.8 < \eta < 5.1$ and $-3.7 < \eta < -1.7$, respectively [46]. A dedicated trigger was used in pp collisions to select high-multiplicity events based on the amplitude in both arrays of the V0 detector. The trigger selected approximately 0.1% of events with the largest multiplicity in the V0

acceptance. The average multiplicity is about 4 times larger than the one in minimum bias collisions. In comparison to minimum-bias collisions, the selection of high-multiplicity events based on forward multiplicity decreases the measurement performed at mid-rapidity, which is due to a suppression of non-flow correlations at mid-rapidity.

Only events with a reconstructed primary vertex Z_{vtx} within ± 10 cm from the nominal interaction point were selected. Removal of background events from e.g. beam interaction with the residual gas molecules in the beam pipe and pileup events was performed based on the information from the Silicon Pixel Detector (SPD) and V0 detectors. A sample of 310×10^6 high-multiplicity pp, 230×10^6 minimum bias p–Pb, 1.3×10^6 Xe–Xe and 55×10^6 Pb–Pb collisions that passed the event selection criteria was used for the analysis.

The charged tracks were reconstructed using the Inner Tracking System (ITS) [47] and the Time Projection Chamber (TPC) [48]. Only tracks with more than 70 clusters in the TPC (out of a maximum of 159) were selected. A selection on pseudorapidity to be within $-0.8 < \eta < 0.8$ ensured a high track reconstruction efficiency. Tracks with a transverse momentum $p_T < 0.2$ GeV/c were rejected due to low tracking efficiency and tracks with $p_T > 3.0$ GeV/c were removed in order to reduce the contribution from jets. A criterion on the maximum distance of closest approach (DCA) of the track to the collision point of less than 2 cm in longitudinal direction and a p_T -dependent selection in the transverse direction, ranging from 0.2 cm at $p_T = 0.2$ GeV/c down to 0.02 cm at $p_T = 3.0$ GeV/c, was applied to ensure an optimum rejection of contamination from secondary particles.

The multi-particle cumulants and symmetric cumulants were calculated from two- and multi-particle azimuthal correlations using the framework introduced in [19], which was recently extended to include the subevent method. The ranges of the subevents were chosen to be $\eta^A = (-0.8, 0)$ and $\eta^B = (0, 0.8)$ for the 2-subevent, and $\eta^A = (-0.8, -0.4)$, $\eta^B = (-0.4, 0.4)$ and $\eta^C = (0.4, 0.8)$ for 3-subevent measurements.

A correction dependent on pseudorapidity η and Z_{vtx} was applied within the generic framework to account for azimuthal non-uniformity. The correction for tracking inefficiencies was obtained from Monte Carlo simulations as a function of p_T , η and Z_{vtx} from generated particles and from tracks reconstructed after particle propagation through the detector description based on a GEANT3 simulation [49]. The systematic uncertainty was estimated by varying the event and track selection criteria, and by evaluating the analysis performance using Monte Carlo simulations with the so-called closure test. The contribution to the systematic uncertainty from the event selection was examined by narrowing the selection on Z_{vtx} to ± 5 cm. The track reconstruction biases were evaluated by tightening the selection criteria on the DCA in both the longitudinal and transverse directions, by increasing the required minimum number of TPC clusters in the track reconstruction, and by comparing the results to those obtained with other types of tracks, which have different requirements regarding the role of the ITS. The uncertainty from the Monte Carlo closure test was estimated by comparing calculations at the event generator level (i.e. invoking neither the detector geometry, nor the reconstruction algorithm) with simulation output after the full reconstruction. The individual contributions were summed in quadrature to form the systematic uncertainties. Results were obtained as a function of the number of tracks with the same selection as was applied to the measurements, which were corrected for losses due to detector inefficiencies with Monte Carlo simulations to obtain the true number of created charged particles $N_{\text{ch}}(|\eta| < 0.8)$.

Figure 1 presents the measurements of anisotropic flow coefficients $v_n\{k\}$ of harmonic n , obtained from k -particle correlations, in pp, p–Pb, Xe–Xe and Pb–Pb collisions. The collision energies are similar except for pp collisions where, however, no collision energy dependence of the integrated v_n is expected, as was shown in [25]. Figure 1 (a) shows v_2 , v_3 and v_4 measured using two-particle ($k = 2$) cumulants with a pseudorapidity gap $|\Delta\eta|$ greater than 1.4 or 1.0 units, chosen to suppress non-flow contributions. Due to the limited statistics of the pp data sample, the $|\Delta\eta|$ gap in case of v_3 and v_4 was reduced to 1 unit, consistently across all collision systems. A pronounced multiplicity dependence of v_2 is ob-

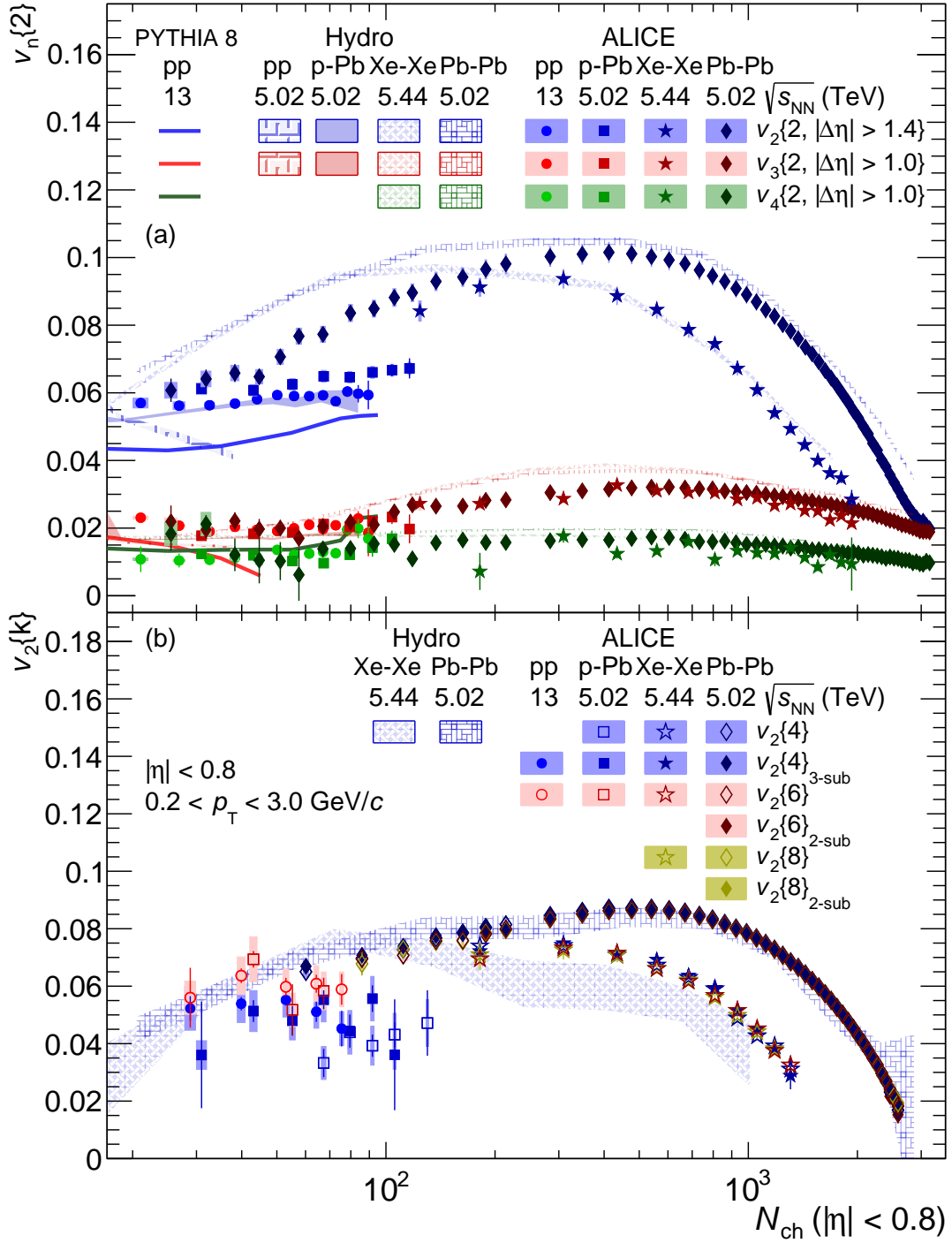


Fig. 1: Multiplicity dependence of $v_n\{k\}$ for pp, p-Pb, Xe-Xe and Pb-Pb collisions. Statistical uncertainties are shown as vertical lines and systematic uncertainties as filled boxes (coloured online). Data are compared with PYTHIA 8 [50] simulations (solid lines) of pp collisions at $\sqrt{s} = 13$ TeV and IP-Glasma+MUSIC+UrQMD [29, 51] calculations of pp, p-Pb, Pb-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV and Xe-Xe collisions at $\sqrt{s_{NN}} = 5.44$ TeV (filled bands - coloured online). (a) v_2, v_3 and v_4 measured using two-particle cumulants with a pseudorapidity gap greater than 1.4 or 1.0 units. (b) v_2 measured using multi-particle cumulants, with the 3-subevent method for the four-particle cumulant, and 2-subevent method for higher order cumulants in Pb-Pb collisions.

served in the flow dominated collision systems (Pb-Pb and Xe-Xe) as a result of the medium response

to the anisotropic shape of the initial overlap region of the colliding nuclei and the effectivity to transform the initial eccentricities to final flow coefficients, which decreases at low multiplicities. The Pb–Pb data exhibit larger values of v_2 than the Xe–Xe data, except at low multiplicity ($N_{\text{ch}}(|\eta| < 0.8) < 200$) where the two measurements become compatible. An ordering of $v_2 > v_3 > v_4$ is observed in large collision systems except for the region with the largest multiplicities, where $v_2 \approx v_3$. At low multiplicity, the magnitudes of v_n are similar to those measured in pp and p–Pb collisions. The measurements from Xe–Xe and Pb–Pb collisions are compared with recent calculations containing IP-Glasma initial conditions, a MUSIC hydrodynamic model, and the UrQMD model for hadronic rescatterings [29, 51], shown with filled bands, where the width of the band represents the statistical uncertainty of the model. The calculations qualitatively describe all the v_n measurements at large and intermediate multiplicities, while they tend to overestimate the v_2 in the low multiplicity region. In small collision systems, all the v_n coefficients exhibit a weak dependence on multiplicity. This trend cannot be explained solely by model calculations without collective effects, as demonstrated by the comparison with predictions from PYTHIA 8.210 Monash 2013 with colour reconnection [50], computed with a similar multiplicity definition as the experimental results from pp collisions. The qualitatively similar ordering of v_n in pp and p–Pb collisions with the one observed in large collision systems is not described by PYTHIA either. These observations suggest the presence of effects other than just non-flow correlations at multiplicities larger than about 2–3 times the minimum bias value of $\langle N_{\text{ch}} \rangle \approx 10$ in pp and $\langle N_{\text{ch}} \rangle \approx 24$ in p–Pb collisions within $|\eta| < 0.8$ and $0.2 < p_T < 3.0$ GeV/c. Such conclusions are supported in p–Pb collisions by the comparison with the IP-Glasma+MUSIC+UrQMD model calculations, which qualitatively describes the measurements, noting that further non-flow suppression with a larger $|\Delta\eta|$ separation in the experimental results or improvements in the phenomenological description could help to reach a quantitative agreement. Nevertheless, the model reveals a strong decrease of v_n with multiplicity in pp collisions, which is in stark contrast with the weak dependence in data. However, a simultaneous description of pp, p–Pb and Pb–Pb collision measurements was achieved in [30] with the hydrodynamic superSONIC model with hadronic rescatterings from UrQMD, which suggests that hydrodynamic collectivity prevails in both small and large collision systems.

Figure 1 (b) shows measurements of elliptic flow $v_2\{k\}$ using multi-particle cumulants with $k = 4, 6$ and 8. Measurements of $v_2\{4\}$ calculated with the 3-subevent method are also presented, and in Pb–Pb collisions the v_2 from six- and eight-particle cumulants with the 2-subevent method is shown in addition. Compared to two-particle correlations, multi-particle cumulants are less influenced by non-flow effects, since the latter usually only involve a few particles. No further non-flow suppression was observed by increasing an $|\Delta\eta|$ gap between the subevents up to 0.2 in the multi-particle cumulant measurements. Characteristic patterns of long-range multi-particle correlations, such as consistent results from the standard and subevent methods ($v_2\{4\} \approx v_2\{4\}_{3\text{-sub}}$, $v_2\{6\} \approx v_2\{6\}_{2\text{-sub}}$ and $v_2\{8\} \approx v_2\{8\}_{2\text{-sub}}$), and compatible measurements of v_2 with multi-particle cumulants ($v_2\{4\} \approx v_2\{6\} \approx v_2\{8\}$) are found in both Xe–Xe and Pb–Pb collisions, signalling a negligible contribution from non-flow correlations and the dominance of collective effects. Moreover, a good agreement of $v_2\{4\}$ calculated with the IP-Glasma+MUSIC+UrQMD [29, 51] model is found for Pb–Pb collisions down to $N_{\text{ch}} \approx 200$, below which $v_2\{2\}$ measurements are overestimated by the model. The hydrodynamic model prediction, which does not include any tuning of its parameters to other collision systems, underestimates the $v_2\{4\}$ from Xe–Xe collisions by about 15–20% in the presented multiplicity range. In p–Pb collisions, further non-flow suppression with the 3-subevent method leads to a decrease of the cumulant $c_2\{4\} > c_2\{4\}_{3\text{-sub}}$, which due to the relation $v_2\{4\} = \sqrt[4]{-c_2\{4\}}$ leads to an increase $v_2\{4\} < v_2\{4\}_{3\text{-sub}}$. The 3-subevent method allows for a measurement of a real-valued $v_2\{4\}_{3\text{-sub}}$ at a lower N_{ch} than the standard $v_2\{4\}$ measurement, making it possible to study collectivity at even lower multiplicities. An indication of multi-particle correlations in p–Pb collisions is unveiled with consistent results of v_2 measured with four- and six-particle cumulants. In pp collisions, significant contributions from non-flow effects to the four-particle cumulant ($c_2\{4\} > 0$) prevent the extraction of a real-valued $v_2\{4\}$. However, a measurement of the real-valued

$v_2\{4\}_{3\text{-sub}}$ is possible with the 3-subevent method. Similarly as for $v_2\{2, |\Delta\eta| > 1.4\}$, the $v_2\{4\}_{3\text{-sub}}$ exhibits only a weak dependence on multiplicity. Our results confirm the existence of long-range multi-particle correlations in pp and p–Pb collisions at multiplicities $N_{\text{ch}}(|\eta| < 0.8) \geq 30$. PYTHIA calculations, which do not contain genuine long-range multi-particle correlations, do not allow real values of $v_2\{4\}$ to be measured even with the subevent method [42]. The iEBE-VISHNU hydrodynamic calculations [31] can quantitatively describe all available two-particle correlation measurements in pp collisions. However, they cannot reproduce the four-particle cumulants with the currently used initial state model parameters, not even on a qualitative level. Another model with initial-state calculations predicts the measured multi-particle cumulants without involving final-state interactions [39]. Therefore, with v_n measurements alone it is not completely clear whether the origin of the apparent collectivity observed in small collision systems is the same as in large collision systems.

Further information about the origin of the observed collectivity can be extracted from the correlations between different orders of v_n^2 and v_m^2 represented by symmetric cumulants $SC(m, n)$. Figure 2 presents the multiplicity dependence of $SC(m, n)$ measured with the 3-subevent method. A positive $SC(4, 2)_{3\text{-sub}}$ is observed in large collision systems over the entire multiplicity range, similar to what was measured previously in Pb–Pb collisions at 2.76 TeV [14, 17] without the subevent method. The trend is reproduced by the IP-Glasma+MUSIC+UrQMD [29, 51] calculations. A positive $SC(4, 2)_{3\text{-sub}}$ is similarly observed both in pp and p–Pb collisions, as was found in [52]. The measurements in pp collisions are compared with PYTHIA 8 [50], which indicates a decrease of $SC(4, 2)_{3\text{-sub}}$ with decreasing multiplicity different from what is seen in data. Calculations [39, 53] with initial state correlations or parton-escape mechanism can qualitatively, or even semi-quantitatively, describe the p–Pb data. It should be noted that the results from the initial state model [39] were calculated with respect to variables that cannot be directly compared to experimental data.

An anti-correlation between v_2^2 and v_3^2 is implied by the negative $SC(3, 2)_{3\text{-sub}}$ observed in Xe–Xe and Pb–Pb collisions, similarly as in [14, 17]. Considering the linear response of v_n to initial state eccentricities ε_n (for $n = 2, 3$), the negative $SC(3, 2)_{3\text{-sub}}$ also implies an anti-correlation between initial state ε_2 and ε_3 , and thus provides tight constraints on the initial conditions in Xe–Xe and Pb–Pb collisions. There is a hint of a change to a positive sign of $SC(3, 2)_{3\text{-sub}}$ in Pb–Pb collisions at $N_{\text{ch}} \approx 100$. This tendency seems to be maintained and extended to lower multiplicities in small collision systems, suggesting a common positive correlation between v_2^2 and v_3^2 among collision systems of different sizes. Such a behaviour is not observed in the measurements from small collision systems with a larger η acceptance [52], where $SC(3, 2)_{3\text{-sub}}$ remains negative in the whole multiplicity range. Among possible reasons for this can be different contributions from non-flow effects. The IP-Glasma+MUSIC+UrQMD [29, 51] calculations for Xe–Xe and Pb–Pb collisions reproduce the negative correlation at large multiplicities. This negative sign persists in simulations down to the lowest multiplicities. PYTHIA 8 [50] fails to describe the results from pp collisions quantitatively, but it does reproduce the trend of the data. No hydrodynamic computation of $SC(m, n)$ in small systems is currently available. Nevertheless, calculations based on initial state correlations in [38, 39] reflect the crossing from negative to positive $SC(3, 2)$ in p–Pb collisions, whereas a positive correlation is observed in pp collisions [38].

While the $SC(m, n)$ encodes information on both the magnitude of and correlation between the flow coefficients, in the absence of non-flow the latter can be accessed directly by dividing $SC(m, n)_{3\text{-sub}}$ by the corresponding flow coefficients $\langle v_m^2 \rangle \langle v_n^2 \rangle$ from Figs. 2 and 1, respectively. The normalised ratio indicates that the correlation between flow coefficients is similar between different collision systems at the same N_{ch} , and reveals a large increase in magnitude in the correlation strength for collisions with N_{ch} below ≈ 100 compared to higher multiplicities. While this may be indicative of a different fluctuation pattern in low multiplicity interactions, it has to be noted that non-flow effects likely persist in this region based on the observed finite values of PYTHIA calculations in Figs. 1 and 2. Such effects make the interpretation of an increase of the normalised ratio significantly less straightforward and requires further

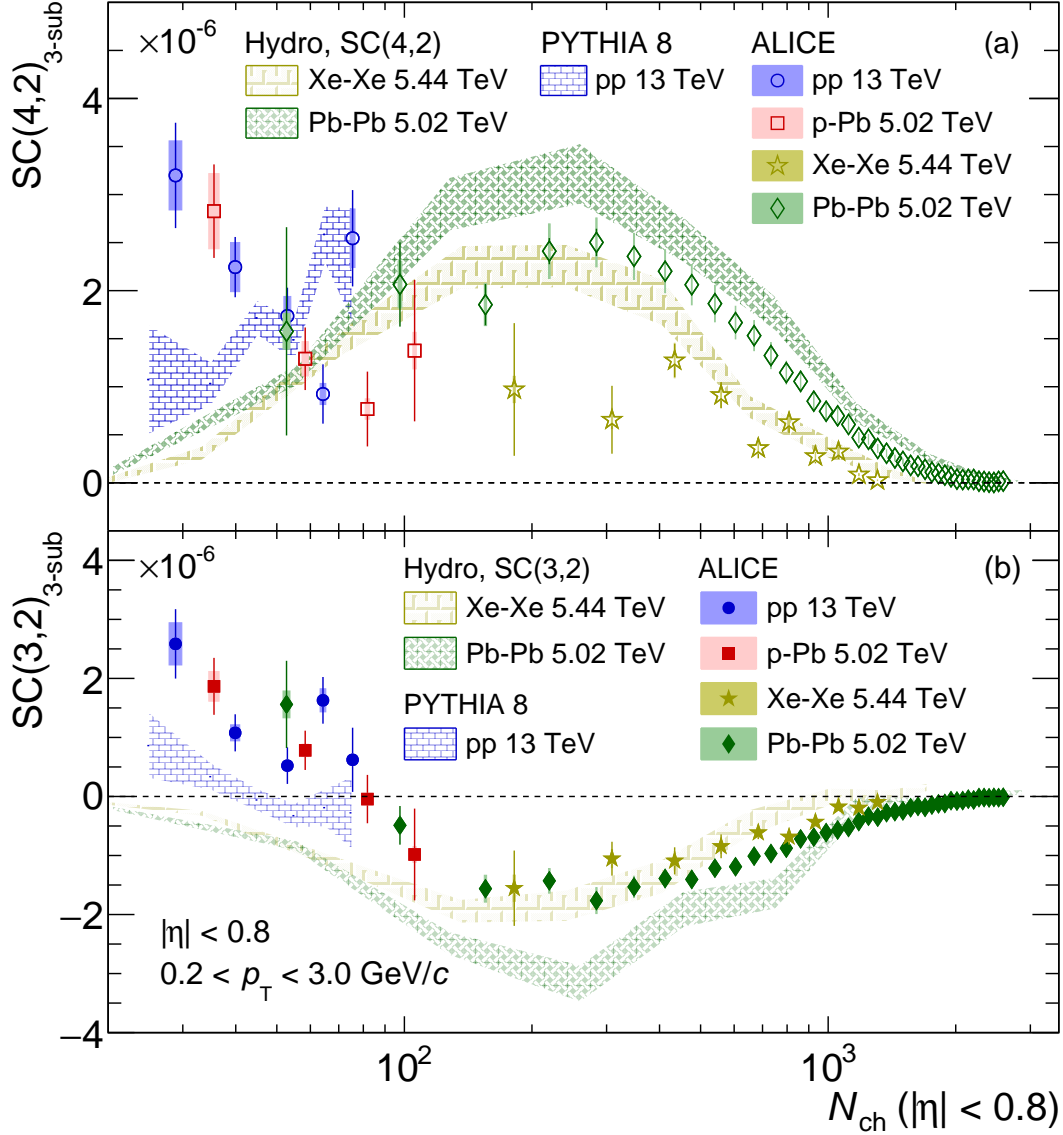


Fig. 2: Multiplicity dependence of the symmetric cumulant (a) $SC(3,2)_{3-sub}$ and (b) $SC(4,2)_{3-sub}$ for pp, p-Pb, Xe-Xe and Pb-Pb collisions. Statistical uncertainties are shown as vertical lines and systematic uncertainties as filled boxes (coloured online). The measurements in large collision systems are compared with the IP-Glasma+MUSIC+UrQMD [29, 51] calculations and results in pp collisions are compared with the PYTHIA 8 model [50].

studies. The normalised ratio $SC(m,n)_{3-sub}/\langle v_m^2 \rangle \langle v_n^2 \rangle$ is shown in Fig. A.1 of supplemental materials where it is accompanied with more discussion.

In summary, we have presented the measurements of flow coefficients $v_n\{k\}$ and symmetric cumulants $SC(m,n)_{3-sub}$ as a function of produced particle multiplicity in small (pp, p-Pb) and large (Xe-Xe, Pb-Pb) collision systems. In pp and p-Pb collisions, a non-trivial ordering $v_2 > v_3 > v_4$ and a weak dependence of v_n on the multiplicity, is observed. The values of v_n from pp and p-Pb collisions are compatible with heavy-ion collisions at low multiplicity. These first ALICE measurements of v_2 using multi-particle cumulants in small collision systems are found to be compatible with each other after suppression of non-flow contributions with the subevent method. Positive values of $SC(4,2)_{3-sub}$ are seen

in all four collision systems (pp, p–Pb, Xe–Xe and Pb–Pb). The observed anti-correlation between v_2^2 and v_3^2 measured with $SC(3,2)_{3\text{-sub}}$ in large collision systems seems to evolve into a positive correlation at low multiplicity. A similar sign change is indicated also in pp and p–Pb collisions. Thus, the different systems exhibit a similar $SC(m,n)$ at the same N_{ch} , and below $N_{ch} < 100$ reveal a large variation of the correlation strength and/or an increasing contribution of non-flow. The measurements in pp collisions can not be reproduced by the PYTHIA 8 model. The hydrodynamic description with the IP-Glasma+MUSIC+UrQMD calculations shows rather good agreement with data in Pb–Pb, Xe–Xe and p–Pb collisions, but fails to describe the measurements in pp collisions, where applicable. The presented data provide new information to help us understand the origin of the observed collectivity, and by future comparisons with phenomenological calculations with various mechanisms will disentangle the correct approach to modelling of small systems.

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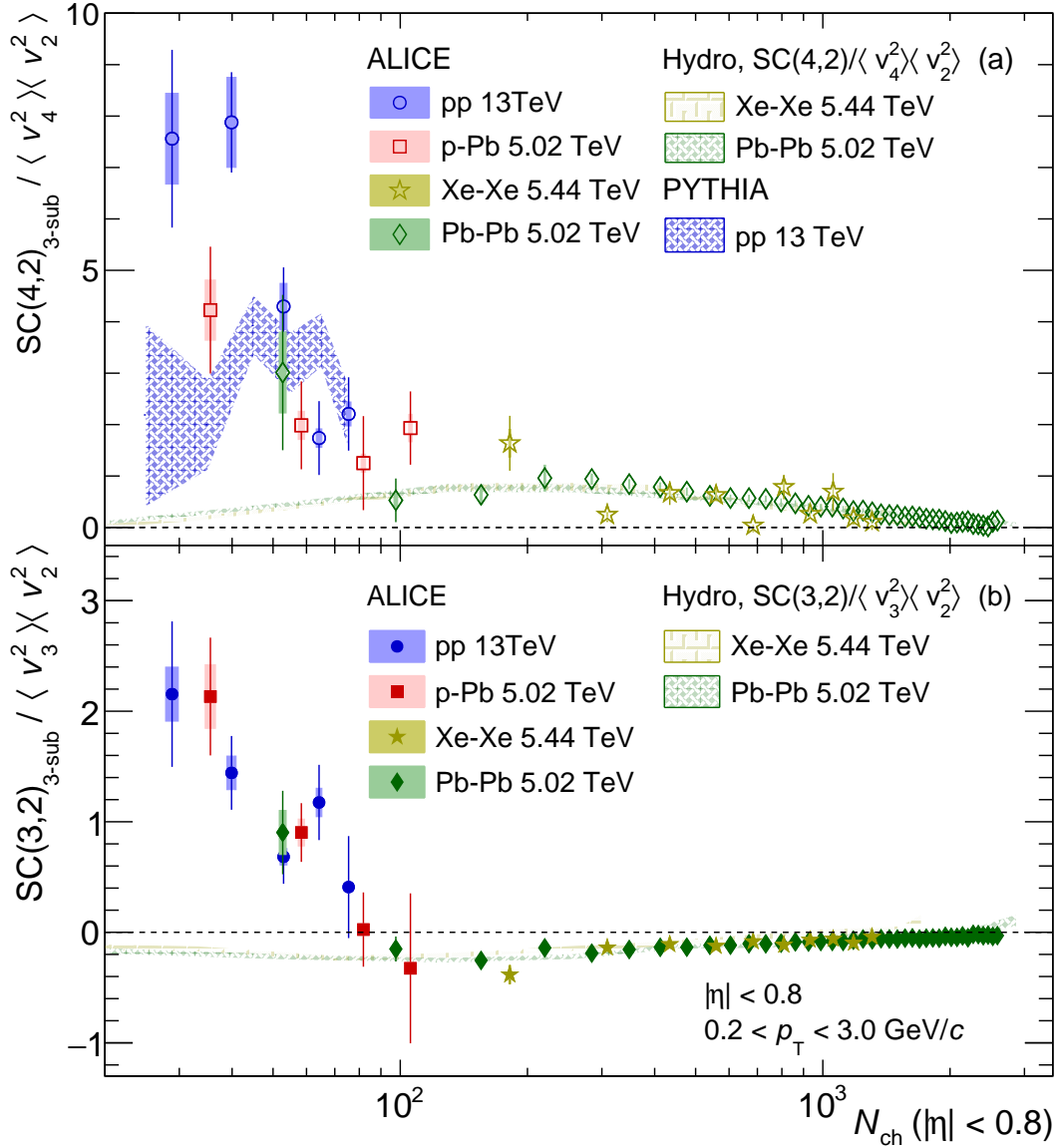


Fig. A.1: Multiplicity dependence of (a) $SC(3,2)_{3\text{-sub}}/\langle v_3^2 \rangle \langle v_2^2 \rangle$ and (b) $SC(4,2)_{3\text{-sub}}/\langle v_4^2 \rangle \langle v_2^2 \rangle$ for pp, p-Pb, Xe-Xe and Pb-Pb collisions. Observables in the denominator are obtained from the two-particle cumulants with a pseudorapidity gap $|\Delta\eta| > 1.4$ for v_2 and $|\Delta\eta| > 1.0$ for higher harmonics. Statistical uncertainties are shown as vertical lines and systematic uncertainties as filled boxes (coloured online). The results are compared with PYTHIA 8 [50] and the IP-Glasma+MUSIC+UrQMD [29, 51] calculations.

A Supplemental material

Figure A.1 shows the multiplicity dependence of the normalized ratio $SC(m,n)_{3\text{-sub}}/\langle v_n^2 \rangle \langle v_m^2 \rangle$, where the measurements in the denominator were obtained with a $|\Delta\eta|$ gap. In the absence of non-flow contamination, this observable directly reflects the correlation between v_n^2 and v_m^2 , independent of the magnitude of the individual flow coefficients v_n and v_m .

The ratio $SC(m,n)_{3\text{-sub}}/\langle v_m^2 \rangle \langle v_n^2 \rangle$ complements observations made in Fig. 2 and reveals a strong increase of the correlation strength at $N_{\text{ch}} < 100$ compared to (mid)central Pb-Pb and Xe-Xe collisions. This increase may be either due to different patterns of fluctuations or an increasing non-flow contribu-

tion towards small multiplicities, as suggested by the finite correlation strength seen in PYTHIA [50] simulations in both Fig. 1 and Fig. 2.

The data in Fig. A.1 are compared with PYTHIA 8 [50] and the IP-Glasma+MUSIC+UrQMD [29, 51] model calculations. The PYTHIA 8 calculations seem to underestimate the $SC(4,2)_{3\text{-sub}}/\langle v_4^2 \rangle \langle v_2^2 \rangle$ observed in pp collisions for $N_{\text{ch}} < 30$. The $SC(3,2)_{3\text{-sub}}/\langle v_2^2 \rangle \langle v_3^2 \rangle$ value cannot be extracted from PYTHIA calculations due to small values of v_3 . The IP-Glasma+MUSIC+UrQMD calculation describes the data at high $N_{\text{ch}} > 100$, while for $N_{\text{ch}} < 100$ the correlations seen in data are much stronger, and in case of $SC(3,2)_{3\text{-sub}}/\langle v_2^2 \rangle \langle v_3^2 \rangle$ are of an opposite sign.

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