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Rapidity dependence of antideuteron coalescence in pp collisions at $\sqrt{s} = 13$ TeV with ALICE

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Abstract

The production yields of antideuterons and antiprotons are measured in pp collisions at a center-of-mass energy of $\sqrt{s} = 13$ TeV, as a function of transverse momentum (p_{T}) and rapidity (y), for the first time up to $|y| = 0.7$. The measured spectra are used to study the p_{T} and rapidity dependence of the coalescence parameter B_2, which quantifies the coalescence probability of antideuterons. The p_{T} and rapidity dependence of the obtained B_2 is extrapolated for $p_{\mathrm{T}} > 1.7$ GeV/c and $|y| > 0.7$ using the phenomenological antideuteron production model implemented in PYTHIA 8.3 as well as a baryon coalescence afterburner model based on EPOS 3. Such measurements are of interest to the astrophysics community, since they can be used for the calculation of the flux of antinuclei from cosmic rays, in combination with coalescence models.

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1 Introduction

Antinuclei such as antideuterons and antihelium still elude detection in space and, if discovered, would likely open the door to the indirect detection of weakly interacting dark-matter (DM) candidates [1, 2]. Other scenarios speculate about the existence of anticlouds or antistars as another possible explanation of antinuclei in space [3]. Indeed, the possible presence of antinuclei in our Galaxy could be explained either by reactions of high-energy cosmic rays (CRs) with the interstellar medium (ISM) or by more exotic sources such as decays/annihilations of DM candidates. Proton–proton (pp) and proton–nucleus (p–A) collisions are notably interesting in such a context, since the collisions between CRs and the ISM are the most relevant sources for the formation of nuclei in the Galaxy, because both, CRs and the ISM, consist mostly of hydrogen ($\sim 90\%$) and helium ($\sim 9\%$), and only in small percentage of heavier nuclei. The observation of a significant antimatter excess with respect to the expected background of antimatter produced in ordinary cosmic-ray interactions would represent a signal for dark-matter annihilation in the galactic halo or for the existence of antimatter islands in our Universe [3, 4].

Information on the production of antinuclei from accelerator experiments, particularly in small collision systems (i.e., pp and p–A collisions), is essential for the theoretical description of the background constituted by antinuclei from cosmic-ray collisions with ISM. Since the ISM is almost at rest and the CRs have kinetic energies peaked at around 300 GeV, the values of relevant center-of-mass energies per nucleon–nucleon collision for the antideuteron production range from $\sqrt{s_{\mathrm{NN}}}$ $\sim$17 GeV to several TeV [5].

The goal of the existing experimental programs at different accelerator facilities is to obtain antinucleus production cross sections for pp and p–A reactions over a broad range of kinetic energies, to cover as much as possible the energy range spanned by the cosmic rays. To this end, several experiments have measured the production of (anti)nuclei in different collision systems and center-of-mass energies. These include the low collision energies of the AGS [6–9], SPS [10], and RHIC [11–16], and the TeV scale energies of the Large Hadron Collider (LHC) [17–36]. Such measurements are crucial for correct interpretations of any future measurement in satellite and balloon-borne experiments, such as AMS-02 [37], GAPS [38], or BESS-Polar [39]. An ultimate goal of the experimental programs at accelerators is to pin down the microscopic production process of antinuclei in hadronic collisions or DM decay. In fact, the production mechanism of light (anti)nuclei in high-energy hadronic collisions is still not clear and much debated in the scientific community. The plethora of experimental data collected in the last decade is typically described using two different phenomenological models: the statistical hadronization model (SHM) and the baryon coalescence approach. In the SHM [40–46], light (anti)nuclei, as well as other hadron species, are assumed to be emitted by a source in local thermal and hadrochemical equilibrium with their abundances being fixed at the moment of the chemical freeze-out of the system created in the collision, at a temperature of $T_{\mathrm{chem}} \sim 156$ MeV [47]. This model provides an excellent description of the measured hadron yields in central nucleus–nucleus collisions [44], but struggles to reproduce the evolution of the ratios between integrated yields of nuclei and those of protons with the multiplicity of particles produced in the collision [33, 34]. In the coalescence model [48–54], multi-baryon states are assumed to be formed by coalescence of baryons that are close in phase space at kinetic freeze-out, which occurs at a later time than the chemical freeze-out. In the simplest implementations of coalescence, only the momentum correlations are considered and the bound states are formed if the difference in momentum among the nucleons lies below a given threshold, namely the coalescence momentum p_0. In the state-of-the-art implementations of the coalescence approach, the quantum-mechanical properties of the constituent baryons and of the final bound states are taken into account and the coalescence probability is calculated from the overlap between the wave functions of individual (point-like) baryons and the Wigner density of the final-state cluster [55].

From the experimental point of view, the coalescence probability is related to the coalescence parameter B_A, which is obtained from the ratio of the invariant yield of the nuclei with mass number A and that of the protons raised to the power of A, assuming protons and neutrons to have the same transverse momentum distributions as they are isospin partners, and $p_{\rm T}^{\rm p} = p_{\rm T}^{A}/A$

$$B_A = \left(\frac{1}{2\pi p_{\rm T}^{A}}\left(\frac{{\rm d}^2N}{{\rm d}y{\rm d}p_{\rm T}}\right)_A\right) \Big/ \left(\frac{1}{2\pi p_{\rm T}^{\rm p}}\left(\frac{{\rm d}^2N}{{\rm d}y{\rm d}p_{\rm T}}\right)_{\rm p}\right)^A . \quad (1)$$

The production of (anti)nuclei at the LHC was so far measured in the midrapidity region $|y| < 0.5$, and then, combined with several coalescence models, employed to predict the flux of antinuclei from CR interactions at forward rapidity. However, the possible impact of a rapidity dependence of the production yield and of the coalescence probability of antinuclei was suggested in Ref. [56]. Previous measurements of the production of (anti)protons and (anti)deuterons in different rapidity intervals have been carried out in heavy-ion collisions at RHIC energies by the BRAHMS Collaboration [13]. In this Letter we report the first measurements of the rapidity dependence of the production yield of $\overline{\rm p}$ and $\overline{\rm d}$ and of the coalescence parameter B_2 (obtained from Eq. 1 with $A = 2$) carried out in pp collisions at LHC energies.

The results presented in this Letter contribute to the understanding of the impact that the extrapolation of the production yield and coalescence probability at forward rapidity, constrained to experimental information, has on the flux of antinuclei from cosmic rays, which can be obtained using hadronic production mechanisms based on coalescence. By measuring the production of antideuterons and antiprotons as a function of rapidity in pp collisions at $\sqrt{s}$ =13 TeV up to $|y| = 0.7$, the extrapolation at forward rapidity and high $p_{\rm T}$ $/A$ of the coalescence parameter B_2 is performed using several coalescence models. Given the sensitivity of the cosmic-ray experiments discussed above, i.e., AMS-02, GAPS, and BESS-Polar, and the current experimental uncertainties on the production measurements, the direct impact of these results on cosmological DM research is related to a better understanding of the antideuteron background flux [57], which needs to be precisely modeled to interpret future measurements of antideuteron CR flux correctly.

2 Experimental apparatus

ALICE is one of the four large experiments at the LHC and is dedicated to the study of hadronic collisions at ultra-relativistic energies. A detailed description of the ALICE apparatus and its performance can be found in Refs. [58, 59]. In the following, only the sub-detector systems used for this analysis are described.

Trajectories of charged particles are reconstructed in the ALICE central barrel, which covers the pseudorapidity interval of $|\eta| < 0.9$, with the Inner Tracking System (ITS) [58], the Time Projection Chamber (TPC) [60], and the Time-Of-Flight (TOF) detector [61]. These detectors are located inside a solenoidal magnet, which generates a highly homogeneous magnetic field of 0.5 T, parallel to the beam line.

The ITS consists of six cylindrical layers of silicon detectors and is used for the determination of primary and secondary vertices, and for charged-particle tracking. The TPC is a gas detector used for charged-particle track reconstruction and momentum determination, and particle identification via the measurement of the specific energy loss (dE/dx) of particles in the detector gas. The dE/dx resolution depends on the event multiplicity and is about 5%–6.5% for minimum-ionizing particles crossing the full volume of the TPC [59]. The particle identification is extended at high $p_{\rm T}$ using the TOF detector, which is located at a radial distance of 3.7 m from the nominal interaction point. It measures the arrival time of particles relative to the event collision time provided by the TOF detector itself or by the T0 detectors. The T0 consists of two arrays of Cherenkov counters, T0A and T0C, located on opposite sides of the interaction point, covering the pseudorapidity regions $4.6 < \eta < 4.9$ and $-3.3 < \eta < -3.0$. A weighted average is performed when both T0 and TOF detectors have measured the start time [62]. The TOF time resolution is 56 ps [61].

Collision events are triggered by two plastic scintillator arrays, V0A and V0C [63], located at asymmetric positions, one on each side of the interaction point, covering the pseudorapidity regions $-3.7 < \eta < -1.7$ and $2.8 < \eta < 5.1$. Each V0 array consists of four rings in the radial direction, with each ring comprising eight cells with the same azimuthal size. The V0 detector is used to define the minimum-bias (MB) trigger (requiring coincident signals in the V0 detectors to be synchronous with the bunch-crossing time defined by the LHC clock). The V0 is also used to reject background events like beam–gas interactions, collisions with de-bunched protons, or with mechanical structures of the beam line [63].

3 Data analysis

The analyzed data sample was collected in 2016, 2017, and 2018 during the LHC pp run at $\sqrt{s} = 13$ TeV. Events with multiple vertices identified from track segments in the two innermost layers of the ITS are tagged as pile-up and removed from the analysis [59]. In order to ensure full geometrical acceptance in the ITS for $|\eta| < 0.9$ and reject background collisions, the coordinate of the primary vertex along the beam axis is required to be within 10 cm from the nominal interaction point. A total number of approximately 1.7 billion MB pp events were analyzed, with a total integrated luminosity of about 22 nb^{-1} [64].

The reconstructed tracks are required to fulfill the same set of quality criteria of the corresponding analyses of antiprotons and antideuteron production at midrapidity [27, 65], with the only difference that in the case of the present analysis no selections in pseudorapidity are applied, besides the ones enforced by the detector acceptance ($|\eta| <0.9$ in the ALICE central barrel). The antiproton (antideuteron) identification is done at $p_{\mathrm{T}} < 0.7$ GeV/c ($p_{\mathrm{T}} < 1.2$ GeV$/c$) by requiring that its energy loss per unit of track length measured by the TPC is within $3\sigma_{\mathrm{d}E/\mathrm{d}x}$ from the expected average for antiprotons (antideuterons), where $\sigma_{\mathrm{d}E/\mathrm{d}x}$ is the $\mathrm{d}E/\mathrm{d}x$ resolution. For $p_{\mathrm{T}} > 0.7$ GeV/c for antiprotons and $p_{\mathrm{T}} > 1.2$ GeV$/c$ for antideuterons, the $\mathrm{d}E/\mathrm{d}x$ signal of the TPC is complemented by the time-of-flight measured by the TOF detector. The antiproton (antideuteron) signal is extracted from a fit to the $\mathrm{n}\sigma^{\mathrm{TOF}} = (\Delta t - \Delta t_{\mathrm{p(d)}})/\sigma_{\mathrm{TOF}}$ distribution, where Δt is the measured time-of-flight, $\Delta t_{\mathrm{p(d)}}$ its expected value for protons (deuterons), and σ_{TOF} the resolution on the time-of-flight measurement. The fit function consists of a Gaussian with an exponential tail for the signal and the sum of two exponential functions for the background. The raw signal yield is extracted by integrating the signal function in the asymmetric interval $[-3\sigma_{\mathrm{TOF}} + \mu_0, +3.5\sigma_{\mathrm{TOF}} + \mu_0]$, where μ_0 is the mean of the Gaussian. The acceptance of the ITS, TPC, and TOF detectors ($|\eta| <0.9$) limits the rapidity coverage of the measurement. Hence, the analysis is performed in seven rapidity intervals, 0.2 units wide, from -0.7 to 0.7. At $|y|$ larger than 0.7, the acceptance of the TPC and TOF detectors allows for the reconstruction of deuterons only at high p_{T} (for $|y| = 0.8$, the minimum p_{T} of deuterons that can be reconstructed is ~ 3.5 GeV/c). However, in the high p_{T} region ($p_{\mathrm{T}} > 3.5$ GeV/c) deuterons cannot be identified any longer because the background due to mismatched tracks in the TOF is dominant with respect to the actual signal.

The raw antiproton and antideuteron p_{T} spectra are corrected for the acceptance and reconstruction efficiency and, only for antiprotons, for the fraction of secondary antiprotons produced by feed-down of weak decays of $\overline{\Lambda}$ hyperons. Both corrections are calculated with Monte Carlo (MC) simulations in which antinuclei are embedded into pp collision events generated using PYTHIA 8.1 with the Monash 2013 tune [66]. Antinuclei are generated with uniform p_{T} and rapidity distributions in ($0 < p_{\mathrm{T}} < 10$) GeV$/c$ and $-1 < y < 1$. The particle interactions in the experimental apparatus are simulated using GEANT 4 as transport package [67]. The acceptance $\times$ efficiency, in each p_{T} and rapidity interval, is calculated as the ratio of reconstructed and generated antiprotons and antinuclei in the simulation, where the same track selection and PID criteria as those used in the data are applied to the reconstructed sample in MC.

The fraction of secondary antiprotons from weak decays is estimated in each p_{T} and rapidity interval by fitting the distributions of the measured distance of closest approach to the primary vertex in the transverse plane (DCA_{xy}), using the template method [18]. The DCA_{xy} template distributions of primary (produced in the collision at the primary vertex) and secondary antiprotons from feed-down are taken from MC simulations. The resulting primary fraction of antiprotons has a mild p_T dependence, ranging from ∼80% at low p_T to ∼90% at high p_T.

The sources of systematic uncertainties on the transverse momentum distributions considered for this analysis are related to (i) track selection and particle identification, (ii) detector material budget, (iii) TPC–ITS and TPC–TOF track matching efficiencies, (iv) hadronic interaction of antiparticles with the detector material, and (v) signal loss (to account for antiparticles lost in rejected events). The uncertainties related to (i) are estimated for the present analysis and discussed below, while the other ones are inherited, as relative uncertainties, from previous similar analyses. The uncertainties due to (ii) and (v) are taken from Ref. [27], and those due to (iii) and (iv) are inherited from Ref. [30].

In order to assess the systematic uncertainties related to (i), the analysis is repeated using 50 settings with different track selection criteria. These criteria are defined by randomly sampling the analysis parameters from uniform probability distributions. The systematic uncertainty is determined by calculating the standard deviation of the distribution of the fully corrected yields, in each p_T and rapidity interval, similarly to what was done in Ref. [27]. For the contribution related to the particle identification, additionally, the difference between the yield extracted using the bin counting method, after the subtraction of the background counts estimated from the fit, and the integral of the fit function is included in the systematic uncertainties. In the case of antiprotons, also the contribution related to the secondary fraction of antiprotons from weak decays is included in (i), hence the uncertainty is larger for antiprotons than for antideuterons. For this uncertainty, the DCA selection criteria (both in the perpendicular plane DCA_{xy}, and along the beam axis DCA_z) are also changed 50 times, and each time the corresponding primary fraction is used for correction. The total systematic uncertainties are obtained by summing in quadrature all the individual contributions, which are summarized in Table 1. Within the studied rapidity range, the systematic uncertainties have been found to be independent of the analyzed rapidity interval.

Table 1: Summary of the systematic uncertainties on the transverse momentum distributions.

Source of uncertainty	$\overline{\mathrm{p}}$ low p_T	$\overline{\mathrm{p}}$ high p_T	$\overline{\mathrm{d}}$ low p_T	$\overline{\mathrm{d}}$ high p_T
Tracking and PID	4.5%	6%	1.5%	3%
Material budget	<1%	<1%	<1%	<1%
Matching efficiencies	3%	3%	1%	2.5%
Hadronic interaction	2%	2%	3%	4%
Signal loss	1%	1%	1%	1%
Total	6%	7%	4%	6%

4 Results

The p_T-differential yields per unit of rapidity ($\mathrm{d}^2N/\mathrm{d}p_T\mathrm{d}y$) of antiprotons and antideuterons in MB events and analyzed rapidity intervals are shown in Fig. 1. The results are consistent with those measured in $|y| < 0.5$ [27, 65]. The final transverse-momentum distributions are fitted using the Lévy-Tsallis function [68] to extrapolate the yields to the unmeasured regions. The Lévy-Tsallis fits of the spectra in the high-rapidity intervals ($|y| \geq 0.5$) are constrained to the spectrum in the lowest rapidity interval ($|y| < 0.1$), such that the normalization is always adjusted during the fit while the shape is constrained to the spectrum at low rapidity, which covers a broader p_T range. This procedure assumes that the shape of the p_T spectra is independent of rapidity, and this assumption was validated by the rather stable

χ^2/NDF values of the fits across the studied rapidity interval. In order to obtain the integrated yields, the integrals of the data points of the transverse momentum spectra in the measured regions are summed to the integrals of the fit functions at low (down to 0) and high (up to 10 GeV/c) p_T. The resulting p_T-integrated yields (dN/dy) of antiprotons and antideuterons are shown in Fig. 2. The measured integrated yields of both species show a flat trend with rapidity, up to $|y| = 0.7$. The antiproton integrated yields as a function of rapidity are compared with the predictions of two event generators, namely PYTHIA 8 [66], with two different tunes (described in the following), and EPOS 3 [69, 70]. The measured antideuteron yields are compared with the predictions of three models, based on different assumptions and making use of different event generators. Two of these models use PYTHIA 8 [66] as event generator, the third EPOS 3 [69, 70]. All models shown in Fig. 2 are normalized to data, by dividing the yields obtained with the simulations by a factor corresponding to the ratio between the measured yield and the simulated one in the lowest rapidity interval ($|y| < 0.1$), since the natural abundances of particles are not well reproduced by the event generators.

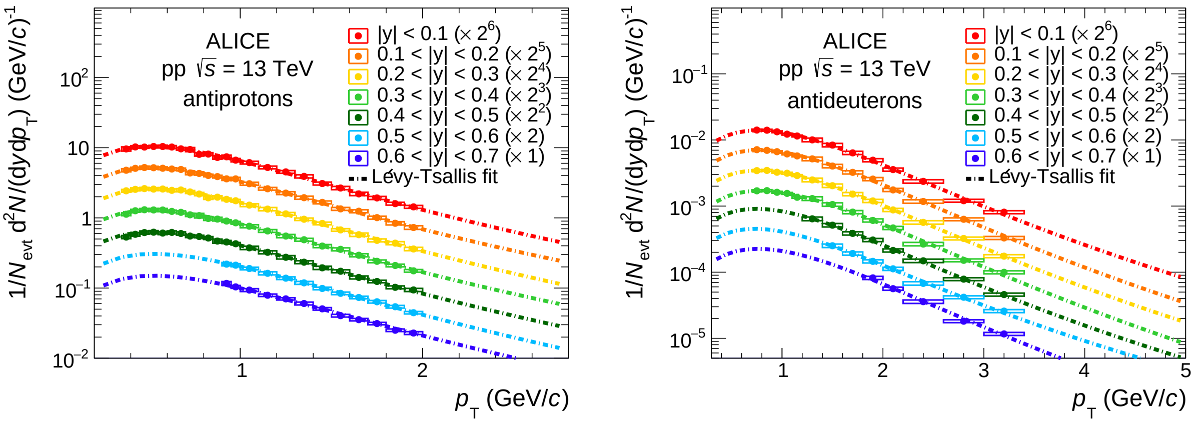

Figure 1: Antiproton (left panel) and antideuteron (right panel) p_T-differential yields for different rapidity intervals, in MB events. Statistical and systematic uncertainties are represented by vertical bars and boxes, respectively. The statistical uncertainties are smaller than the size of the markers in the reported scale and, hence, not visible. The dash-dotted lines represent the fit of the spectra executed with a Lévy-Tsallis function.

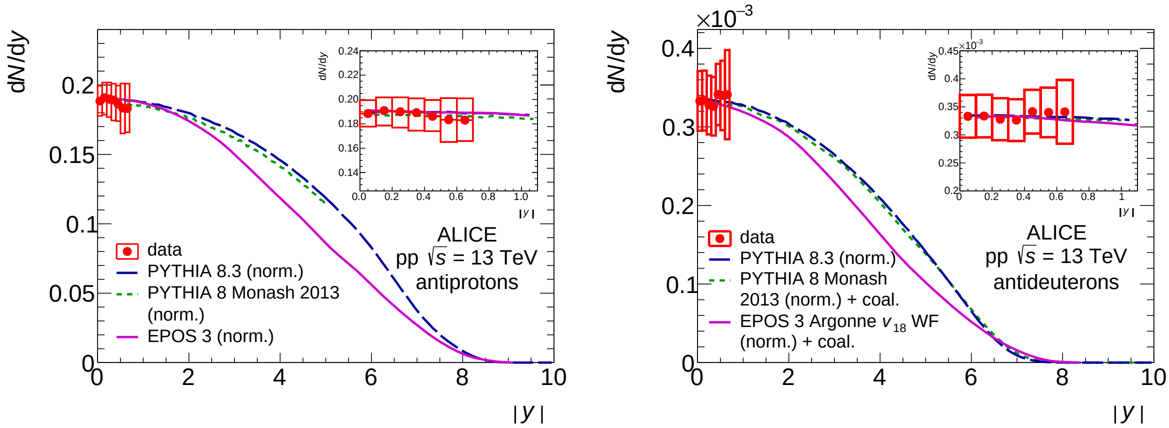

Figure 2: Integrated yields of antiprotons (left) and antideuterons (right) as a function of rapidity, compared with the corresponding predictions of three models (see text for details). Statistical and systematic uncertainties are represented by vertical bars and boxes, respectively. The statistical uncertainties are smaller than the size of the markers in the reported scale and, hence, not visible. In the insets of the figures, a zoom in the low-rapidity region is displayed. The integrated yields estimated by models are normalized to the measured ones, see text for details.

In the first approach, already used in Ref. [35] and described in detail in the corresponding Supplemental Material [71], the phase space distributions of antinucleons are generated with PYTHIA 8 with the Monash 2013 tune [66], and their $p_{\rm T}$ spectra are re-weighted to match the measured antiproton $p_{\rm T}$ spectra. In this simple model, all spatial correlations are ignored and the antideuteron is formed if an antiproton and an antineutron have a momentum difference below a given coalescence momentum $\Delta p < p_0$ in the antideuteron rest frame. The best estimate of the coalescence momentum for this model is $p_0 = (285 \pm 1)$ MeV/c [35]. The second approach is a reaction-based model [72], where antideuterons are generated by ordinary nuclear reactions between antinucleons produced in the collision with parameterized energy-dependent cross sections tuned on available experimental data [73], as implemented in the MC event generator PYTHIA 8.3. Both, the simple coalescence and the reaction-based model, are based on the simplistic assumption that coalescence can happen only if antiprotons and antineutrons are close in momentum space, while all spatial correlations are neglected, and the nucleus wave function does not come into play. The third approach is a state-of-the-art coalescence model [55] based on the Wigner function formalism, applied as an afterburner to the nucleons generated with the EPOS 3 event generator. In such a model, the quantum-mechanical properties of the nucleus are taken into account through its Wigner function. Antideuterons are formed following a probability calculated on an event-by-event basis. The coalescence model is very sensitive to the source size, i.e., the size of the particle-emitting source that can be measured using femtoscopy with two-particle correlations [74]. Since the one implemented in EPOS 3 does not reproduce the measurements [55], a parameterization of the transverse-mass ($m_{\rm T}$) dependence of the measured source size for high-multiplicity events ($r_0 = 1.249$ fm [75]) scaled to the minimum bias value ($r_0 = 1.18$ fm [76]) was used. It has to be noted that the average values of transverse mass, $\langle m_{\rm T} \rangle$, of the high multiplicity sample (corresponding to the interval of 0–0.17% in multiplicity) and of the minimum bias one (full multiplicity 0–100%) differ by $\sim 10\%$. Nonetheless, the source size is assumed to be independent of rapidity. Note also that the integrated yields predicted by the models are scaled to the data, while the shapes of the $p_{\rm T}$ distributions are not modified. Therefore, the models are no longer sensitive to the magnitude of the source size used for the calculations, but only to its $m_{\rm T}$ trend. Additionally, the predictions are obtained using the Argonne v_{18} wavefunction which was shown to give the best predictions for the antideuteron momentum distribution [55]. In this way, one obtains predictions using a realistic coalescence model, which considers not only the momentum distributions of nucleons but also their spatial correlations, as well as the quantum-mechanical nature of the coalescence process. The integrated yields of antiprotons and antideuterons as a function of rapidity obtained with the event generators and with the three models previously described show a similar trend, rather flat at low rapidities ($|y| < 1.5$) and decreasing starting from rapidity about 1.5, up to the kinematic limit of $|y| \sim 9.5$, related to the collision system and energy. However, while the predictions of the models based on PYTHIA 8 and simple coalescence are roughly equal in the full rapidity interval (0–9.5), a large difference, for rapidity larger than 1.5, is observed between these models and the third one in which the distributions of the antinucleons are obtained using EPOS 3 and a state-of-the-art model for deuteron formation via coalescence.

Using Eq. 1, the coalescence parameters B_2 in the different rapidity classes are calculated and shown in Fig. 3, in comparison with predictions from PYTHIA 8.3 (left panel) and from EPOS 3 with the coalescence model of Ref. [55] used as afterburner (right panel). In the models, the antiproton $p_{\rm T}$ distributions are reweighted to the measured ones, while the antideuterons are only scaled to the integrated measured yields, and their $p_{\rm T}$ shapes are not changed. The rising trend of the coalescence parameters with $p_{\rm T}/A$ is well reproduced by the two models. Since larger coalescence parameters correspond to smaller source sizes (at higher $p_{\rm T}$), this result reflects the observed decreasing trend of the source size with $m_{\rm T}$ [74].

Finally, selecting intervals of $p_{\rm T}/A$, it is interesting to look at the trend of the coalescence parameter as a function of rapidity. The rapidity dependence of B_2 is shown in Fig. 4 and compared with predictions from PYTHIA 8.3 (left panel) and from a coalescence model [55] used as afterburner of EPOS 3 (right panel). As for the yields, the coalescence parameter is flat in the rapidity region investigated by the measurements. The models reproduce the measured trend, with an agreement within 2σ on average.

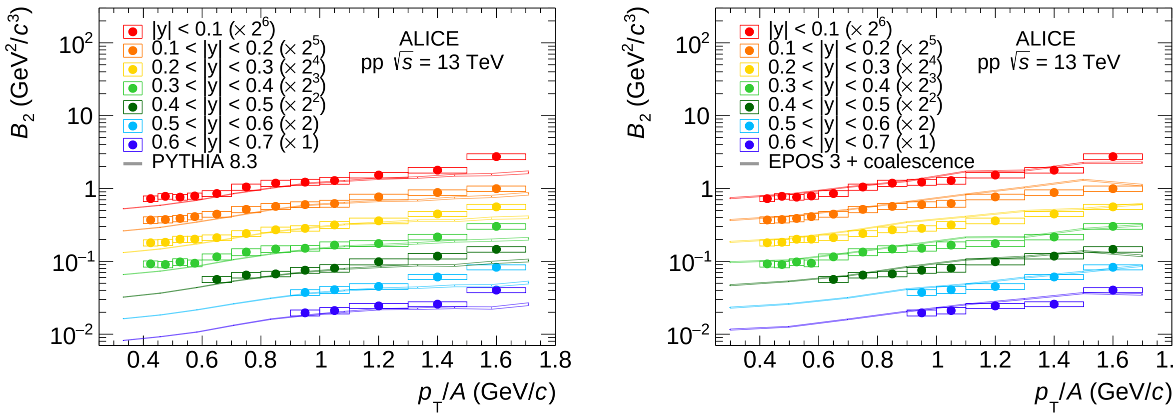

Figure 3: Coalescence parameter B_2 as a function of the transverse momentum per nucleon for different rapidity intervals. Data are compared to model predictions from PYTHIA 8.3 (left panel) and from a coalescence model [55] used as afterburner of EPOS 3 (right panel), shown as colored lines. Statistical and systematic uncertainties on the data points are represented by vertical bars and boxes, respectively. The statistical uncertainties of the data points are smaller than the size of the markers in the reported scale and, hence, not visible.

Some deviations beyond the 2σ level are present at high p_T/A (> 1.5 GeV/c) for the model based on PYTHIA 8.3. A similar trend of B_2, flat as a function of rapidity, was also observed by the BRAHMS Collaboration in central Au–Au collisions at $\sqrt{s_{NN}} = 200$ GeV [13].

The models shown in Figs. 3 and 4 have similar statistical uncertainties, of about 10%, which are reflected in the width of the colored bands.

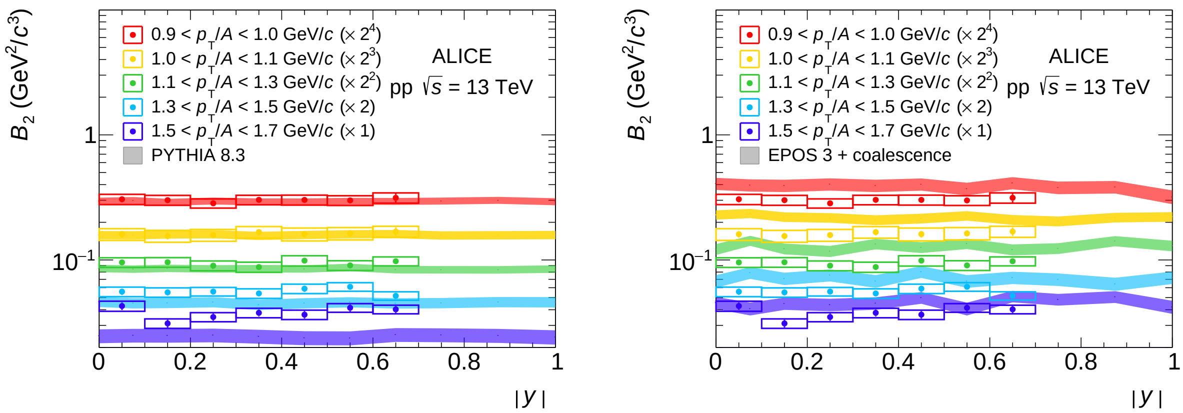

Figure 4: Coalescence parameter B_2 as a function of rapidity, for fixed transverse momentum per nucleon values. Statistical and systematic uncertainties on the data points are represented by vertical bars and boxes, respectively. The statistical uncertainties are smaller than the size of the markers in the reported scale and, hence, not visible. Bands show the model predictions from PYTHIA 8.3 (left panel) and from a coalescence model [55] used as afterburner of EPOS 3 (right panel), the width reflecting the statistical uncertainties.

The flux of antideuterons from cosmic rays as a function of rigidity, reported in Fig. 1 of Ref. [77], is predicted with the model of Ref. [4] using the coalescence parameter B_2 measured at midrapidity in pp collisions at $\sqrt{s} = 7$ TeV with ALICE [78]. The rapidity and p_T dependencies of the coalescence parameters predicted by PYTHIA 8.3 and EPOS 3, scaled to match the experimental measurements in the first p_T interval reported here, are extrapolated at high p_T ($p_T > 1.7$ GeV/c) and forward rapidity ($|y| > 0.7$) and provided to the authors of Refs. [4, 77]. In these models, the B_2 distributions are substantially flat as a function of rapidity and increase monotonically as a function of p_T /A. No difference is found in the shape of the antideuteron flux shown in Fig. 1 of Ref. [77], either using the different models presented in this Letter or using the B_2 measured at midrapidity in pp collisions at $\sqrt{s} = 7$ TeV [78]. A constant B_2 as a function of rapidity is a common assumption in several CR models. However, in the coalescence models presented in this work, the trend with rapidity of B_2 is investigated and its effect on the predictions of the CR antideuteron flux is verified explicitly, for the first time. Additionally, in Ref. [77], the author investigates the role of the rapidity extrapolation in predicting the cosmic ray flux of antinuclei (i.e., antiprotons, antideuterons, and antihelium-3), by setting the production cross section to zero in different kinematical regions. The results for the antideuteron flux from CR, presented in Fig. 1 of Ref. [77], show that at low kinetic energies, corresponding to rigidity $R \leq 5$ GV (being the rigidity defined as $R = \frac{pc}{Ze}$, where p is the momentum, c is the speed of light, and Ze is the charge of the cosmic-ray particle), about 90% of the flux is due to antideuterons with rapidity larger than 0.5. This kinematical region is also the one where the CR flux is expected to be dominated by the DM signal, according to many cosmological models [5], and hence of high impact for astrophysics. Therefore, it is important to measure the coalescence parameter in the rapidity region $|y| > 0.5$, in order to use experimental measurements as input for the astrophysical models. Moreover, the results of the antideuteron flux show that LHC experiments (such as LHCb [79–81] in its fixed-target configuration, and the future ALICE 3 facility [82]) are expected to fill the rapidity gap ($0.5 < |y| < 1.5$) providing measurements of production yields of antinuclei. These measurements would be of important use to the astrophysics community for the calculation of the flux of antinuclei from cosmic rays, to avoid relying on the model extrapolation at forward rapidities.

5 Summary

In this Letter, the p_{T}-differential yields of antideuterons and antiprotons as a function of rapidity are measured in pp collisions at $\sqrt{s} = 13$ TeV, for the first time up to $|y| = 0.7$. Using the measured spectra, the p_{T} and rapidity dependence of the coalescence parameter B_2 is explored. Both the antiproton and antideuteron integrated yields, and the B_2 are found to be independent of rapidity in the measured range. The rapidity dependence of the obtained yields and B_2 is extrapolated for $|y| > 0.7$ using phenomenological antideuteron production models based on coalescence, making use of PYTHIA and EPOS 3 as event generators. The resulting predictions of B_2 as a function of p_{T} and rapidity are used as input for the model of Ref. [4], to investigate the role of the rapidity dependence of the coalescence parameter in the predictions of the flux of antinuclei from cosmic rays [77]. No impact on the predictions of the flux of antideuterons from CRs is seen when including the flat rapidity dependence of B_2 predicted with the models presented in this Letter, with respect to the predictions that use the B_2 measured at midrapidity in pp collisions at $\sqrt{s} = 7$ TeV [78]. However, as pointed out in Ref. [77], at low kinetic energies ($E_{\mathrm{kin}} < 1$ GeV/nucleon) about 90% of the flux is due to antideuterons with rapidity between 0.5 and 1.5. Hence, the rapidity-dependent measurements of the production yields and coalescence probability of antinuclei in the region $0.5 < |y| < 1.5$ are a crucial input for CR models. These measurements have a fundamental impact on astrophysical indirect searches for dark matter, as the antinuclei produced from CR interactions are the dominant background source in the region of low kinetic energies. This measurement can be extended in the rapidity interval of interest by other facilities, present or future, such as LHCb [79, 81] and ALICE 3 [82].

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