

G04p-93/
041

GSII

**GSII-93-41
PREPRINT
Juni 1993**

**SUBTHRESHOLD ANTIPROTON PRODUCTION IN
HEAVY-ION COLLISIONS**

C. SPIELES, A. JAHNS, H. SORGE, H. STÖCKER, W. GREINER

Subthreshold Antiproton Production In Heavy Ion Collisions

C. Spieles^a, A. Jahns^{a,b}, H. Sorge^c, H. Stöcker^a, W. Greiner^a

^a Institut für Theoretische Physik

Johann Wolfgang Goethe-Universität, Frankfurt a. M., Germany

^b Gesellschaft für Schwerionenforschung, Darmstadt, Germany

^c Los Alamos National Laboratory, New Mexico, USA

February 12, 1993

Abstract

We present a RQMD calculation of antiproton yields and their momentum distribution in $Ne + NaF$ collisions at $2\text{ GeV}/u$. The antiprotons can be produced below threshold due to multi-step excitations for which meson-baryon interactions play a considerable role. In this system the annihilation probability for an initially produced antiproton is predicted to be about 65%.

Subthreshold particle production provides a useful tool to extract information about the hot high density phase of a heavy ion collision. This process is energetically forbidden in free NN -collisions. So it can give clues about cooperative contributions of many participant nucleons. Subthreshold production of \bar{p} has first been observed experimentally for $Si + Si$ at BEVALAC [1]. The yields were found to be inconsistent with calculations which take into account the Fermi-momentum of the nucleons (three orders of magnitude larger than expected)[2]. Experimental studies of subthreshold antiproton production have

been done recently with a 2.0 GeV/u and 1.8 GeV/u Ne beam from the heavy ion synchrotron (SIS) at GSI and will be continued[3]. Various mechanisms of antibaryon enrichment in general have been proposed, e.g. hadronic multi-step processes[4][5][6], strong mean field effects[7] and meson-meson interaction like $\rho\rho \rightarrow p\bar{p}$ [8] as well as string-string interaction[9] and the quark-gluon plasma [10][11]. Consequently, there is a need for learning about the weight of all those contributions. Here we want to demonstrate the effect of hadronic multi-step excitation to \bar{p} production and the annihilation rate. For our investigations we used a microscopic phase space approach – the relativistic quantum molecular dynamics (RQMD 1.07)[12] – which includes the excitation of Δ and N^* resonances, the formation and decay of colour strings and the mesonic degrees of freedom. The number of produced baryons in our model – based on the LUND scheme[13] – depends strongly on the excited mass of the resonance (baryonic resonances are treated as a string if $m > 2 GeV$, mesonic resonances if $m > 1 GeV$). So the excitation mechanism in elementary hadronic collisions is a crucial feature for the production. The new mass of an excited nonstrange baryon is chosen according to a Breit-Wigner distribution and projected on 20 different isospin $\frac{1}{2}$ and $\frac{3}{2}$ resonance states with corresponding lifetime. The average \bar{p} -formation time τ is about 1.5 fm/c (see [6]). For the $N\bar{N}$ annihilation cross section we implemented a parametrisation according to available experimental data of the elementary reaction. In our calculations the effective potentials in the RQMD are switched off (pure cascade mode). We performed the Monte-Carlo calculation in two steps, to cope with the extremely low production probabilities in subthreshold collisions: First we triggered on highly excited baryon masses above 2.8 GeV , then we rerun those events forcing a decay into an antibaryon but weighting them with the elementary production probability according to the string-fragmentation scheme.

As we have shown already for AGS-energies[6], heavy particles are not dominantly created in the first collisions of a nucleus-nucleus reaction but in secondary interactions. This looks the same at SIS-energies. RQMD calculations exhibit a scenario of cumulative scattering of participants and secondaries, a multi-step process of successive binary collisions like $\pi.N$, $\rho.N$, $\eta.N^*$ etc. as well as $\Delta\Delta$ etc.. Baryons and mesons get more and more excited and decay or

rescatter subsequently. Eventually the invariant mass of a single resonance-resonance interaction exceeds the threshold for antiproton production. Note that the mechanism responsible for the *enhancement* of the antiproton yields at AGS is mandatory to get *any* antiproton production at all for SIS-energies.

We investigated $Ne + NaF$ and $d + Cu$ reactions at $2.0 \text{ GeV}/u$. RQMD calculations show that the excitation of highly excited strings is strongly dependent on the number of participating target and projectile nucleons. The probability of e.g. $\Delta\Delta$ -interactions decreases drastically in asymmetric systems like dA. RQMD predicts an enhancement of several orders of magnitude in the antiproton yields when $Ne + NaF$ is compared to $d + Cu$.

Additionally one has to take into account the specific reabsorption rate of the antiprotons. Previous calculations by means of transport theory matched the data for $Sz + Sz$ at $2.1 \text{ GeV}/u$ within a factor of 2, if the absorption is neglected[4]. Other calculations for $Nz + Ni$ at $1.85 \text{ GeV}/u$ predict an annihilation of about a factor 1000, but these calculations can also fit the data[5]. The latter result is in striking contrast to our own investigations. Systematic studies of antiproton production and annihilation at $10 - 15 \text{ GeV}/u$ in the framework of RQMD established survival probabilities near 10% in Au+Au[14]. In $Ne + NaF$ the absorption rate is still 45%. The values vary with the number of participating nucleons, the impact parameter and –rather weakly– the beam energy. For the system $Ne + NaF$ at SIS-energies ($2 \text{ GeV}/u$) RQMD predicts an absorption rate of 65% (Fig. 1). If the antiproton formation time is set to zero, absorption increases up to 90% at $2 \text{ GeV}/u$.

Antiprotons are produced rather isotropically in momentum space. This is due to the multiple collisions that are necessary for the excitation of heavy resonances in $Ne + NaF$ at $2 \text{ GeV}/u$. This result does not change due to absorption. Therefore we predict that future experiments at finite angles will exhibit a similar transverse momentum distribution as the observed p_{\perp} distributions[3]. Figure 2 shows the rapidity integrated transverse momentum spectrum for initial and final antiprotons in minimum bias reactions of $Ne + NaF$. Strong annihilation is predicted at low p_{\perp} . This could be easily observed by experiments at finite angles. The mean p_{\perp} of the final yield ($\langle p_{\perp} \rangle = 285 \text{ GeV}/c$) is only slightly higher than that of the initial yield ($\langle p_{\perp} \rangle = 270 \text{ GeV}/c$).

Figure 3 shows the collision spectrum for $Ne + NaF$ at $1.9 \text{ GeV}/u$. The contributions of baryon-baryon (BB) and meson-baryon (MB) collisions add to 115 binary interactions in a central collision of which 80% are BB, 20% MB and $< 1\%$ MM. However, a BB collision requires a \sqrt{s} of four nucleon masses to produce an anti-baryon whereas a MB collision requires only three nucleon masses. If the collisions above the particular threshold are considered, the MB collisions contribute with 60% and BB with 40%. We do not find evidence for any antiproton production via the MM channel (especially $\rho\rho$) at SIS-energies.

In conclusion we state a significant discrepancy in the antiproton absorption rates as predicted by different microscopic models. Further experimental and theoretical studies are needed for clarification.

References

- [1] J. B. Carroll et al., Phys. Rev. Lett. 62 (1989) 1829
- [2] A. Shor et al., Nucl. Phys. A 514 (1990) 717
- [3] A. Gillitzer et al., Proc. XXX Int. Winter Meeting on Nuclear Physics, Bormio 1992 and Proc. XXXI Int. Winter Meeting on Nuclear Physics, Bormio 1993 and approved proposal extension at GSI(1993)
- [4] G. Batko, W. Cassing, U. Mosel, K. Niita, Gy. Wolf, Phys. Lett. B256 (1991) 331
- [5] S. Teis, W. Cassing, T. Maruyama, U. Mosel, GSI annual report 1992, S.133
- [6] A. Jahns, H. Stöcker, W. Greiner, H. Sorge, Phys. Rev. Lett. 68 (1992) 2895
- [7] J. Schaffner, I. N. Mishustin, L. M. Satarov, H. Stöcker, W. Greiner, Z. Phys. A341, 47 (1991)
- [8] C. M. Ko, L. H. Xia, Phys. Rev. C 40 (1989) R1118
- [9] H. Sorge, M. Berenguer, H. Stöcker, W. Greiner, Phys. Lett. B289, 6 (1992)
- [10] K. A. Olive, Phys. Lett. 89B, 299 (1980)
- [11] U. Heinz, P. R. Subramanian, H. Stöcker, W. Greiner, J. Phys. G12, 1237 (1986)
- [12] H. Sorge, H. Stöcker, W. Greiner, Ann. Phys. (N.Y.) 192 (1989) 266
- [13] B. Anderson, G. Gustafson, G. Ingelmann, T. Sjöstrand, Phys. Rep. 97(1983) 31
- [14] A. Jahns, C. Spieles, R. Mattiello, H. Stöcker, W. Greiner, H. Sorge, Phys. Lett. B in print

Figure Captions:

Figure 1

Rapidity distribution of initial (dashed line) and final antiprotons with $\tau = 0 fm/c$ (dotted line) and the default value $\tau \approx 1.5 fm/c$ (full line) in minimum bias reactions of $Ne + NaF$ at $2 GeV/u$. Initial \bar{p} 's correspond to an infinite formation time.

Figure 2

Spectrum of transverse momentum of initial (dashed line) and final (full line) antiprotons in minimum bias reactions of $Ne + NaF$ at $2 GeV/u$.

Figure 3

Collision spectrum of central $Ne + NaF$ collisions at $1.9 GeV/u$. The contribution of meson-baryon (dashed line) and baryon-baryon (full line) collisions are shown. The thresholds for antiproton production are indicated.

Figure 1

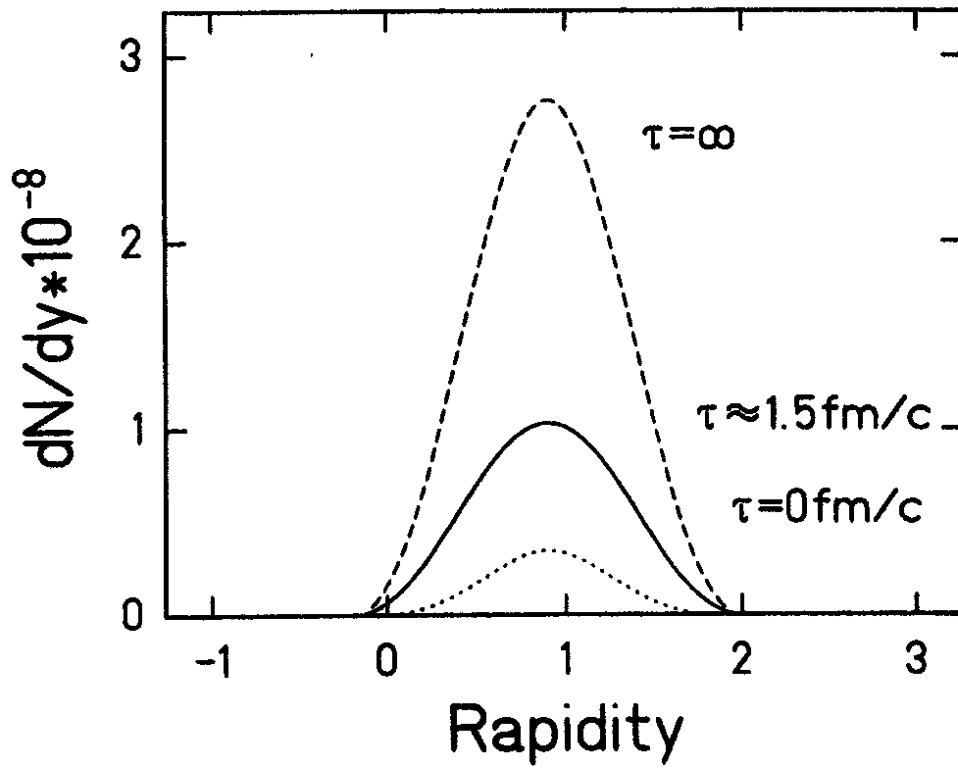


Figure 2

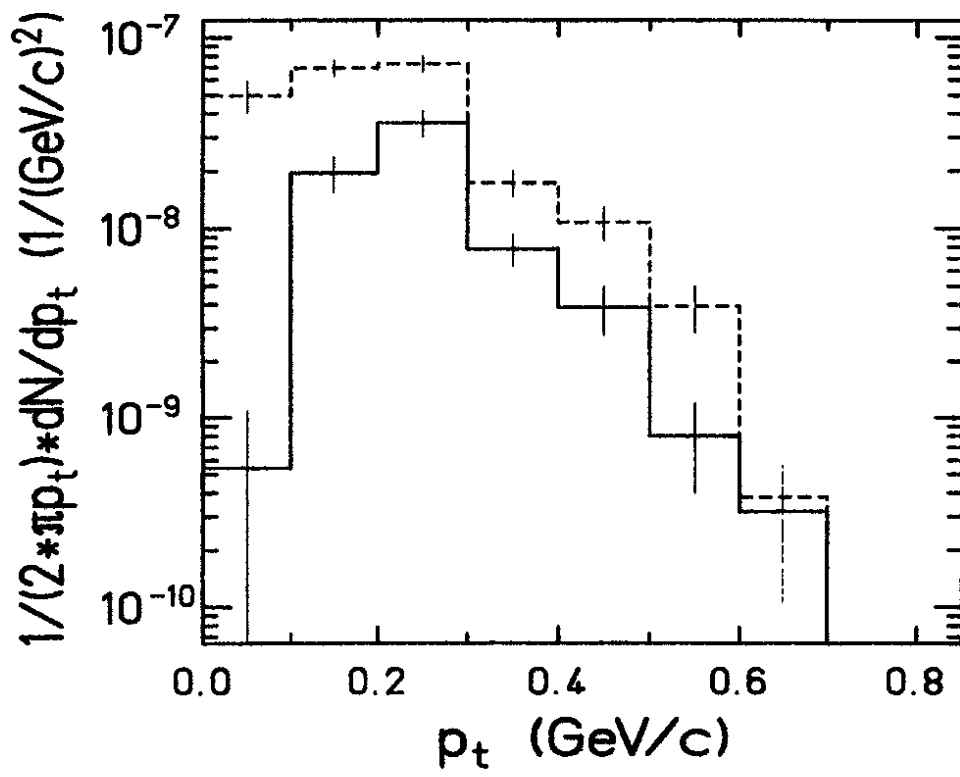


Figure 3

