



CERN-EP-2020-287
20 December 2020

Higher harmonic non-linear flow modes of charged hadrons in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV

ALICE Collaboration*

Abstract

Anisotropic flow coefficients, v_n , non-linear flow mode coefficients, $\chi_{n,mk}$, and correlations among different symmetry planes, $\rho_{n,mk}$ are measured in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV. Results obtained with multi-particle correlations are reported for the transverse momentum interval $0.2 < p_{\text{T}} < 5.0$ GeV/ c within the pseudorapidity interval $0.4 < |\eta| < 0.8$ as a function of collision centrality. The v_n coefficients and $\chi_{n,mk}$ and $\rho_{n,mk}$ are presented up to the ninth and seventh harmonic order, respectively. Calculations suggest that the correlations measured in different symmetry planes and the non-linear flow mode coefficients are dependent on the shear and bulk viscosity to entropy ratios of the medium created in heavy-ion collisions. The comparison between these measurements and those at lower energies and calculations from hydrodynamic models places strong constraints on the initial conditions and transport properties of the system.

arXiv:2002.00633v1 [nucl-ex] 3 Feb 2020

© 2020 CERN for the benefit of the ALICE Collaboration.

Reproduction of this article or parts of it is allowed as specified in the CC-BY-4.0 license.

*See Appendix B for the list of collaboration members

1 Introduction

One of the primary goals in the ultra-relativistic heavy-ion collision programs at the Relativistic Heavy Ion Collider (RHIC) and the Large Hadron Collider (LHC) is to study the nuclear matter at extreme conditions. The pressure gradients in the strongly interacting matter, known as the Quark–Gluon Plasma (QGP), are believed to drive the hydrodynamic expansion observed through anisotropy in multi-particle correlations in high energy collisions at RHIC and the LHC [1, 2]. The anisotropic expansion of the medium, commonly referred to as anisotropic flow [1], can be characterized by a Fourier decomposition of the azimuthal particle distribution with respect to the common symmetry planes [3, 4]

$$\frac{dN}{d\varphi} \propto 1 + 2 \sum_{n=1}^{\infty} v_n \cos(n(\varphi - \psi_n)), \quad (1)$$

where the flow coefficient v_n and the symmetry plane ψ_n are, respectively, the magnitude and orientation of the n -th order complex flow vector $V_n \equiv v_n e^{in\psi_n}$. Anisotropic flow measurements through two- and multi-particle azimuthal correlations [5–12] have provided important information on the medium response and in particular its transport coefficients such as the shear viscosity to entropy density ratio (η/s), bulk viscosity to entropy density ratio (ζ/s) and the equation of state [13]. Studies have shown the relativistic hydrodynamic nature of the medium [1, 2, 14–21], with η/s close to the AdS/CFT minimum $1/(4\pi)$ [22].

The initial state eccentricity, determined from the energy density profile, is obtained from the definition [23]

$$\varepsilon_n e^{in\Phi_n} = -\{r^n e^{in\varphi}\} / \{r^n\}, \quad n \geq 2, \quad (2)$$

where the curly brackets denote the average over the transverse plane, i.e. $\{\dots\} = \int dx dy e(x, y, \tau_0)(\dots)$, r is the distance to the system's center of mass, φ is the azimuthal angle, $e(x, y, \tau_0)$ is the energy density at the initial time τ_0 , and Φ_n is the participant plane angle (see Refs. [24, 25]). Hydrodynamic models demonstrate that the second and the third harmonic flow coefficients exhibit an almost linear dependence on the initial eccentricity coefficients ε_n [26]. Considering that the anisotropic expansion is a result of a hydrodynamic evolution governed by η/s , a measurement of the second and third harmonics combined with hydrodynamic calculations can constrain the properties of the medium. Several estimates for the limits of η/s were determined through measurements of elliptic flow coefficient v_2 [27–32] and their comparison with hydrodynamic calculations. Consequently, the early constraints placed the value of η/s between 0.08 to 0.16 [33–35]. However, the limited sensitivity of the elliptic flow to η/s and the large uncertainty in the initial state anisotropy inhibit a precise determination of the value of η/s [34, 36–38], and its temperature dependence, which was recently shown to be explorable during the second run of LHC [39, 40]. In addition, part of the anisotropic flow can also originate from the hadronic phase [41–43]. It has been shown in [43, 44] that the inclusion of the temperature dependent bulk viscosity $\zeta/s(T)$ in hydrodynamic simulation lead to a better description of the average transverse momentum of charged hadrons and on the elliptic flow coefficient. The effects of bulk viscosity should be considered when extracting any transport coefficient from the data [45–47].

Flow harmonics of order $n \geq 3$ reveal finer details of initial conditions [6, 7, 9, 10, 12], enabling to constrain η/s better [39, 40, 48, 49]. Higher flow harmonics $n > 3$ do not exhibit a linear response to the initial anisotropy [26] as a finite contribution is induced by the initial state anisotropy of the lower orders [50, 51]. For example, the fourth order flow vector V_4 gains contributions not only from the fourth order flow (linear flow mode), but also from the second order flow (non-linear flow mode). Starting from the V_n estimators studied in [50], the flow can be expressed as an vector sum of the linear and non-linear

modes

$$\begin{aligned}
 V_4 &= V_{4L} + \chi_{4,22} V_2^2, \\
 V_5 &= V_{5L} + \chi_{5,23} V_2 V_3, \\
 V_6 &= V_{6L} + \chi_{6,222} V_2^3 + \chi_{6,33} V_3^2 + \chi_{6,24} V_2 V_{4L}, \\
 V_7 &= V_{7L} + \chi_{7,223} V_2^2 V_3 + \chi_{7,34} V_3 V_{4L} + \chi_{7,25} V_2 V_{5L}, \\
 V_8 &= V_{8L} + \chi_{8,2222} V_2^4 + \chi_{8,233} V_2 V_3^2 + \mathcal{E}(V_{4L}, V_{5L}, V_{6L}),
 \end{aligned} \tag{3}$$

where $\chi_{n,mk}$ is called non-linear flow mode coefficient, characterizing the non-linear flow mode induced by the lower order harmonics. The high order linear component is denoted by V_{nL} , while the many higher order linear couplings are depicted by $\mathcal{E}(\dots)$ for V_8 . The V_{nL} is linearly related to a cumulant-defined anisotropy [52]

$$\varepsilon'_4 e^{i4\Phi'_4} = \varepsilon_4 e^{i4\Phi_4} + \frac{3\langle r^2 \rangle^2}{\langle r^4 \rangle} \varepsilon_2 e^{i4\Phi_2} \tag{4}$$

as opposed to the relation $v_n \propto \varepsilon_n$, where v_n is the magnitude of the total contribution and ε_n is given by Eq. (2).

In earlier measurements performed by ALICE [53], the non-linear flow mode coefficients were measured up to the sixth harmonic order in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV. It was indicated that the coefficients $\chi_{5,23}$ and $\chi_{6,33}$ are sensitive not only to η/s , but also to the distinctive energy density profiles generated by different initial conditions. It was observed that the hydrodynamic models with their respective initial conditions Monte-Carlo (MC)-Glauber [54, 55], MC-KLN [33, 56], and IP-Glasma [57]), are unable to reproduce these measurements, which indicates that the model tuning and η/s parameterization require further work.

In this paper, the measurements of high order flow coefficients in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV are presented. The flow coefficients v_n are measured up to the ninth harmonic, v_9 , extending the previous measurements of v_2 – v_6 [58]. The data recorded during the 2015 heavy-ion run of the LHC allow the measurements of non-linear flow mode and correlations between symmetry planes to be extended. A total of six non-linear flow mode coefficients are measured, including the non-linear flow mode coefficient $\chi_{7,223}$, for which the sensitivity to η/s is expected to be significantly stronger than for the lower odd-harmonic coefficient $\chi_{5,23}$ [37, 59]. The results are compared with those in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV [53] and various state of the art hydrodynamical calculations.

2 Formalism and Observables

In order to separate the linear and non-linear contributions from Eq. (3), one assumes the respective contributions to be uncorrelated [60]. For example for the fourth order V_4 , by mean-squaring the equations in Eq. (3) and setting $\langle (V_2^*)^2 V_{4L} \rangle \simeq \langle V_2^2 V_{4L}^* \rangle \simeq 0$, the linear part can be derived

$$\underbrace{\langle |V_{4L}|^2 \rangle}_{v_{4L}}^{\frac{1}{2}} = \left(\underbrace{\langle |V_4|^2 \rangle}_{v_4^2} - \underbrace{\chi_{4,22}^2 \langle |V_2|^4 \rangle}_{v_{4,\text{NL}}^2} \right)^{\frac{1}{2}}. \tag{5}$$

Here $\langle \rangle$ denotes an average over all events and $*$ the complex conjugate. The magnitudes of the linear and non-linear flow coefficients are denoted with v_{4L} and $v_{4,\text{NL}}$, respectively.

The observables of the non-linear response mode are constructed from the projections of flow vectors on to the symmetry planes of lower harmonics [61, 62]. For $n = 4$, the magnitude of the non-linear response mode is given by

$$v_{4,22} = \frac{\Re \langle V_4 (V_2^*)^2 \rangle}{\sqrt{\langle |V_2|^4 \rangle}} \approx \langle v_4 \cos(4\psi_4 - 4\psi_2) \rangle, \tag{6}$$

where $v_{4,22}^2 \equiv v_{4,\text{NL}}^2 \equiv \chi_{4,22}^2 \langle |V_2|^4 \rangle$. The right-hand side approximation holds if the low ($n = 2, 3$) and high order flow is weakly correlated. Only the fourth harmonic is shown here and the complete list of other harmonics are provided in Appendix A.

The contributions from short-range correlations unrelated to the common symmetry plane, commonly referred to as “non-flow”, are suppressed by using the subevent method where the event is divided into two subevents separated by a pseudorapidity gap [4]. The underlying multi-particle correlation coefficient for subevent A is $v_{4,22}^{\text{A}} = \langle \langle \cos(4\varphi_1^{\text{A}} - 2\varphi_2^{\text{B}} - 2\varphi_3^{\text{B}}) \rangle \rangle / \langle \langle \cos(2\varphi_1^{\text{A}} + 2\varphi_2^{\text{A}} - 2\varphi_3^{\text{B}} - 2\varphi_4^{\text{B}}) \rangle \rangle^{1/2}$ as determined using Eq. (6),¹ and a similar treatment is applied for the subevent B, for which $v_{4,22}^{\text{B}}$ is obtained by swapping B for A in the aforementioned expression. The final result is then the average of the results from subevents A and B.

The symmetry plane correlations are defined as the ratio between the magnitude of the non-linear flow modes and flow coefficients [63]. For $n = 4$, one obtains

$$\rho_{4,22} = \frac{v_{4,22}}{v_4} \approx \langle \cos(4\psi_4 - 4\psi_2) \rangle. \quad (7)$$

A value close to zero indicates weakly correlated symmetry planes, while a value reaching one implies a strong correlation. The correlations between symmetry planes reflect those of the corresponding initial state participant planes [53, 64], therefore providing valuable information on the evolution of the QGP. Correlations between symmetry planes have been previously studied using the event-plane method [64, 65]. However, these results depend on the event-plane resolution [66], which complicates the comparison between data and theoretical calculations. Recently, the ALICE Collaboration has measured symmetry-plane correlations [53]. It was found that the correlations of symmetry planes of higher harmonics with second and third order symmetry planes increased for less central collisions. Furthermore, the comparison with hydrodynamic calculations revealed the importance of final-state collective dynamics in addition to the initial-state density fluctuations [33] as it is known that the observation of correlated final state symmetry planes implies the existence of fluctuations in the initial state eccentricity vectors.

The fourth non-linear flow mode coefficient, with the aforementioned assumptions, is given by [59]

$$\chi_{4,22} = \frac{v_{4,22}}{\sqrt{\langle v_2^4 \rangle}}. \quad (8)$$

3 Experimental Setup and Data Analysis

The data sample consists of about 42 million minimum bias Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV, recorded by ALICE [67, 68] during the 2015 heavy-ion run at the LHC. Detailed descriptions of the detector can be found in [67, 69, 70]. The trigger plus crossing of beam is provided by signals from the two scintillator arrays, V0A and V0C [67, 71], covering the pseudorapidity intervals $2.8 < \eta < 5.1$ and $-3.7 < \eta < -1.7$, respectively. A primary vertex position less than 10 cm in beam direction from the nominal interaction point is required. Pile-up events are removed by correlating the V0 multiplicity with the multiplicity from the first Silicon Pixel Detector (SPD) [67, 72] layer. To further remove pile-up events, information from the Time-of-Flight (TOF) [73] detector is used: the multiplicity estimates from the SPD are correlated with those imposed with a TOF readout requirement. The centrality of the collision is determined using information from the V0 arrays. Further details on the centrality determination in ALICE are given in [74]. Only collisions with centrality ranges between 0% to 60% are used in the analysis.

¹For practical usage, the self-correlation is recursively removed from three- and four-particle correlations, resulting in modified equations.

The track reconstruction is based on combined information from the Time Projection Chamber (TPC) [67, 75] and the Inner Tracking System (ITS) [67, 72]. To avoid contributions from secondary particles, the tracks are required to have a distance of closest approach to the primary vertex of less than 3.2 cm and 2.4 cm in the longitudinal and transverse directions, respectively. Furthermore, each track is required to have at least 70 TPC space points out of the maximum 159, and the average χ^2 per degree of freedom of the track fit to the TPC space points to be less than 2. In order to counteract the effects of track reconstruction efficiency and contamination from secondary particles [76], a HIJING simulation [77, 78] with GEANT3 [79] detector model is employed to construct a p_T -dependent track weighting correction. The track reconstruction efficiency is approximately 65% at $p_T = 0.2$ GeV/ c and 80% at $p_T > 1.0$ GeV/ c , while the contamination from secondaries is less than 10% and 5%, respectively. Only particle tracks within the transverse momentum interval $0.2 < p_T < 5.0$ GeV/ c and pseudorapidity range $0.4 < |\eta| < 0.8$ are considered. A pseudorapidity gap $|\Delta\eta| > 0.8$ is used to suppress the non-flow. The observables in this analysis are measured with multi-particle correlations obtained using the generic framework for anisotropic flow analysis [80].

4 Systematic Uncertainties

The systematic uncertainties are estimated by varying criteria for selecting the events and tracks. The systematic evaluation is done by independently varying the selection criteria, and the results given by this variation are then compared to the default criteria given in Sec. 3. The total uncertainty is obtained by assuming that the individual sources are uncorrelated, which are then quadratically summed.

Summaries of the systematic uncertainties are given in Tabs. 1–4. Uncertainties stemming from the event selection criteria are estimated by changing the rejection based on the vertex position from 10 cm to 8 cm and by adjusting the pile-up rejection criteria. It is found that the contribution to the uncertainty is generally negligible, below 1%. An alternative centrality determination is employed using the event multiplicity estimates from the SPD layers. The uncertainty related to the centrality determination is less than 2% for all observables, except for v_7 to v_9 for which the uncertainty increases to 10%.

Table 1: Systematic uncertainties of the flow coefficients. The uncertainties are given in percents and are categorized into four groups: event selection, centrality determination, tracking and non-flow. The overall systematic uncertainty is obtained by summing in quadrature the uncertainties from each source.

Type	v_2	v_3	v_4	v_5	v_6	v_7	v_8	v_9
Event Selection								
z -vertex cut	< 0.1	< 0.1	< 0.1	0.5	1.2	1.6	1.8	1.7
Pile-up rejection	< 0.1	< 0.1	< 0.1	0.2	0.8	1.3	1.7	2.0
Centrality Determination								
SPD	0.6	0.3	0.3	1.1	3.9	6.6	9.1	11.5
Tracking								
Magnetic field polarity	0.1	0.1	1.7	2.4	4.1	6.8	10.5	15.2
Tracking mode	0.1	0.2	< 0.1	2.4	5.4	7.2	7.6	6.8
Number of TPC space points	0.7	1.2	1.4	1.5	1.6	1.7	1.7	1.8
Space charge distortion	< 0.1	< 0.1	< 0.1	0.2	0.7	1.2	1.7	2.3
Non-flow								
Charge combinations (−−/++)	1.1	0.7	0.8	2.9	6.2	9.3	12.3	15.2
Overall	1.5	1.4	2.4	4.9	10.3	15.4	20.4	25.6

The ALICE detector can be operated with either positive or negative solenoid magnetic field polarity. The polarity of the field affects the direction of the charged particle curvature, while also subjecting the structural materials of the detector itself to either a positive or negative magnetic field. The default data set is composed of events recorded with both polarities. The results produced with exclusively either

Table 2: Relative systematic uncertainties of the harmonic projections $v_{n,mk}$.

Type	$v_{4,22}$	$v_{5,23}$	$v_{6,222}$	$v_{6,33}$	$v_{6,24}$	$v_{7,223}$
Event Selection						
z-vertex cut	0.1	0.1	0.2	0.3	0.2	0.1
Pile-up rejection	< 0.1	0.1	0.4	0.5	0.4	< 0.1
Centrality Determination						
SPD	1.5	0.7	0.3	0.3	0.7	1.4
Tracking						
Magnetic field polarity	0.5	0.5	1.9	3.2	4.4	5.5
Tracking mode	0.1	0.4	1.4	1.7	1.1	< 0.1
Number of TPC space points	3.8	2.3	1.5	1.4	2.1	3.5
Space charge distortion	0.2	0.1	1.8	4.0	6.7	9.9
Non-flow						
Charge combinations (––/++)	4.2	4.7	5.8	7.4	9.6	14.3
Overall	5.9	5.3	6.7	9.3	12.7	18.6

Table 3: Relative systematic uncertainties of the symmetry plane correlations $\rho_{n,mk}$.

Type	$\rho_{4,22}$	$\rho_{5,23}$	$\rho_{6,222}$	$\rho_{6,33}$	$\rho_{6,24}$	$\rho_{7,223}$
Event Selection						
z-vertex cut	0.1	0.3	0.1	0.2	0.8	2.5
Pile-up rejection	0.1	0.3	0.1	0.3	1.0	2.2
Centrality Determination						
SPD	0.9	0.3	0.7	0.9	1.2	1.5
Tracking						
Magnetic field polarity	< 0.1	1.8	6.8	10.1	13.8	18.0
Tracking mode	0.1	0.3	0.8	2.6	6.1	11.2
Number of TPC space points	< 0.1	0.7	0.1	< 0.1	1.0	2.8
Space charge distortion	0.2	0.2	1.5	3.5	6.7	11.1
Non-flow						
Charge combinations (––/++)	3.1	3.6	3.6	5.6	8.7	12.9
Overall	3.3	4.2	7.9	12.4	18.8	27.5

Table 4: Relative systematic uncertainties of the non-linear flow mode coefficients $\chi_{n,mk}$.

Type	$\chi_{4,22}$	$\chi_{5,23}$	$\chi_{6,222}$	$\chi_{6,33}$	$\chi_{6,224}$	$\chi_{7,223}$
Event Selection						
z-vertex cut	< 0.1	0.1	0.3	0.3	0.3	0.1
Pile-up rejection	< 0.1	0.1	0.5	0.6	0.5	0.1
Centrality Determination						
SPD	0.2	0.6	1.0	1.0	0.7	0.1
Tracking						
Magnetic field polarity	0.6	0.2	2.5	4.1	5.1	5.5
Tracking mode	< 0.1	0.2	1.4	1.7	1.2	0.2
Number of TPC space points	< 0.1	0.2	0.5	0.7	0.9	1.1
Space charge distortion	0.2	0.1	1.9	4.4	7.1	10.1
Non-flow						
Charge combinations (––/++)	0.2	1.5	7.7	12.0	14.4	15.0
Overall	0.7	1.7	8.5	13.6	17.0	19.0

negative or positive magnetic field configurations deviate from the default by up to 15% in case of flow coefficients, and 28% for $\rho_{7,223}$. In order to estimate the non-flow contributions from resonance decays,

the like-sign technique [2] which correlates exclusively either positively or negatively charged particles, is employed. The difference with respect to the results obtained by selecting all charged particles is assigned as a systematic uncertainty. In general, this uncertainty ranges from 2% to 15%. The effect from the space charge distortions in the TPC drift volume because of the higher interaction rates is estimated by comparing results from different regions of the TPC, one for $\eta > 0$ and the other $\eta < 0$. The maximum uncertainty is evaluated less than 15%. The track reconstruction related uncertainty, referred to as tracking mode, is evaluated by comparing the results obtained with tracks for which the requirement for the number of hits in the ITS layers is changed. In this case, the uncertainty is generally less than 15%, and a maximum 20% is evaluated for $\rho_{7,223}$. Furthermore, the track selection criteria is tightened by increasing the minimum number of the TPC space points from 70 to 90, resulting in uncertainties around 1% to 3%.

5 Results

In this section, the measurements of the flow coefficients, the non-linear modes, symmetry plane correlations and the non-linear flow mode coefficients are presented. They are compared with hydrodynamic calculations with various settings [25, 57, 81, 82]. The first calculation is based on an event-by-event viscous hydrodynamic model with EKRT initial conditions [25, 81] using a value of $\eta/s = 0.20$ (param0) and a temperature dependent $\eta/s(T)$ (param1). For both configurations, ζ/s is set to zero. The visualization of the model parameters can be found in Fig. 1. The second calculation employs the iEBE-VISHNU hybrid model [83] with AMPT [84–86] and T_RENTo [87] initial conditions. The $\eta/s = 0.08$ and $\zeta/s = 0$ are taken for param2, while the $\eta/s(T)$ and $\zeta/s(T)$ (param3), extracted using Bayesian analysis [45] (except for the normalization factors) from a fit to the final multiplicities of the charged hadron spectra in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$, are used for the T_RENTo initial conditions. The third calculation uses the MUSIC model [88] with IP-Glasma [89] initial conditions with a value of $\eta/s = 0.095$ and $\zeta/s(T)$ (param4). Each of the $\eta/s(T)$ parameterizations is adjusted to reproduce the measured charged hadron multiplicity, the low- p_{T} region of the charged-hadron spectra, and v_n from central to mid-peripheral collisions up to the fourth harmonic at RHIC and the LHC [25, 44, 57, 84, 90–92]. The model configurations are summarized in Tab. 5.

In Fig. 2, the measurements of the flow coefficients from v_2 to v_9 are presented. The first two coefficients up to v_6 have been extensively measured at RHIC and LHC [5–12], and more recently also v_7 [49]. The present analysis reports the first results on higher harmonic coefficients from v_7 to v_9 in ALICE, where v_8 and v_9 are measured for the first time at the LHC energies. The coefficients exhibit a modest centrality dependence, and their magnitude is similar to that of v_7 within statistical and systematic uncertainties.

Figure 2 also shows the comparison between the measured v_n and model calculations. The hydrodynamic calculations qualitatively reproduce the v_n measurements, and the overall model depiction is very good for v_2 and v_3 . For $n \geq 4$ however, the calculations show noticeable overestimations, especially in mid-peripheral collisions. For v_5 and v_6 , the data are well described by EKRT+param0, showing a better agreement than the temperature dependent EKRT+param1. The data are also described by AMPT+param2, for which the agreement for v_5 and v_6 is good in mid-central and mid-peripheral collisions. IP-Glasma+param4 and T_RENTo+param3 overestimate the measurements by a factor of 1.5~2. Values of v_7 are well estimated by AMPT+param2, and v_8 by both AMPT+param2 and T_RENTo+param3 within uncertainties. In other cases, the data are overestimated by the other models.

To study the dependence on the harmonic orders of the anisotropy coefficients [97], Fig. 3 shows values of different coefficients as a function of the order of the harmonic, for all centralities. This presentation is particularly well suited in visualizing the harmonic dependence, and the acoustic scaling [97] observed across the harmonic orders. The decrease in v_n with increasing harmonic order up to $n = 7$ indicates viscous damping [97]. This means that the higher frequency waveform propagating through the medium

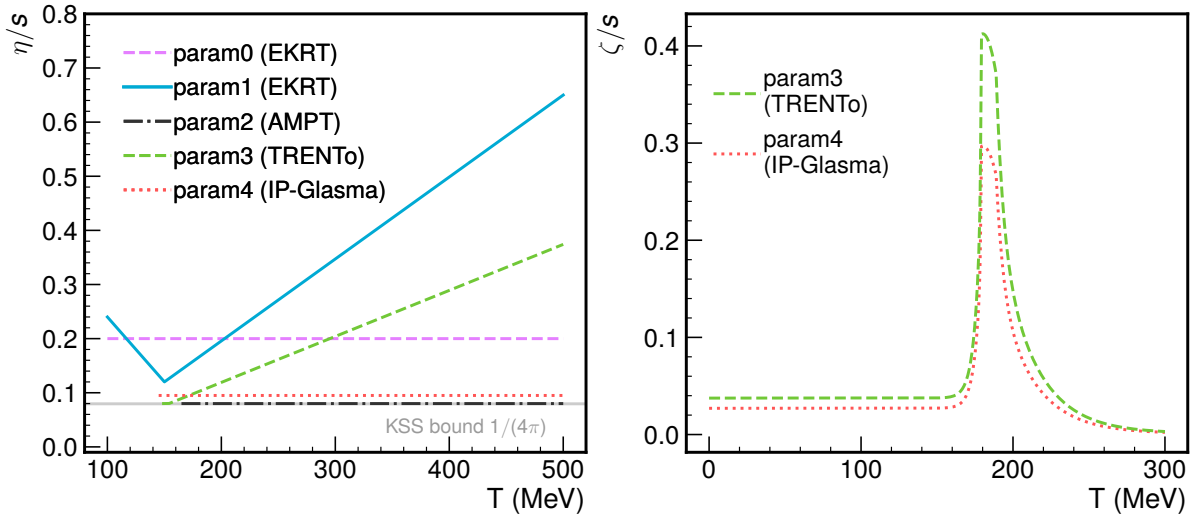


Fig. 1: The five different parameterizations of η/s and ζ/s used for the different hydrodynamic model calculations are shown in the left and right panel. Note that the functional form of $\zeta/s(T)$ is the same for param3 and param4 and taken from Eq. 5 in [45] motivated by Refs. [43, 93–95]. For the parameters with $T_{\text{R}}\text{ENTo}$ initial condition, the ones based on identified yields are taken from Table 4 in [45]. The ζ/s normalization factor used with IP-Glasma ($T_{\text{R}}\text{ENTo}$) initial conditions is 0.9 (1.25). The models with $\zeta/s = 0$ are not shown on the right.

Table 5: Hydrodynamic model configurations. Shown are the key components such as initial condition models, and η/s and ζ/s parameterizations. With $T_{\text{R}}\text{ENTo}$ initial conditions, an entropy deposition parameter $p = 0$ [82] is used for all calculations.

Model	Hydrodynamic code	Initial conditions	η/s	ζ/s
EKRT+param0 [25, 81]	EbyE [25, 96]	EKRT [25, 81]	0.20	0
EKRT+param1 [25, 81]	EbyE [25, 96]	EKRT [25, 81]	$\eta/s(T)$ [25]	0
AMPT+param2 [82]	iEBE-VISHNU [83]	AMPT [84–86]	0.08	0
$T_{\text{R}}\text{ENTo}$ +param3 [82]	iEBE-VISHNU [83]	$T_{\text{R}}\text{ENTo}(p = 0)$ [87]	$\eta/s(T)$ [45, 82]	$\zeta/s(T)$ [45, 82]
IP-Glasma+param4 [57]	MUSIC [88]	IP-Glasma [89]	0.095	$\zeta/s(T)$ [57]

should get more damped until freeze-out takes place. In [98, 99] the viscosity effect is explained as the main contributor to the observed damping. It is speculated, that another driving factor is the phase of the oscillation itself, which also contributes to the magnitude at the time of freeze-out. The measurements show that there is a hint of $v_9 > v_8$, as also predicted in the acoustic model [97].

Figure 4 presents the higher order flow coefficients as well as their linear and non-linear flow modes up to the seventh order as a function of centrality. For the flow harmonics v_4 and v_5 , presented in panels (a) and (b), respectively, the non-linear contributions are small in central collisions, where the linear contribution is dominant. A weak centrality dependence is observed for the linear component. In case of v_4 , significant contributions from the second order arise in less central collisions. The v_5 coefficient, on the other hand, is largely induced by the low order v_2 and v_3 , indicated by the large $v_{5,23}$.

Panels (c) and (d) of Fig. 4 show the flow modes of v_6 and v_7 . Only the non-linear flow modes of v_6 and v_7 are presented. The $v_{6,222}$ increases from zero to approximately half of the total v_6 in mid-central collisions, while the other mode, $v_{6,33}$, has a much weaker centrality dependence. The relatively large magnitude of these flow modes imply strong contributions from the second and third order harmonics. Finally, $v_{6,24}$ follows the trend of the total magnitude. The magnitude of $v_{6,24}$ comes close to the total, which in turn suggests not only strong contributions from the second harmonic order, but also the fourth one. The $v_{6,24}$ induced by the fourth order is seen to be the dominant contribution to the sixth order from 10% centrality classes and higher. For the seventh order v_7 , there are three non-linear contributions, of which $v_{7,223}$ is measured. The centrality dependence is similar as with the v_6 coefficient, and there is a

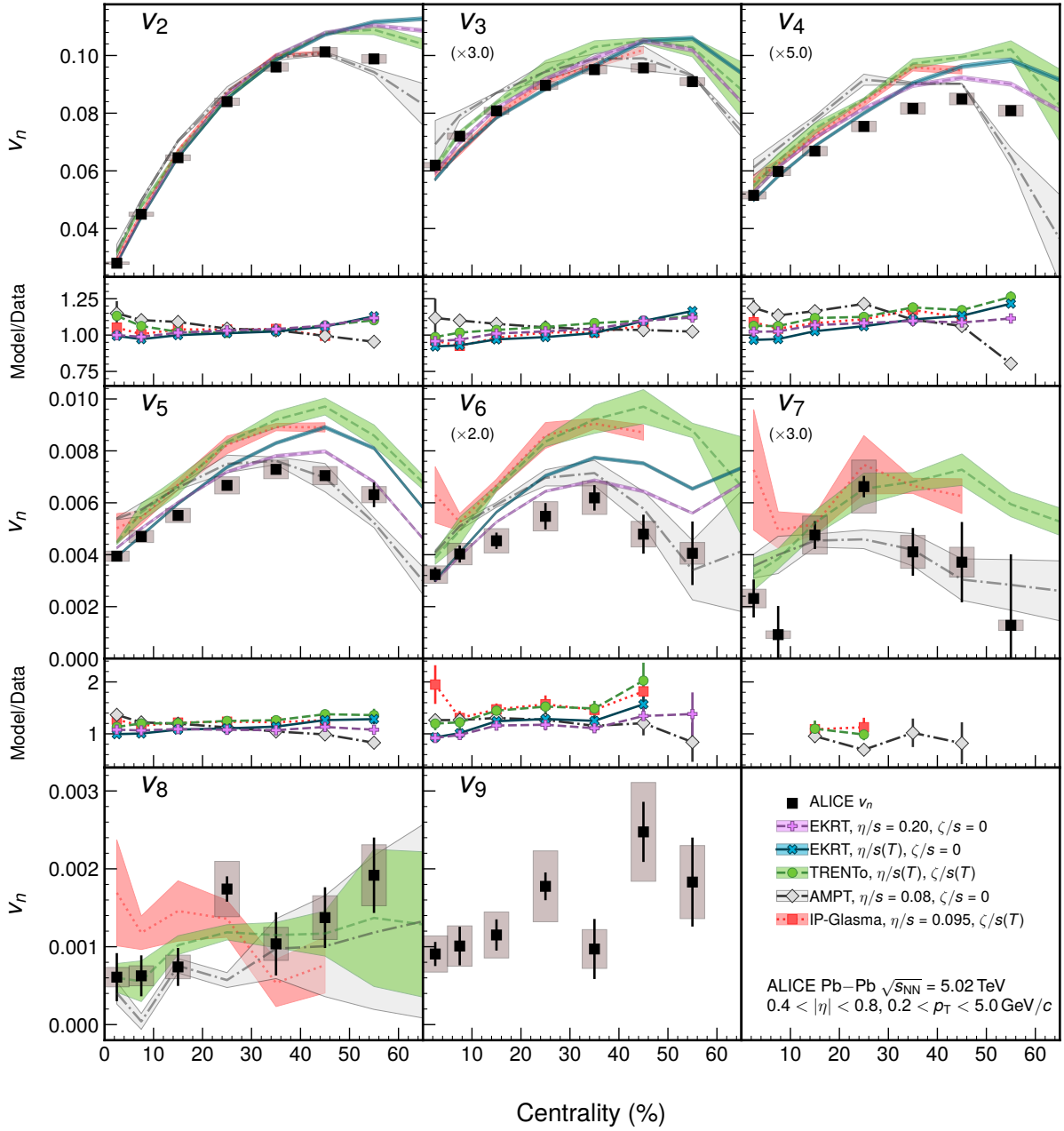


Fig. 2: Flow harmonics up to the ninth order as a function of centrality, along with five different hydrodynamic calculations shown as color bands, each representing different configurations. For the black markers representing the measured data points, the systematic uncertainty is indicated by the gray patches around the markers. The bands indicate the extent of the uncertainty of the corresponding calculation. On the bottom part of each panel, the ratios between model calculations and the data are shown with symbols. Ratios with uncertainties larger than 1 are not shown in the ratio panel. For some panels, the points are scaled by an indicated factor for better visibility across the panels.

similar general trend as for the lower order harmonics among the non-linear flow modes.

The coefficients $\rho_{n,mk}$, quantifying the correlations amongst different symmetry planes, are presented as a function of centrality in Fig. 5. Except for $\rho_{6,33}$, all coefficients indicate an increase in correlation between symmetry planes with increasing centrality class of the collision. The measurements generally agree with the ones obtained at the lower energy. The $\rho_{6,222}$ is the only coefficient for which an energy

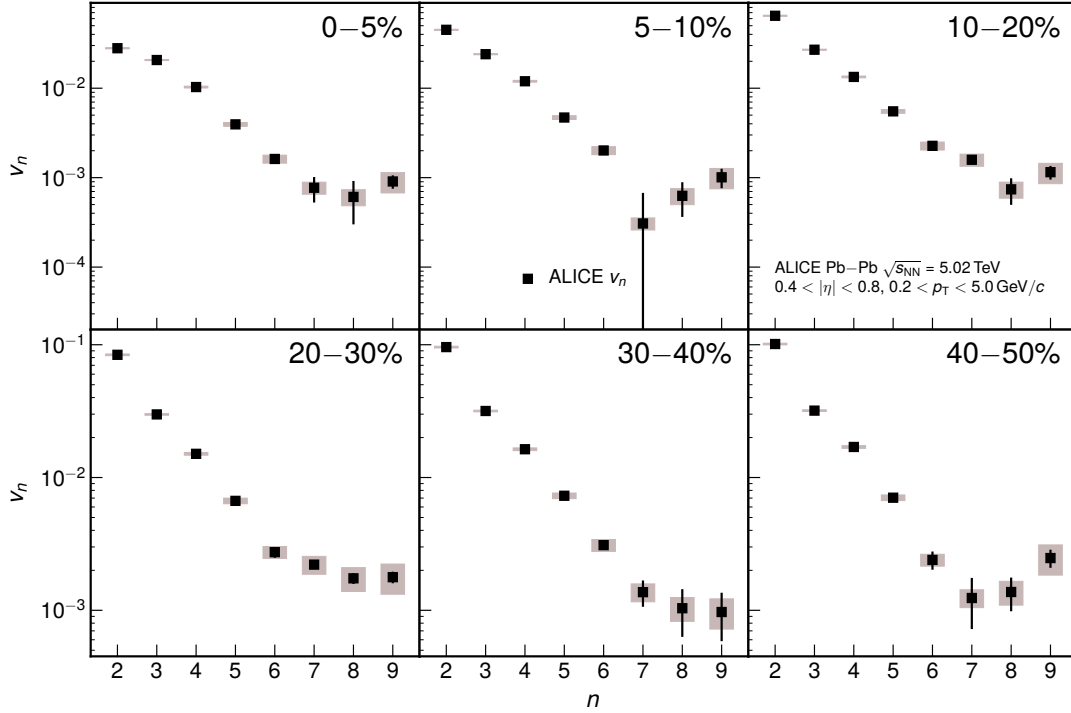


Fig. 3: v_n as a function of the harmonic order n for various centrality intervals.

dependence can be observed. The hydrodynamic calculations reproduce the measurements within the large theoretical uncertainties. For $\rho_{4,22}$, $\rho_{5,23}$, and $\rho_{6,222}$, `TRnTO+param3` however underestimates the data in mid-central collisions.

Finally, the non-linear flow mode coefficients are presented in Fig. 6. Six coefficients are measured, of which four are compared with the lower beam energy results available in [53]. For $\chi_{4,22}$ and $\chi_{5,23}$, the centrality dependence and overall magnitude agree well with the results from the lower beam energy. The centrality dependence of the new data is similar to the previous results: a larger value in more central collisions, decreasing close to unity towards 50% centrality.

All of the non-linear flow mode coefficients for the sixth harmonic agree with the previous measurements. The centrality dependence of $\chi_{6,222}$ is similar to the ones of the lower order coefficients, and the overall magnitude similar to $\chi_{4,22}$. As for $\chi_{6,33}$, no clear centrality dependence is observed within the current experimental uncertainties. Whereas the previous measurements are unable to distinguish between the magnitudes of $\chi_{6,222}$ and $\chi_{6,33}$, the current results show that $\chi_{6,222} > \chi_{6,33}$ across the whole centrality interval. For $\chi_{7,223}$, the overall magnitude is larger than for the other non-linear flow mode coefficients.

The hydrodynamic calculations for the non-linear flow mode coefficients show slightly more variation compared to the symmetry-plane correlations. As seen from the panels of Fig. 6, one observes the reproduction of the data points by `EKRT+param0` up to the modes of the sixth harmonic, and `TRnTO+param3` in all harmonics. The `EKRT+param1` calculations slightly overestimate the centrality dependence of the non-linear flow mode coefficients. It can be seen that the parameterizations of the EKRT presented here imply $\chi_{n,mk}$ across all harmonic orders to have sensitivity to η/s , whereas in the previous calculations in [53], weak η/s dependence was found for $\chi_{4,22}$ and $\chi_{6,222}$. The fifth order coefficient $\chi_{5,23}$ is expected to be quite sensitive to η/s in central collisions as can be seen from the difference of the predicted values from `EKRT+param0` and `EKRT+param1`. The `AMPT+param2` calculations underestimate the magnitude of some of the measured non-linear flow mode coefficients in various centrality classes, especially $\chi_{5,23}$, $\chi_{7,223}$ as well as $\chi_{6,24}$. The `IP-Glasma+param4` calculations overestimate the measurements in some centrality intervals.

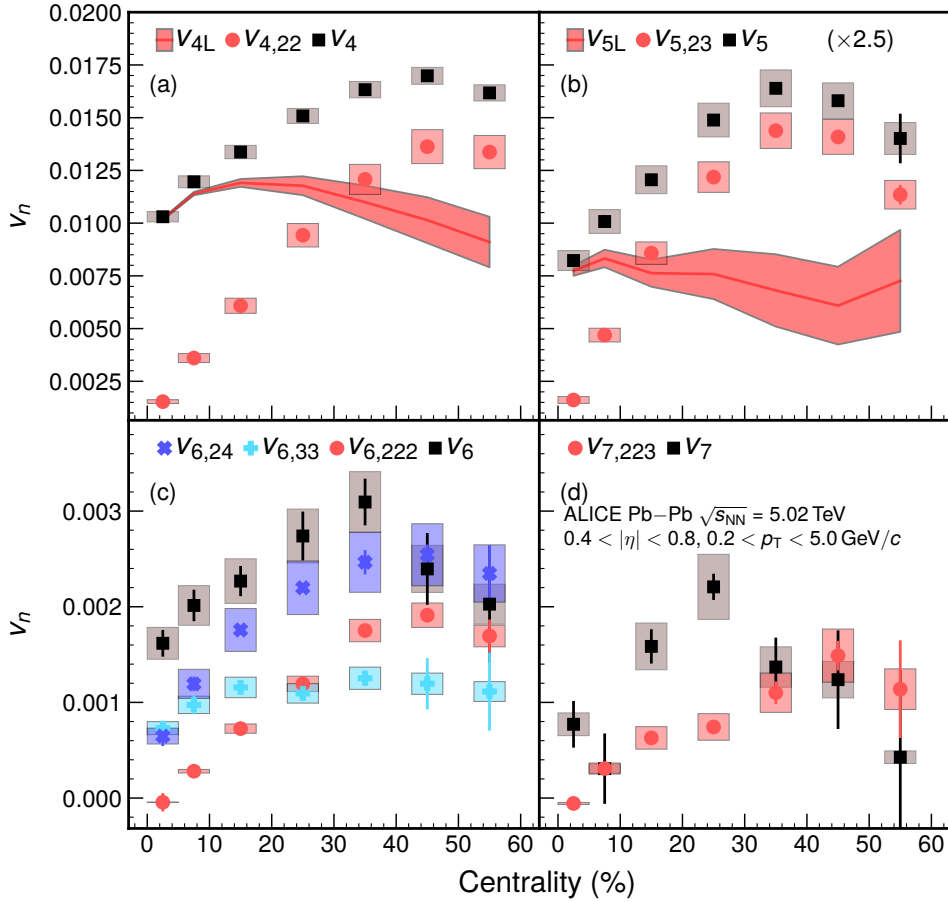


Fig. 4: Linear and non-linear flow modes as a function of centrality. The total contribution measured in Pb–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV is shown as black squares. Various non-linear contributions are presented in different red and blue colors, while the linear part, extracted from the aforementioned contributions, is shown as a red band. For panel (b), the data points are scaled by 2.5 for better visibility across the panels.

The agreement between data and the model calculations is quantified by calculating the reduced χ^2/N_{dof} defined as

$$\chi^2/N_{\text{dof}} = \frac{1}{N_{\text{dof}}} \sum_{i=1}^{N_{\text{dof}}} \frac{(y_i - f_i)^2}{\sigma_i^2}, \quad (9)$$

where y_i and f_i are the values for data and calculations, respectively, and $\sigma_i^2 = \sigma_{i,\text{stat}}^2 + \sigma_{i,\text{syst}}^2 + \sigma_{f_i,\text{stat}}^2$ is the quadratic uncertainty in terms of statistical measurement $\sigma_{i,\text{stat}}$, model uncertainties $\sigma_{f_i,\text{stat}}$, and systematic uncertainties $\sigma_{i,\text{syst}}$ in centrality bin i . Here N_{dof} represents the number of data points across the centrality interval.

The χ^2/N_{dof} for the flow coefficients are presented in Fig. 7, panel (a). It is observed that IP-Glasma+param4 gives the best description of v_2 and v_3 compared to the other models, indicated by the overall low value of χ^2/N_{dof} . However, the overall performance of IP-Glasma+param4 is considerably worse at $n \geq 4$, for which the data are overestimated, as seen in Fig. 2. For v_4 to v_6 , EKRT+param0 gives the lowest value of χ^2/N_{dof} . In the case of EKRT+param1, the χ^2/N_{dof} is slightly worse than EKRT+param0. The χ^2/N_{dof} of TRenTo+param3 is very close to that of IP-Glasma+param4, indicating a comparable description of data between the two model configurations. At low harmonic orders, TRenTo+param3 performs slightly worse than IP-Glasma+param4. For $n \geq 4$, description of the data between these two models are comparable except for $n = 8$, where TRenTo+param3 clearly has a better magnitude and centrality depiction. Notably this can be seen for v_8 where the χ^2/N_{dof} value is the lowest across all

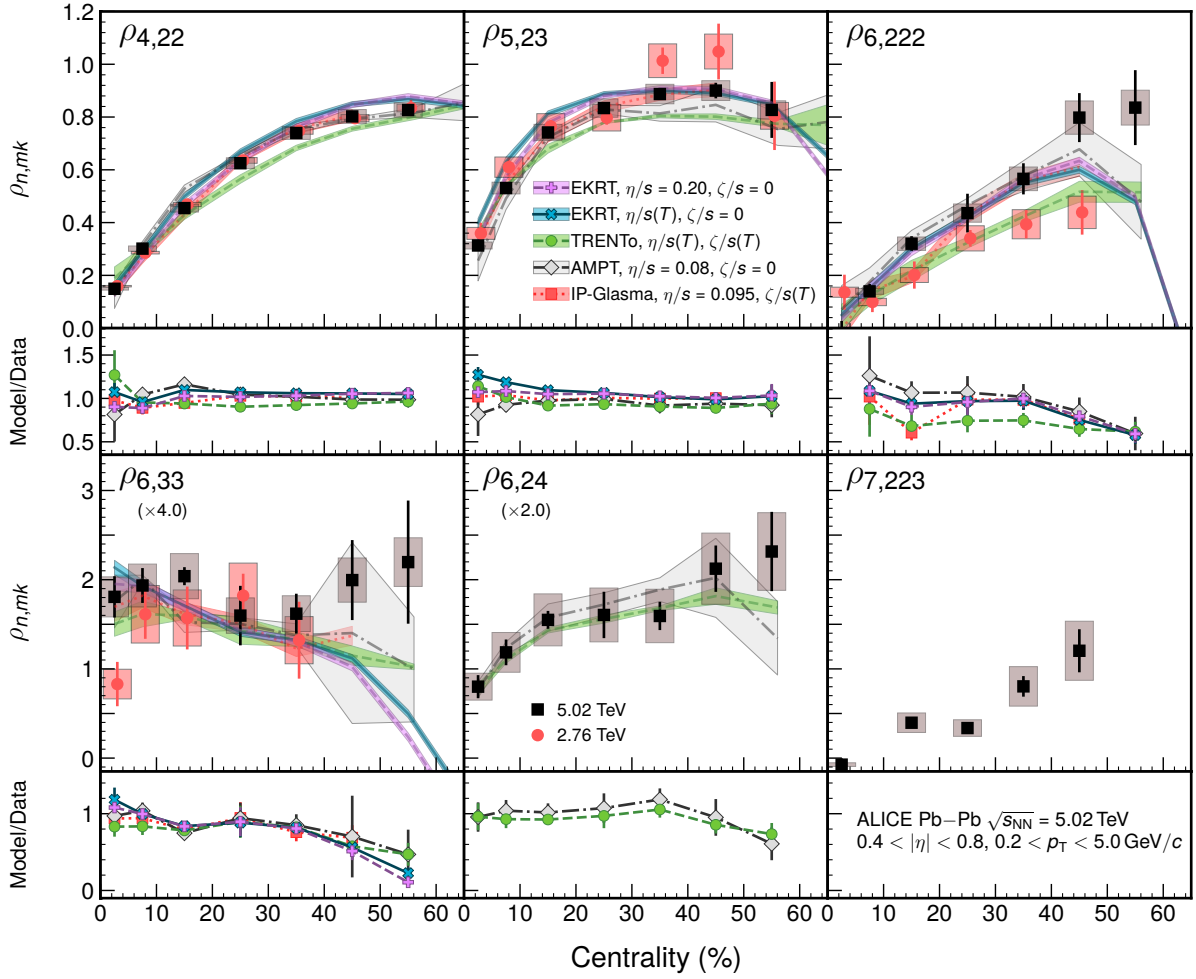


Fig. 5: Symmetry-plane correlations as a function of centrality in Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV (black markers) compared with those in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV [53], along with five different hydrodynamic calculations shown as color bands. On the bottom part of each panel, the ratios between model calculations and the data are shown. For some panels, the data points are scaled by an indicated factor for better visibility.

models. Finally, the performance of AMPT+param2 can be considered good within the reported χ^2/N_{dof} values. It is noted that the magnitude of v_7 is best depicted by AMPT+param2 amongst the three models used.

The performance of the models with respect to the symmetry-plane correlations is quantified in panel (b) of Fig. 7. IP-Glasma+param4 has by far the best estimates of $\rho_{n,mk}$ for $\rho_{4,22}$ and $\rho_{5,23}$. For other models, the model depiction is comparable. In low harmonic orders, EKRT+param0 shows good agreement with the data, as well as AMPT+param2, which has the best agreement in higher harmonic orders. For TRenTo+param3, the agreement is slightly worse for $\rho_{5,23}$ and $\rho_{6,222}$.

The panel (c) of Fig. 7 shows the χ^2/N_{dof} for non-linear flow mode coefficients. As seen also in Fig. 6, TRenTo+param3 consistently provides the most successful overall description of the data. For other models the data are more frequently over- or underestimated. TRenTo+param3 estimates $\chi_{n,mk}$ better than it does the v_n coefficients, for which significant overestimation was present at almost every harmonic order (see Fig. 2). For EKRT+param0 the agreement is good, but the calculation over- or underestimates in some cases especially in most central or mid-peripheral collisions. Most of the observables are better described by the calculations using EKRT+param0 with a const $\eta/s = 0.2$ as compared to results

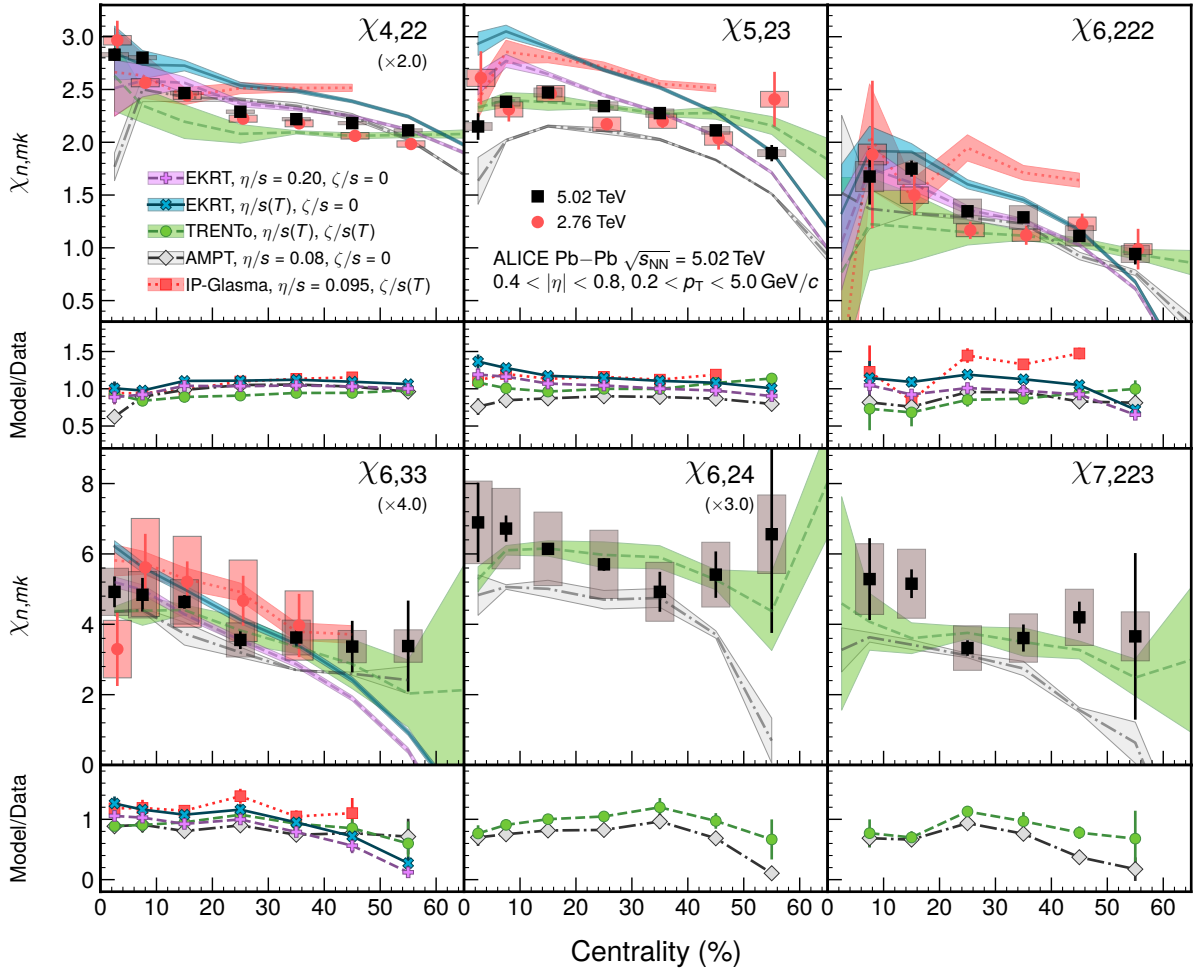


Fig. 6: Non-linear mode coefficients as a function of centrality in Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV (black markers) compared with those from $\sqrt{s_{NN}} = 2.76$ TeV (red markers) [53], along with five different hydrodynamic calculations shown as color bands. On the bottom part of each panel, the ratios between model calculations and the data are shown. For some panels, the points are scaled by an indicated factor for better visibility across the panels.

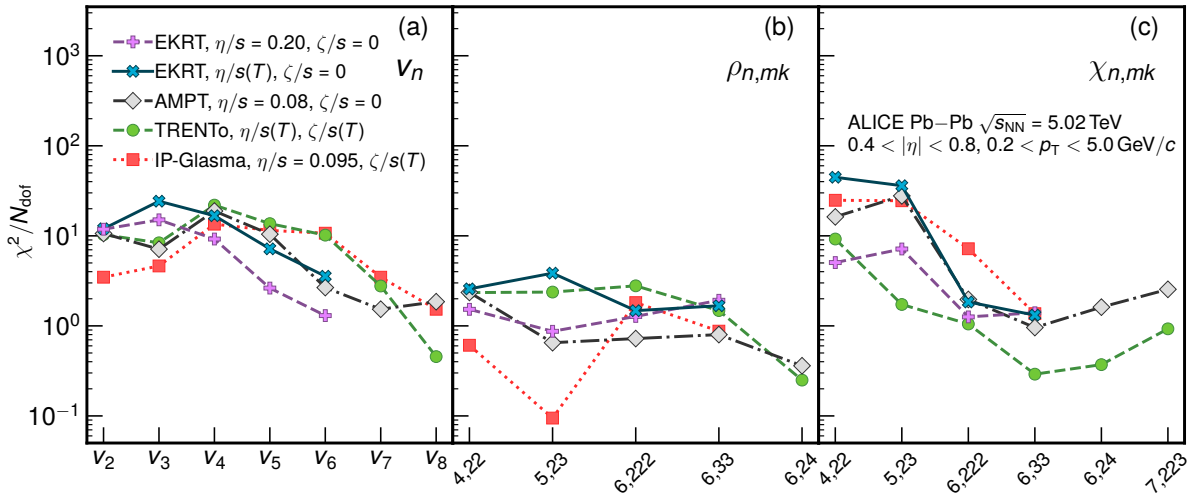


Fig. 7: Overview of various model comparisons with data, quantified by χ^2/N_{dof} . Lower χ^2/N_{dof} represents a better overall description for a given observable.

from EKRT+param1 which uses a temperature dependent η/s value. AMPT+param2 performs worse for low-order harmonics as it overpredicts the data in central and mid-central collisions. Of the five configurations, IP-Glasma+param4 describes the data worse in all harmonic orders.

The deviation of the calculated results from the measured value of each observable is of the same order of magnitude for the different models. Even where the model results show gross agreement with overall features in data, the values of χ^2/N_{dof} vary considerably from one harmonic order to another. Considering the χ^2/N_{dof} to be a goodness-of-fit estimate to validate any model, these variations suggest that the sensitivity of the different observables on the initial conditions, η/s , and ζ/s are reflected differently in the model calculations. Since the current uncertainties in the model calculations are large for higher order harmonics, the absolute χ^2 test should not be over-interpreted.

6 Summary

The measurements of anisotropic flow coefficients (v_n), non-linear flow mode coefficients ($\chi_{n,mk}$), and correlations among different symmetry planes ($\rho_{n,mk}$) in Pb–Pb collisions at $\sqrt{s_{NN}}=5.02$ TeV are presented. The anisotropic flow coefficients are measured up to v_9 , where v_8 and v_9 are measured for the first time at LHC energies. It is observed that v_n decreases as n increases, observing n -ordered damping up to $n=7$. The v_n is found to be enhanced for $n > 7$. The non-linear contribution becomes dominant towards peripheral collisions in all harmonic orders. The strength of the non-linear flow mode and the symmetry plane correlations depends also on harmonic orders. The non-linear flow mode coefficients show a clear centrality and harmonic order dependencies and the strongest non-linear mode coefficients is observed for the fifth and seventh harmonic orders.

These results are compared with various hydrodynamic model calculations with different initial conditions, as well as different parameterizations of η/s and ζ/s . None of the models presented in this paper simultaneously describe the v_n coefficients, $\chi_{n,mk}$, or $\rho_{n,mk}$. Based on the model and data comparisons, among all the models, the event-by-event viscous hydrodynamic model with EKRT initial conditions and a constant $\eta/s=0.2$ is observed to describe the data best, as far as the harmonics up to the sixth order are concerned. As a result further tuning is required to find the accurate parameterization of η/s and ζ/s . It is found that the different order harmonic observables presented in this paper have different sensitivities to the initial conditions and the system properties. These results allow further model parameters to be optimized and the initial conditions and the transport properties of nuclear matter in ultra-relativistic heavy-ion collisions to be better constrained.

Acknowledgements

The ALICE Collaboration would like to thank all its engineers and technicians for their invaluable contributions to the construction of the experiment and the CERN accelerator teams for the outstanding performance of the LHC complex. The ALICE Collaboration gratefully acknowledges the resources and support provided by all Grid centres and the Worldwide LHC Computing Grid (WLCG) collaboration. The ALICE Collaboration acknowledges the following funding agencies for their support in building and running the ALICE detector: A. I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation (ANSL), State Committee of Science and World Federation of Scientists (WFS), Armenia; Austrian Academy of Sciences, Austrian Science Fund (FWF): [M 2467-N36] and Nationalstiftung für Forschung, Technologie und Entwicklung, Austria; Ministry of Communications and High Technologies, National Nuclear Research Center, Azerbaijan; Conselho Nacional de Desenvolvimento Científico e Tecnológico (CNPq), Financiadora de Estudos e Projetos (Finep), Fundação de Amparo à Pesquisa do Estado de São Paulo (FAPESP) and Universidade Federal do Rio Grande do Sul (UFRGS), Brazil; Ministry of Education of China (MOEC), Ministry of Science & Technology of China (MSTC) and National Natural Science Foundation of China (NSFC), China; Ministry of Science and Education

and Croatian Science Foundation, Croatia; Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Cubaenergía, Cuba; Ministry of Education, Youth and Sports of the Czech Republic, Czech Republic; The Danish Council for Independent Research | Natural Sciences, the VILLUM FONDEN and Danish National Research Foundation (DNRF), Denmark; Helsinki Institute of Physics (HIP), Finland; Commissariat à l’Energie Atomique (CEA), Institut National de Physique Nucléaire et de Physique des Particules (IN2P3) and Centre National de la Recherche Scientifique (CNRS) and Région des Pays de la Loire, France; Bundesministerium für Bildung und Forschung (BMBF) and GSI Helmholtzzentrum für Schwerionenforschung GmbH, Germany; General Secretariat for Research and Technology, Ministry of Education, Research and Religions, Greece; National Research, Development and Innovation Office, Hungary; Department of Atomic Energy Government of India (DAE), Department of Science and Technology, Government of India (DST), University Grants Commission, Government of India (UGC) and Council of Scientific and Industrial Research (CSIR), India; Indonesian Institute of Science, Indonesia; Centro Fermi - Museo Storico della Fisica e Centro Studi e Ricerche Enrico Fermi and Istituto Nazionale di Fisica Nucleare (INFN), Italy; Institute for Innovative Science and Technology , Nagasaki Institute of Applied Science (IIST), Japanese Ministry of Education, Culture, Sports, Science and Technology (MEXT) and Japan Society for the Promotion of Science (JSPS) KAKENHI, Japan; Consejo Nacional de Ciencia (CONACYT) y Tecnología, through Fondo de Cooperación Internacional en Ciencia y Tecnología (FONCICYT) and Dirección General de Asuntos del Personal Académico (DGAPA), Mexico; Nederlandse Organisatie voor Wetenschappelijk Onderzoek (NWO), Netherlands; The Research Council of Norway, Norway; Commission on Science and Technology for Sustainable Development in the South (COMSATS), Pakistan; Pontificia Universidad Católica del Perú, Peru; Ministry of Science and Higher Education and National Science Centre, Poland; Korea Institute of Science and Technology Information and National Research Foundation of Korea (NRF), Republic of Korea; Ministry of Education and Scientific Research, Institute of Atomic Physics and Ministry of Research and Innovation and Institute of Atomic Physics, Romania; Joint Institute for Nuclear Research (JINR), Ministry of Education and Science of the Russian Federation, National Research Centre Kurchatov Institute, Russian Science Foundation and Russian Foundation for Basic Research, Russia; Ministry of Education, Science, Research and Sport of the Slovak Republic, Slovakia; National Research Foundation of South Africa, South Africa; Swedish Research Council (VR) and Knut & Alice Wallenberg Foundation (KAW), Sweden; European Organization for Nuclear Research, Switzerland; Suranaree University of Technology (SUT), National Science and Technology Development Agency (NSDTA) and Office of the Higher Education Commission under NRU project of Thailand, Thailand; Turkish Atomic Energy Agency (TAEK), Turkey; National Academy of Sciences of Ukraine, Ukraine; Science and Technology Facilities Council (STFC), United Kingdom; National Science Foundation of the United States of America (NSF) and United States Department of Energy, Office of Nuclear Physics (DOE NP), United States of America.

References

- [1] J.-Y. Ollitrault, “Anisotropy as a signature of transverse collective flow,” *Phys. Rev.* **D46** (1992) 229–245.
- [2] S. A. Voloshin, A. M. Poskanzer, and R. Snellings, “Collective phenomena in non-central nuclear collisions,” *Landolt-Bornstein* **23** (2010) 293–333, arXiv:0809.2949 [nucl-ex].
- [3] S. Voloshin and Y. Zhang, “Flow study in relativistic nuclear collisions by Fourier expansion of Azimuthal particle distributions,” *Z. Phys.* **C70** (1996) 665–672, arXiv:hep-ph/9407282 [hep-ph].
- [4] A. M. Poskanzer and S. A. Voloshin, “Methods for analyzing anisotropic flow in relativistic nuclear collisions,” *Phys. Rev.* **C58** (1998) 1671–1678, arXiv:nucl-ex/9805001 [nucl-ex].

- [5] ALICE Collaboration, K. Aamodt *et al.*, “Elliptic flow of charged particles in Pb–Pb collisions at 2.76 TeV,” *Phys. Rev. Lett.* **105** (2010) 252302, arXiv:1011.3914 [nucl-ex].
- [6] ALICE Collaboration, K. Aamodt *et al.*, “Higher harmonic anisotropic flow measurements of charged particles in Pb–Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV,” *Phys. Rev. Lett.* **107** (2011) 032301, arXiv:1105.3865 [nucl-ex].
- [7] ALICE Collaboration, B. Abelev *et al.*, “Anisotropic flow of charged hadrons, pions and (anti-)protons measured at high transverse momentum in Pb–Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV,” *Phys. Lett.* **B719** (2013) 18–28, arXiv:1205.5761 [nucl-ex].
- [8] ALICE Collaboration, B. Abelev *et al.*, “Elliptic flow of identified hadrons in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV,” *JHEP* **06** (2015) 190, arXiv:1405.4632 [nucl-ex].
- [9] ALICE Collaboration, J. Adam *et al.*, “Higher harmonic flow coefficients of identified hadrons in Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV,” *JHEP* **09** (2016) 164, arXiv:1606.06057 [nucl-ex].
- [10] ALICE Collaboration, J. Adam *et al.*, “Anisotropic flow of charged particles in Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV,” *Phys. Rev. Lett.* **116** no. 13, (2016) 132302, arXiv:1602.01119 [nucl-ex].
- [11] ALICE Collaboration, S. Acharya *et al.*, “Searches for transverse momentum dependent flow vector fluctuations in Pb–Pb and p–Pb collisions at the LHC,” *JHEP* **09** (2017) 032, arXiv:1707.05690 [nucl-ex].
- [12] ALICE Collaboration, S. Acharya *et al.*, “Anisotropic flow of identified particles in Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV,” *JHEP* **09** (2018) 006, arXiv:1805.04390 [nucl-ex].
- [13] D. A. Teaney, “Viscous Hydrodynamics and the Quark Gluon Plasma,” arXiv:0905.2433 [nucl-th].
- [14] D. Molnar and M. Gyulassy, “Saturation of elliptic flow and the transport opacity of the gluon plasma at RHIC,” *Nucl. Phys.* **A697** (2002) 495–520, arXiv:nucl-th/0104073 [nucl-th]. [Erratum: *Nucl. Phys.*A703,893(2002)].
- [15] D. Teaney, “The Effects of viscosity on spectra, elliptic flow, and HBT radii,” *Phys. Rev.* **C68** (2003) 034913, arXiv:nucl-th/0301099 [nucl-th].
- [16] R. A. Lacey, N. N. Ajitanand, J. M. Alexander, P. Chung, W. G. Holzmann, M. Issah, A. Taranenko, P. Danielewicz, and H. Stoecker, “Has the QCD Critical Point been Signaled by Observations at RHIC?,” *Phys. Rev. Lett.* **98** (2007) 092301, arXiv:nucl-ex/0609025 [nucl-ex].
- [17] H.-J. Drescher, A. Dumitru, C. Gombeaud, and J.-Y. Ollitrault, “The Centrality dependence of elliptic flow, the hydrodynamic limit, and the viscosity of hot QCD,” *Phys. Rev.* **C76** (2007) 024905, arXiv:0704.3553 [nucl-th].
- [18] Z. Xu, C. Greiner, and H. Stoecker, “PQCD calculations of elliptic flow and shear viscosity at RHIC,” *Phys. Rev. Lett.* **101** (2008) 082302, arXiv:0711.0961 [nucl-th].
- [19] D. Molnar and P. Huovinen, “Dissipative effects from transport and viscous hydrodynamics,” *J. Phys.* **G35** (2008) 104125, arXiv:0806.1367 [nucl-th].
- [20] U. Heinz and R. Snellings, “Collective flow and viscosity in relativistic heavy-ion collisions,” *Ann. Rev. Nucl. Part. Sci.* **63** (2013) 123–151, arXiv:1301.2826 [nucl-th].

- [21] H. Song, Y. Zhou, and K. Gajdosova, “Collective flow and hydrodynamics in large and small systems at the LHC,” *Nucl. Sci. Tech.* **28** no. 7, (2017) 99, arXiv:1703.00670 [nucl-th].
- [22] P. Kovtun, D. T. Son, and A. O. Starinets, “Viscosity in strongly interacting quantum field theories from black hole physics,” *Phys. Rev. Lett.* **94** (2005) 111601, arXiv:hep-th/0405231 [hep-th].
- [23] B. Alver and G. Roland, “Collision geometry fluctuations and triangular flow in heavy-ion collisions,” *Phys. Rev.* **C81** (2010) 054905, arXiv:1003.0194 [nucl-th]. [Erratum: *Phys. Rev.* **C82**,039903(2010)].
- [24] D. Teaney and L. Yan, “Triangularity and dipole asymmetry in heavy ion collisions,” *Phys. Rev.* **C83** (2011) 064904, arXiv:1010.1876 [nucl-th].
- [25] H. Niemi, K. J. Eskola, and R. Paatelainen, “Event-by-event fluctuations in a perturbative QCD + saturation + hydrodynamics model: Determining QCD matter shear viscosity in ultrarelativistic heavy-ion collisions,” *Phys. Rev.* **C93** no. 2, (2016) 024907, arXiv:1505.02677 [hep-ph].
- [26] H. Niemi, G. S. Denicol, H. Holopainen, and P. Huovinen, “Event-by-event distributions of azimuthal asymmetries in ultrarelativistic heavy-ion collisions,” *Phys. Rev.* **C87** no. 5, (2013) 054901, arXiv:1212.1008 [nucl-th].
- [27] **STAR** Collaboration, C. Adler *et al.*, “Identified particle elliptic flow in Au + Au collisions at $\sqrt{s_{NN}} = 130$ GeV,” *Phys. Rev. Lett.* **87** (2001) 182301, arXiv:nucl-ex/0107003 [nucl-ex].
- [28] **STAR** Collaboration, C. Adler *et al.*, “Elliptic flow from two and four particle correlations in Au+Au collisions at $\sqrt{s_{NN}} = 130$ GeV,” *Phys. Rev.* **C66** (2002) 034904, arXiv:nucl-ex/0206001 [nucl-ex].
- [29] **PHENIX** Collaboration, S. S. Adler *et al.*, “Elliptic flow of identified hadrons in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV,” *Phys. Rev. Lett.* **91** (2003) 182301, arXiv:nucl-ex/0305013 [nucl-ex].
- [30] **STAR** Collaboration, J. Adams *et al.*, “Azimuthal anisotropy at RHIC: The First and fourth harmonics,” *Phys. Rev. Lett.* **92** (2004) 062301, arXiv:nucl-ex/0310029 [nucl-ex].
- [31] **STAR** Collaboration, L. Adamczyk *et al.*, “Azimuthal anisotropy in U+U and Au+Au collisions at RHIC,” *Phys. Rev. Lett.* **115** no. 22, (2015) 222301, arXiv:1505.07812 [nucl-ex].
- [32] **PHOBOS** Collaboration, B. Alver *et al.*, “Event-by-Event Fluctuations of Azimuthal Particle Anisotropy in Au + Au Collisions at $\sqrt{s_{NN}} = 200$ GeV,” *Phys. Rev. Lett.* **104** (2010) 142301, arXiv:nucl-ex/0702036 [nucl-ex].
- [33] T. Hirano, U. W. Heinz, D. Kharzeev, R. Lacey, and Y. Nara, “Hadronic dissipative effects on elliptic flow in ultrarelativistic heavy-ion collisions,” *Phys. Lett.* **B636** (2006) 299–304, arXiv:nucl-th/0511046 [nucl-th].
- [34] P. Romatschke and U. Romatschke, “Viscosity Information from Relativistic Nuclear Collisions: How Perfect is the Fluid Observed at RHIC?,” *Phys. Rev. Lett.* **99** (2007) 172301, arXiv:0706.1522 [nucl-th].
- [35] A. K. Chaudhuri, “Saturation of elliptic flow and shear viscosity,” arXiv:0708.1252 [nucl-th].

- [36] H. Song, S. A. Bass, U. Heinz, T. Hirano, and C. Shen, “200 A GeV Au+Au collisions serve a nearly perfect quark-gluon liquid,” *Phys. Rev. Lett.* **106** (2011) 192301, arXiv:1011.2783 [nucl-th]. [Erratum: *Phys. Rev. Lett.* 109,139904(2012)].
- [37] M. Luzum and J.-Y. Ollitrault, “Extracting the shear viscosity of the quark-gluon plasma from flow in ultra-central heavy-ion collisions,” *Nucl. Phys.* **A904-905** (2013) 377c–380c, arXiv:1210.6010 [nucl-th].
- [38] C. Shen, S. A. Bass, T. Hirano, P. Huovinen, Z. Qiu, H. Song, and U. Heinz, “The QGP shear viscosity: Elusive goal or just around the corner?,” *J. Phys.* **G38** (2011) 124045, arXiv:1106.6350 [nucl-th].
- [39] ALICE Collaboration, J. Adam *et al.*, “Correlated event-by-event fluctuations of flow harmonics in Pb-Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV,” *Phys. Rev. Lett.* **117** (2016) 182301, arXiv:1604.07663 [nucl-ex].
- [40] ALICE Collaboration, S. Acharya *et al.*, “Systematic studies of correlations between different order flow harmonics in Pb-Pb collisions at $\sqrt{s_{\text{NN}}} = 2.76$ TeV,” *Phys. Rev.* **C97** no. 2, (2018) 024906, arXiv:1709.01127 [nucl-ex].
- [41] P. Bozek, “Flow and interferometry in 3+1 dimensional viscous hydrodynamics,” *Phys. Rev.* **C85** (2012) 034901, arXiv:1110.6742 [nucl-th].
- [42] J.-B. Rose, J.-F. Paquet, G. S. Denicol, M. Luzum, B. Schenke, S. Jeon, and C. Gale, “Extracting the bulk viscosity of the quark gluon plasma,” *Nucl. Phys.* **A931** (2014) 926–930, arXiv:1408.0024 [nucl-th].
- [43] S. Ryu, J. F. Paquet, C. Shen, G. S. Denicol, B. Schenke, S. Jeon, and C. Gale, “Importance of the Bulk Viscosity of QCD in Ultrarelativistic Heavy-Ion Collisions,” *Phys. Rev. Lett.* **115** no. 13, (2015) 132301, arXiv:1502.01675 [nucl-th].
- [44] S. Ryu, J.-F. Paquet, C. Shen, G. Denicol, B. Schenke, S. Jeon, and C. Gale, “Effects of bulk viscosity and hadronic rescattering in heavy ion collisions at energies available at the BNL Relativistic Heavy Ion Collider and at the CERN Large Hadron Collider,” *Phys. Rev.* **C97** no. 3, (2018) 034910, arXiv:1704.04216 [nucl-th].
- [45] J. E. Bernhard, J. S. Moreland, S. A. Bass, J. Liu, and U. Heinz, “Applying Bayesian parameter estimation to relativistic heavy-ion collisions: simultaneous characterization of the initial state and quark-gluon plasma medium,” *Phys. Rev.* **C94** no. 2, (2016) 024907, arXiv:1605.03954 [nucl-th].
- [46] A. Dubla, S. Masciocchi, J. M. Pawlowski, B. Schenke, C. Shen, and J. Stachel, “Towards QCD-assisted hydrodynamics for heavy-ion collision phenomenology,” *Nucl. Phys.* **A979** (2018) 251–264, arXiv:1805.02985 [nucl-th].
- [47] J. E. Bernhard, J. S. Moreland, and S. A. Bass, “Bayesian estimation of the specific shear and bulk viscosity of quark-gluon plasma,” *Nature Physics* (2019) .
<https://doi.org/10.1038/s41567-019-0611-8>.
- [48] B. H. Alver, C. Gombeaud, M. Luzum, and J.-Y. Ollitrault, “Triangular flow in hydrodynamics and transport theory,” *Phys. Rev.* **C82** (2010) 034913, arXiv:1007.5469 [nucl-th].
- [49] ATLAS Collaboration, M. Aaboud *et al.*, “Measurement of the azimuthal anisotropy of charged particles produced in $\sqrt{s_{\text{NN}}} = 5.02$ TeV Pb+Pb collisions with the ATLAS detector,” *Eur. Phys. J.* **C78** no. 12, (2018) 997, arXiv:1808.03951 [nucl-ex].

- [50] F. G. Gardim, F. Grassi, M. Luzum, and J.-Y. Ollitrault, “Mapping the hydrodynamic response to the initial geometry in heavy-ion collisions,” *Phys. Rev.* **C85** (2012) 024908, arXiv:1111.6538 [nucl-th].
- [51] F. G. Gardim, J. Noronha-Hostler, M. Luzum, and F. Grassi, “Effects of viscosity on the mapping of initial to final state in heavy ion collisions,” *Phys. Rev.* **C91** no. 3, (2015) 034902, arXiv:1411.2574 [nucl-th].
- [52] D. Teaney and L. Yan, “Event-plane correlations and hydrodynamic simulations of heavy ion collisions,” *Phys. Rev.* **C90** no. 2, (2014) 024902, arXiv:1312.3689 [nucl-th].
- [53] ALICE Collaboration, S. Acharya *et al.*, “Linear and non-linear flow modes in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV,” *Phys. Lett.* **B773** (2017) 68–80, arXiv:1705.04377 [nucl-ex].
- [54] M. L. Miller, K. Reygers, S. J. Sanders, and P. Steinberg, “Glauber modeling in high energy nuclear collisions,” *Ann. Rev. Nucl. Part. Sci.* **57** (2007) 205–243, arXiv:nucl-ex/0701025 [nucl-ex].
- [55] Z. Qiu and U. W. Heinz, “Event-by-event shape and flow fluctuations of relativistic heavy-ion collision fireballs,” *Phys. Rev.* **C84** (2011) 024911, arXiv:1104.0650 [nucl-th].
- [56] H.-J. Drescher and Y. Nara, “Eccentricity fluctuations from the color glass condensate at RHIC and LHC,” *Phys. Rev.* **C76** (2007) 041903, arXiv:0707.0249 [nucl-th].
- [57] S. McDonald, C. Shen, F. Fillion-Gourdeau, S. Jeon, and C. Gale, “Hydrodynamic predictions for Pb+Pb collisions at 5.02 TeV,” *Phys. Rev.* **C95** no. 6, (2017) 064913, arXiv:1609.02958 [hep-ph].
- [58] ALICE Collaboration, S. Acharya *et al.*, “Energy dependence and fluctuations of anisotropic flow in Pb-Pb collisions at $\sqrt{s_{NN}} = 5.02$ and 2.76 TeV,” *JHEP* **07** (2018) 103, arXiv:1804.02944 [nucl-ex].
- [59] L. Yan and J.-Y. Ollitrault, “ v_4, v_5, v_6, v_7 : nonlinear hydrodynamic response versus LHC data,” *Phys. Lett.* **B744** (2015) 82–87, arXiv:1502.02502 [nucl-th].
- [60] D. Teaney and L. Yan, “Non linearities in the harmonic spectrum of heavy ion collisions with ideal and viscous hydrodynamics,” *Phys. Rev.* **C86** (2012) 044908, arXiv:1206.1905 [nucl-th].
- [61] J. Jia and S. Mohapatra, “A Method for studying initial geometry fluctuations via event plane correlations in heavy ion collisions,” *Eur. Phys. J.* **C73** (2013) 2510, arXiv:1203.5095 [nucl-th].
- [62] M. Luzum, “Flow fluctuations and long-range correlations: elliptic flow and beyond,” *J. Phys.* **G38** (2011) 124026, arXiv:1107.0592 [nucl-th].
- [63] Z. Qiu and U. Heinz, “Hydrodynamic event-plane correlations in Pb+Pb collisions at $\sqrt{s} = 2.76$ ATeV,” *Phys. Lett.* **B717** (2012) 261–265, arXiv:1208.1200 [nucl-th].
- [64] ATLAS Collaboration, G. Aad *et al.*, “Measurement of event-plane correlations in $\sqrt{s_{NN}} = 2.76$ TeV lead-lead collisions with the ATLAS detector,” *Phys. Rev.* **C90** no. 2, (2014) 024905, arXiv:1403.0489 [hep-ex].
- [65] CMS Collaboration, S. Chatrchyan *et al.*, “Measurement of higher-order harmonic azimuthal anisotropy in PbPb collisions at $\sqrt{s_{NN}} = 2.76$ TeV,” *Phys. Rev.* **C89** no. 4, (2014) 044906, arXiv:1310.8651 [nucl-ex].

- [66] M. Luzum and J.-Y. Ollitrault, “Eliminating experimental bias in anisotropic-flow measurements of high-energy nuclear collisions,” *Phys. Rev.* **C87** no. 4, (2013) 044907, arXiv:1209.2323 [nucl-ex].
- [67] ALICE Collaboration, K. Aamodt *et al.*, “The ALICE experiment at the CERN LHC,” *JINST* **3** (2008) S08002.
- [68] ALICE Collaboration, B. Abelev *et al.*, “Performance of the ALICE Experiment at the CERN LHC,” *Int. J. Mod. Phys.* **A29** (2014) 1430044, arXiv:1402.4476 [nucl-ex].
- [69] ALICE Collaboration, P. Cortese *et al.*, “ALICE: Physics performance report, volume I,” *J. Phys.* **G30** (2004) 1517–1763.
- [70] ALICE Collaboration, P. Cortese *et al.*, “ALICE: Physics performance report, volume II,” *J. Phys.* **G32** (2006) 1295–2040.
- [71] ALICE Collaboration, E. Abbas *et al.*, “Performance of the ALICE VZERO system,” *JINST* **8** (2013) P10016, arXiv:1306.3130 [nucl-ex].
- [72] ALICE Collaboration, K. Aamodt *et al.*, “Alignment of the ALICE Inner Tracking System with cosmic-ray tracks,” *JINST* **5** (2010) P03003, arXiv:1001.0502 [physics.ins-det].
- [73] ALICE Collaboration, F. Carnesecchi, “Performance of the ALICE Time-Of-Flight detector at the LHC,” *JINST* **14** no. 06, (2019) C06023, arXiv:1806.03825 [physics.ins-det].
- [74] ALICE Collaboration, B. Abelev *et al.*, “Centrality determination of Pb-Pb collisions at $\sqrt{s_{NN}}=2.76$ TeV with ALICE,” *Phys. Rev.* **C88** no. 4, (2013) 044909, arXiv:1301.4361 [nucl-ex].
- [75] J. Alme *et al.*, “The ALICE TPC, a large 3-dimensional tracking device with fast readout for ultra-high multiplicity events,” *Nucl. Instrum. Meth.* **A622** (2010) 316–367, arXiv:1001.1950 [physics.ins-det].
- [76] ALICE Collaboration, “The ALICE definition of primary particles,” <https://cds.cern.ch/record/2270008>.
- [77] X.-N. Wang and M. Gyulassy, “HIJING: A Monte Carlo model for multiple jet production in p p, p A and A A collisions,” *Phys. Rev.* **D44** (1991) 3501–3516.
- [78] M. Gyulassy and X.-N. Wang, “HIJING 1.0: A Monte Carlo program for parton and particle production in high-energy hadronic and nuclear collisions,” *Comput. Phys. Commun.* **83** (1994) 307, arXiv:nucl-th/9502021 [nucl-th].
- [79] R. Brun, F. Bruyant, F. Carminati, S. Giani, M. Maire, A. McPherson, G. Patrick, and L. Urban, “GEANT Detector Description and Simulation Tool,”.
- [80] A. Bilandzic, C. H. Christensen, K. Gulbrandsen, A. Hansen, and Y. Zhou, “Generic framework for anisotropic flow analyses with multiparticle azimuthal correlations,” *Phys. Rev.* **C89** no. 6, (2014) 064904, arXiv:1312.3572 [nucl-ex].
- [81] H. Niemi, K. J. Eskola, R. Paatelainen, and K. Tuominen, “Predictions for 5.023 TeV Pb + Pb collisions at the CERN Large Hadron Collider,” *Phys. Rev.* **C93** no. 1, (2016) 014912, arXiv:1511.04296 [hep-ph].
- [82] W. Zhao, H.-j. Xu, and H. Song, “Collective flow in 2.76 A TeV and 5.02 A TeV Pb+Pb collisions,” *Eur. Phys. J.* **C77** no. 9, (2017) 645, arXiv:1703.10792 [nucl-th].

- [83] C. Shen, Z. Qiu, H. Song, J. Bernhard, S. Bass, and U. Heinz, “The iEBE-VISHNU code package for relativistic heavy-ion collisions,” *Comput. Phys. Commun.* **199** (2016) 61–85, arXiv:1409.8164 [nucl-th]. <https://u.osu.edu/vishnu/>.
- [84] R. S. Bhalerao, A. Jaiswal, and S. Pal, “Collective flow in event-by-event partonic transport plus hydrodynamics hybrid approach,” *Phys. Rev.* **C92** no. 1, (2015) 014903, arXiv:1503.03862 [nucl-th].
- [85] L. Pang, Q. Wang, and X.-N. Wang, “Effects of initial flow velocity fluctuation in event-by-event (3+1)D hydrodynamics,” *Phys. Rev.* **C86** (2012) 024911, arXiv:1205.5019 [nucl-th].
- [86] H.-j. Xu, Z. Li, and H. Song, “High-order flow harmonics of identified hadrons in 2.76 TeV Pb + Pb collisions,” *Phys. Rev.* **C93** no. 6, (2016) 064905, arXiv:1602.02029 [nucl-th].
- [87] J. S. Moreland, J. E. Bernhard, and S. A. Bass, “Alternative ansatz to wounded nucleon and binary collision scaling in high-energy nuclear collisions,” *Phys. Rev.* **C92** no. 1, (2015) 011901, arXiv:1412.4708 [nucl-th]. <http://qcd.phy.duke.edu/trento/>.
- [88] B. Schenke, S. Jeon, and C. Gale, “Elliptic and triangular flow in event-by-event (3+1)D viscous hydrodynamics,” *Phys. Rev. Lett.* **106** (2011) 042301, arXiv:1009.3244 [hep-ph]. <http://www.physics.mcgill.ca/music/>.
- [89] B. Schenke, P. Tribedy, and R. Venugopalan, “Fluctuating Glasma initial conditions and flow in heavy ion collisions,” *Phys. Rev. Lett.* **108** (2012) 252301, arXiv:1202.6646 [nucl-th].
- [90] Z. Qiu, C. Shen, and U. Heinz, “Hydrodynamic elliptic and triangular flow in Pb-Pb collisions at $\sqrt{s} = 2.76$ TeV,” *Phys. Lett.* **B707** (2012) 151–155, arXiv:1110.3033 [nucl-th].
- [91] C. Shen, U. Heinz, P. Huovinen, and H. Song, “Systematic parameter study of hadron spectra and elliptic flow from viscous hydrodynamic simulations of Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV,” *Phys. Rev.* **C82** (2010) 054904, arXiv:1010.1856 [nucl-th].
- [92] C. Shen, U. Heinz, P. Huovinen, and H. Song, “Radial and elliptic flow in Pb+Pb collisions at the Large Hadron Collider from viscous hydrodynamic,” *Phys. Rev.* **C84** (2011) 044903, arXiv:1105.3226 [nucl-th].
- [93] G. S. Denicol, T. Kodama, T. Koide, and P. Mota, “Effect of bulk viscosity on Elliptic Flow near QCD phase transition,” *Phys. Rev.* **C80** (2009) 064901, arXiv:0903.3595 [hep-ph].
- [94] F. Karsch, D. Kharzeev, and K. Tuchin, “Universal properties of bulk viscosity near the QCD phase transition,” *Phys. Lett.* **B663** (2008) 217–221, arXiv:0711.0914 [hep-ph].
- [95] J. Noronha-Hostler, J. Noronha, and C. Greiner, “Transport Coefficients of Hadronic Matter near $T(c)$,” *Phys. Rev. Lett.* **103** (2009) 172302, arXiv:0811.1571 [nucl-th].
- [96] E. Molnar, H. Niemi, and D. H. Rischke, “Numerical tests of causal relativistic dissipative fluid dynamics,” *Eur. Phys. J.* **C65** (2010) 615–635, arXiv:0907.2583 [nucl-th].
- [97] E. Shuryak, “The sounds of the Little and Big Bangs,” *Universe* **3** no. 4, (2017) 75, arXiv:1710.03776 [hep-ph].
- [98] P. Staig and E. Shuryak, “The Fate of the Initial State Fluctuations in Heavy Ion Collisions. III The Second Act of Hydrodynamics,” *Phys. Rev.* **C84** (2011) 044912, arXiv:1105.0676 [nucl-th].

- [99] R. A. Lacey, Y. Gu, X. Gong, D. Reynolds, N. N. Ajitanand, J. M. Alexander, A. Mwai, and A. Taranenko, “Is anisotropic flow really acoustic?,” [arXiv:1301.0165](#) [nucl-ex].
- [100] J. Qian, U. W. Heinz, and J. Liu, “Mode-coupling effects in anisotropic flow in heavy-ion collisions,” *Phys. Rev.* **C93** no. 6, (2016) 064901, [arXiv:1602.02813](#) [nucl-th].

A List of Observables

In this section the complete list of the measured observables is presented. By root-mean-squaring the equations in Eq. (3), one obtains a starting point for the definitions presented in this section. Provided that the linear and non-linear parts are uncorrelated, the following harmonic projections are obtained

$$\begin{aligned}
 v_{4,22} &= \frac{\Re\langle V_4(V_2^*)^2 \rangle}{\sqrt{\langle |V_2|^4 \rangle}} & v_{5,23} &= \frac{\Re\langle V_5 V_2^* V_3^* \rangle}{\sqrt{\langle |V_2|^2 |V_3|^2 \rangle}} \\
 &\approx \langle v_4 \cos(4\psi_4 - 4\psi_2) \rangle, & &\approx \langle v_5 \cos(5\psi_5 - 3\psi_3 - 2\psi_2) \rangle, \\
 v_{6,222} &= \frac{\Re\langle V_6(V_2^*)^3 \rangle}{\sqrt{\langle |V_2|^6 \rangle}} & v_{6,24} &= \frac{\Re\langle V_6 V_2^* V_4^* \rangle}{\sqrt{\langle |V_2|^2 |V_4|^2 \rangle}} \\
 &\approx \langle v_6 \cos(6\psi_6 - 6\psi_2) \rangle, & &\approx \langle v_6 \cos(6\psi_6 - 4\psi_4 - 2\psi_2) \rangle, \\
 v_{6,33} &= \frac{\Re\langle V_6(V_3^*)^2 \rangle}{\sqrt{\langle |V_3|^4 \rangle}} & v_{7,223} &= \frac{\Re\langle V_7(V_2^*)^2 V_3^* \rangle}{\sqrt{\langle |V_2|^4 |V_3|^2 \rangle}} \\
 &\approx \langle v_6 \cos(6\psi_6 - 6\psi_3) \rangle, & &\approx \langle v_7 \cos(7\psi_7 - 4\psi_2 - 3\psi_3) \rangle, \\
 v_{8,233} &= \frac{\Re\langle V_8 V_2^* (V_3^*)^2 \rangle}{\sqrt{\langle |V_2|^2 |V_3|^4 \rangle}} \\
 &\approx \langle v_8 \cos(8\psi_8 - 2\psi_2 - 6\psi_3) \rangle,
 \end{aligned} \tag{A.1}$$

with $v_{4,22}^2 = \chi_{4,22}^2 \langle |V_2|^4 \rangle$, $v_{5,23}^2 = \chi_{5,23}^2 \langle |V_2|^2 |V_3|^2 \rangle$, \dots . The rest of the observables we define using the harmonic projections in Eq. (A.1). The Symmetry-plane correlations are defined as

$$\begin{aligned}
 \rho_{4,22} &= \frac{v_{4,22}}{v_4}, & \rho_{5,23} &= \frac{v_{5,23}}{v_5}, \\
 \rho_{6,222} &= \frac{v_{6,222}}{v_6}, & \rho_{7,223} &= \frac{v_{7,223}}{v_7}, \\
 \rho_{6,33} &= \frac{v_{6,33}}{v_6},
 \end{aligned} \tag{A.2}$$

and the non-linear mode coefficients

$$\begin{aligned}
 \chi_{4,22} &= \frac{v_{4,22}}{\sqrt{\langle v_2^4 \rangle}}, & \chi_{5,23} &= \frac{v_{5,23}}{\sqrt{\langle v_2^2 v_3^2 \rangle}}, \\
 \chi_{6,222} &= \frac{v_{6,222}}{\sqrt{\langle v_2^6 \rangle}}, & \chi_{7,223} &= \frac{v_{7,223}}{\sqrt{\langle v_2^4 v_3^2 \rangle}}, \\
 \chi_{6,33} &= \frac{v_{6,33}}{\sqrt{\langle v_3^4 \rangle}}, \\
 \chi_{6,24} &= \Re \frac{\langle V_6 V_2^* V_4^* \rangle \langle v_2^4 \rangle - \langle V_6 (V_2^*)^3 \rangle \langle V_4 (V_2^*)^2 \rangle}{(\langle v_2^2 \rangle \langle v_2^4 \rangle - \langle V_4 (V_2^*)^2 \rangle^2) \langle v_2^2 \rangle}.
 \end{aligned} \tag{A.3}$$

The higher order superpositions in Eq. (3) include the coupling constants for the higher order linear responses. In a more complete treatment [100], the extraction of the higher order non-linear flow mode coefficients are performed by correlating the corresponding superpositions with those of the relevant harmonics, effectively resulting in a more general projection. The results agree with the expressions in

Eq. (8), and generate additional high order linear coupling coefficients

$$\begin{aligned}
 \chi_{6,24} &= \Re \frac{\langle V_6 V_2^* V_4^* \rangle \langle v_2^4 \rangle - \langle V_6 (V_2^*)^3 \rangle \langle V_4 (V_2^*)^2 \rangle}{(\langle v_4^2 \rangle \langle v_2^4 \rangle - \langle V_4 (V_2^*)^2 \rangle^2) \langle v_2^2 \rangle}, \\
 \chi_{7,25} &= \Re \frac{\langle V_7 V_2^* V_5^* \rangle \langle v_2^2 v_3^2 \rangle - \langle V_7 (V_2^*)^2 V_3^* \rangle \langle V_5 V_2^* V_3^* \rangle}{(\langle v_5^2 \rangle \langle v_2^2 v_3^2 \rangle - \langle V_5 V_2^* V_3^* \rangle^2) \langle v_2^2 \rangle}, \\
 \chi_{7,34} &= \Re \frac{\langle V_7 V_3^* V_4^* \rangle \langle v_2^4 \rangle - \langle V_7 (V_2^*)^2 V_3^* \rangle \langle V_4 (V_2^*)^2 \rangle}{(\langle v_4^2 \rangle \langle v_2^4 \rangle - \langle V_4 (V_2^*)^2 \rangle^2) \langle v_3^2 \rangle}.
 \end{aligned} \tag{A.4}$$

B The ALICE Collaboration

S. Acharya¹⁴¹, D. Adamová⁹⁴, A. Adler⁷⁴, J. Adolfsson⁸⁰, M.M. Aggarwal⁹⁹, G. Aglieri Rinella³³, M. Agnello³⁰, N. Agrawal^{10,53}, Z. Ahammed¹⁴¹, S. Ahmad¹⁶, S.U. Ahn⁷⁶, A. Akindinov⁹¹, M. Al-Turany¹⁰⁶, S.N. Alam¹⁴¹, D.S.D. Albuquerque¹²², D. Aleksandrov⁸⁷, B. Alessandro⁵⁸, H.M. Alfanda⁶, R. Alfaro Molina⁷¹, B. Ali¹⁶, Y. Ali¹⁴, A. Alici^{10,26,53}, A. Alkin², J. Alme²¹, T. Alt⁶⁸, L. Altenkamper²¹, I. Altsybeev¹¹², M.N. Anaam⁶, C. Andrei⁴⁷, D. Andreou³³, H.A. Andrews¹¹⁰, A. Andronic¹⁴⁴, M. Angeletti³³, V. Anguelov¹⁰³, C. Anson¹⁵, T. Antičić¹⁰⁷, F. Antinori⁵⁶, P. Antonioli⁵³, N. Apadula⁷⁹, L. Aphecetche¹¹⁴, H. Appelshäuser⁶⁸, S. Arcelli²⁶, R. Arnaldi⁵⁸, M. Arratia⁷⁹, I.C. Arsene²⁰, M. Arslandok¹⁰³, A. Augustinus³³, R. Averbeck¹⁰⁶, S. Aziz⁶¹, M.D. Azmi¹⁶, A. Badalá⁵⁵, Y.W. Baek⁴⁰, S. Bagnasco⁵⁸, X. Bai¹⁰⁶, R. Bailhache⁶⁸, R. Bala¹⁰⁰, A. Balbino³⁰, A. Baldisseri¹³⁷, M. Ball⁴², S. Balouza¹⁰⁴, R. Barbera²⁷, L. Barioglio²⁵, G.G. Barnaföldi¹⁴⁵, L.S. Barnby⁹³, V. Barret¹³⁴, P. Bartalini⁶, K. Barth³³, E. Bartsch⁶⁸, F. Baruffaldi²⁸, N. Bastid¹³⁴, S. Basu¹⁴³, G. Batigne¹¹⁴, B. Batyunya⁷⁵, D. Bauri⁴⁸, J.L. Bazo Alba¹¹¹, I.G. Bearden⁸⁸, C. Beattie¹⁴⁶, C. Bedda⁶³, N.K. Behera⁶⁰, I. Belikov¹³⁶, A.D.C. Bell Hechavarria¹⁴⁴, F. Bellini³³, R. Bellwied¹²⁵, V. Belyaev⁹², G. Bencedi¹⁴⁵, S. Beole²⁵, A. Bercuci⁴⁷, Y. Berdnikov⁹⁷, D. Berenyi¹⁴⁵, R.A. Bertens¹³⁰, D. Berzano⁵⁸, M.G. Besoiu⁶⁷, L. Betev³³, A. Bhasin¹⁰⁰, I.R. Bhat¹⁰⁰, M.A. Bhat³, H. Bhatt⁴⁸, B. Bhattacharjee⁴¹, A. Bianchi²⁵, L. Bianchi²⁵, N. Bianchi⁵¹, J. Bielčák³⁶, J. Bielčiková⁹⁴, A. Bilandzic^{104,117}, G. Biro¹⁴⁵, R. Biswas³, S. Biswas³, J.T. Blair¹¹⁹, D. Blau⁸⁷, C. Blume⁶⁸, G. Boca¹³⁹, F. Bock^{33,95}, A. Bogdanov⁹², L. Boldizsár¹⁴⁵, A. Bolozdynya⁹², M. Bombara³⁷, G. Bonomi¹⁴⁰, H. Borel¹³⁷, A. Borissov^{92,144}, H. Bossi¹⁴⁶, E. Botta²⁵, L. Bratrud⁶⁸, P. Braun-Munzinger¹⁰⁶, M. Bregant¹²¹, M. Broz³⁶, E. Bruna⁵⁸, G.E. Bruno¹⁰⁵, M.D. Buckland¹²⁷, D. Budnikov¹⁰⁸, H. Buesching⁶⁸, S. Bufalino³⁰, O. Bugnon¹¹⁴, P. Buhler¹¹³, P. Buncic³³, Z. Buthelezi^{72,131}, J.B. Butt¹⁴, J.T. Buxton⁹⁶, S.A. Bysiak¹¹⁸, D. Caffarri⁸⁹, A. Caliva¹⁰⁶, E. Calvo Villar¹¹¹, R.S. Camacho⁴⁴, P. Camerini²⁴, A.A. Capon¹¹³, F. Carnesecchi^{10,26}, R. Caron¹³⁷, J. Castillo Castellanos¹³⁷, A.J. Castro¹³⁰, E.A.R. Casula⁵⁴, F. Catalano³⁰, C. Ceballos Sanchez⁵², P. Chakraborty⁴⁸, S. Chandra¹⁴¹, W. Chang⁶, S. Chapeland³³, M. Chartier¹²⁷, S. Chattopadhyay¹⁴¹, S. Chattopadhyay¹⁰⁹, A. Chauvin²³, C. Cheshkov¹³⁵, B. Cheynis¹³⁵, V. Chibante Barroso³³, D.D. Chinellato¹²², S. Cho⁶⁰, P. Chochula³³, T. Chowdhury¹³⁴, P. Christakoglou⁸⁹, C.H. Christensen⁸⁸, P. Christiansen⁸⁰, T. Chujo¹³³, C. Cicalo⁵⁴, L. Cifarelli^{10,26}, F. Cindolo⁵³, G. Clai^{53,ii}, J. Cleymans¹²⁴, F. Colamaria⁵², D. Colella⁵², A. Collu⁷⁹, M. Colocci²⁶, M. Concas^{58,iii}, G. Conesa Balbastre⁷⁸, Z. Conesa del Valle⁶¹, G. Contin^{24,127}, J.G. Contreras³⁶, T.M. Cormier⁹⁵, Y. Corrales Morales²⁵, P. Cortese³¹, M.R. Cosentino¹²³, F. Costa³³, S. Costanza¹³⁹, P. Crochet¹³⁴, E. Cuautle⁶⁹, P. Cui⁶, L. Cunqueiro⁹⁵, D. Dabrowski¹⁴², T. Dahms^{104,117}, A. Dainese⁵⁶, F.P.A. Damas^{114,137}, M.C. Danisch¹⁰³, A. Danu⁶⁷, D. Das¹⁰⁹, I. Das¹⁰⁹, P. Das⁸⁵, P. Das³, S. Das³, A. Dash⁸⁵, S. Dash⁴⁸, S. De⁸⁵, A. De Caro²⁹, G. de Cataldo⁵², J. de Cuveland³⁸, A. De Falco²³, D. De Gruttola¹⁰, N. De Marco⁵⁸, S. De Pasquale²⁹, S. Deb⁴⁹, B. Debjani³, H.F. Degenhardt¹²¹, K.R. Deja¹⁴², A. Deloff⁶⁸⁴, S. Delsanto^{25,131}, W. Deng⁶, D. Devetak¹⁰⁶, P. Dhankher⁴⁸, D. Di Bari³², A. Di Mauro³³, R.A. Diaz⁸⁴, T. Dietel¹²⁴, P. Dillenseger⁶⁸, Y. Ding⁶, R. Divià³³, D.U. Dixit¹⁹, Ø. Djuvsland²¹, U. Dmitrieva⁶², A. Dobrin^{33,67}, B. Dönigus⁶⁸, O. Dordic²⁰, A.K. Dubey¹⁴¹, A. Dubla¹⁰⁶, S. Dudi⁹⁹, M. Dukhishyam⁸⁵, P. Dupieux¹³⁴, R.J. Ehlers^{95,146}, V.N. Eikeland²¹, D. Elia⁵², E. Epple¹⁴⁶, B. Erazmus¹¹⁴, F. Erhardt⁹⁸, A. Erokhin¹¹², M.R. Ersdal²¹, B. Espagnon⁶¹, G. Eulisse³³, D. Evans¹¹⁰, S. Evdokimov⁹⁰, L. Fabbietti^{104,117}, M. Faggin²⁸, J. Faivre⁷⁸, F. Fan⁶, A. Fantoni⁵¹, M. Fasel⁹⁵, P. Fedichio³⁰, A. Feliciello⁵⁸, G. Feofilov¹¹², A. Fernández Téllez⁴⁴, A. Ferrero¹³⁷, A. Ferretti²⁵, A. Festanti³³, V.J.G. Feuillard¹⁰³, J. Figiel¹¹⁸, S. Filchagin¹⁰⁸, D. Finogeev⁶², F.M. Fionda²¹, G. Fiorenza⁵², F. Flor¹²⁵, S. Foertsch²⁹, P. Foka¹⁰⁶, S. Fokin⁸⁷, E. Fragiaco⁵⁹, U. Frankendorf¹⁰⁶, U. Fuchs³³, C. Furget⁷⁸, A. Furs⁶², M. Fusco Girard²⁹, J.J. Gaardhøje⁸⁸, M. Gagliardi²⁵, A.M. Gago¹¹¹, A. Gal¹³⁶, C.D. Galvan¹²⁰, P. Ganoti⁸³, C. Garabatos¹⁰⁶, E. Garcia-Solis¹¹, K. Garg²⁷, C. Gargiulo³³, A. Garibli⁸⁶, K. Garner¹⁴⁴, P. Gasik^{104,117}, E.F. Gauger¹¹⁹, M.B. Gay Ducati⁷⁰, M. Germain¹¹⁴, J. Ghosh¹⁰⁹, P. Ghosh¹⁴¹, S.K. Ghosh³, M. Giacalone²⁶, P. Gianotti⁵¹, P. Giubellino^{58,106}, P. Giubilato²⁸, P. Glässel¹⁰³, A. Gomez Ramirez⁷⁴, V. Gonzalez¹⁰⁶, P. González-Zamora⁴⁴, S. Gorbunov³⁸, L. Görlich¹¹⁸, S. Gotovac³⁴, V. Grabski⁷¹, L.K. Graczykowski¹⁴², K.L. Graham¹¹⁰, L. Greiner⁷⁹, A. Grelli⁶³, C. Grigoras³³, V. Grigoriev⁹², A. Grigoryan¹, S. Grigoryan⁷⁵, O.S. Groettvik²¹, F. Grosa³⁰, J.F. Grosse-Oetringhaus³³, R. Grosso¹⁰⁶, R. Guernane⁷⁸, M. Guittiere¹¹⁴, K. Gulbrandsen⁸⁸, T. Gunji¹³², A. Gupta¹⁰⁰, R. Gupta¹⁰⁰, I.B. Guzman⁴⁴, R. Haake¹⁴⁶, M.K. Habib¹⁰⁶, C. Hadjidakis⁶¹, H. Hamagaki⁸¹, G. Hamar¹⁴⁵, M. Hamid⁶, R. Hannigan¹¹⁹, M.R. Haque^{63,85}, A. Harlanderova¹⁰⁶, J.W. Harris¹⁴⁶, A. Harton¹¹, J.A. Hasenbichler³³, H. Hassan⁹⁵, D. Hatzifotiadou^{10,53}, P. Hauer⁴², S. Hayashi¹³², S.T. Heckel^{68,104}, E. Hellbär⁶⁸, H. Helstrup³⁵, A. Herghelegiu⁴⁷, T. Herman³⁶, E.G. Hernandez⁴⁴, G. Herrera Corral⁹, F. Herrmann¹⁴⁴, K.F. Hetland³⁵, H. Hillemanns³³, C. Hills¹²⁷, B. Hippolyte¹³⁶, B. Hohlweger¹⁰⁴, J. Honermann¹⁴⁴, D. Horak³⁶, A. Hornung⁶⁸, S. Hornung¹⁰⁶, R. Hosokawa¹⁵, P. Hristov³³, C. Huang⁶¹, C. Hughes¹³⁰, P. Huhn⁶⁸, T.J. Humanic⁹⁶, H. Hushnud¹⁰⁹,

L.A. Husova¹⁴⁴, N. Hussain⁴¹, S.A. Hussain¹⁴, D. Hutter³⁸, J.P. Iddon^{33,127}, R. Ilkaev¹⁰⁸, M. Inaba¹³³, G.M. Innocenti³³, M. Ippolitov⁸⁷, A. Isakov⁹⁴, M.S. Islam¹⁰⁹, M. Ivanov¹⁰⁶, V. Ivanov⁹⁷, V. Izucheev⁹⁰, B. Jacak⁷⁹, N. Jacazio⁵³, P.M. Jacobs⁷⁹, S. Jadlovská¹¹⁶, J. Jadlovsky¹¹⁶, S. Jaelani⁶³, C. Jahnke¹²¹, M.J. Jakubowska¹⁴², M.A. Janik¹⁴², T. Janson⁷⁴, M. Jercic⁹⁸, O. Jevons¹¹⁰, M. Jin¹²⁵, F. Jonas^{95,144}, P.G. Jones¹¹⁰, J. Jung⁶⁸, M. Jung⁶⁸, A. Jusko¹¹⁰, P. Kalinac⁶⁴, A. Kalweit³³, V. Kaplin⁹², S. Kar⁶, A. Karasu Uysal⁷⁷, O. Karavichev⁶², T. Karavicheva⁶², P. Karczmarczyk³³, E. Karpechev⁶², U. Kebschull⁷⁴, R. Keidel⁴⁶, M. Keil³³, B. Ketzer⁴², Z. Khabanova⁸⁹, A.M. Khan⁶, S. Khan¹⁶, S.A. Khan¹⁴¹, A. Khanzadeev⁹⁷, Y. Kharlov⁹⁰, A. Khatun¹⁶, A. Khuntia¹¹⁸, B. Kileng³⁵, B. Kim⁶⁰, B. Kim¹³³, D. Kim¹⁴⁷, D.J. Kim¹²⁶, E.J. Kim⁷³, H. Kim^{17,147}, J. Kim¹⁴⁷, J.S. Kim⁴⁰, J. Kim¹⁰³, J. Kim¹⁴⁷, J. Kim⁷³, M. Kim¹⁰³, S. Kim¹⁸, T. Kim¹⁴⁷, T. Kim¹⁴⁷, S. Kirsch^{38,68}, I. Kisel³⁸, S. Kiselev⁹¹, A. Kisiel¹⁴², J.L. Klay⁵, C. Klein⁶⁸, J. Klein⁵⁸, S. Klein⁷⁹, C. Klein-Bösing¹⁴⁴, M. Kleiner⁶⁸, A. Kluge³³, M.L. Knichel³³, A.G. Knospe¹²⁵, C. Kobdaj¹¹⁵, M.K. Köhler¹⁰³, T. Kollegger¹⁰⁶, A. Kondratyev⁷⁵, N. Kondratyeva⁹², E. Kondratyuk⁹⁰, J. König⁶⁸, P.J. Konopka³³, L. Koska¹¹⁶, O. Kovalenko⁸⁴, V. Kovalenko¹¹², M. Kowalski¹¹⁸, I. Králik⁶⁴, A. Kravčáková³⁷, L. Kreis¹⁰⁶, M. Krivda^{64,110}, F. Krizek⁹⁴, K. Krizkova Gajdosova³⁶, M. Krüger⁶⁸, E. Kryshen⁹⁷, M. Krzewicki³⁸, A.M. Kubera⁹⁶, V. Kučera⁶⁰, C. Kuhn¹³⁶, P.G. Kuijter⁸⁹, L. Kumar⁹⁹, S. Kundu⁸⁵, P. Kurashvili⁸⁴, A. Kurepin⁶², A.B. Kurepin⁶², A. Kuryakin¹⁰⁸, S. Kuschpil⁹⁴, J. Kvapil¹¹⁰, M.J. Kweon⁶⁰, J.Y. Kwon⁶⁰, Y. Kwon¹⁴⁷, S.L. La Pointe³⁸, P. La Rocca²⁷, Y.S. Lai⁷⁹, R. Langoy¹²⁹, K. Lapidus³³, A. Lardeux²⁰, P. Larionov⁵¹, E. Laudi³³, R. Lavicka³⁶, T. Lazareva¹¹², R. Lea²⁴, L. Leardini¹⁰³, J. Lee¹³³, S. Lee¹⁴⁷, F. Lehas⁸⁹, S. Lehner¹¹³, J. Lehrbach³⁸, R.C. Lemmon⁹³, I. León Monzón¹²⁰, E.D. Lesser¹⁹, M. Lettrich³³, P. Lévai¹⁴⁵, X. Li¹², X.L. Li⁶, J. Lien¹²⁹, R. Lietava¹¹⁰, B. Lim¹⁷, V. Lindenstruth³⁸, S.W. Lindsay¹²⁷, C. Lippmann¹⁰⁶, M.A. Lisa⁹⁶, A. Liu¹⁹, J. Liu¹²⁷, S. Liu⁹⁶, W.J. Llope¹⁴³, I.M. Lofnes²¹, V. Loginov⁹², C. Loizides⁹⁵, P. Loncar³⁴, J.A.L. Lopez¹⁰³, X. Lopez¹³⁴, E. López Torres⁸, J.R. Luhder¹⁴⁴, M. Lunardon²⁸, G. Luparello⁵⁹, Y.G. Ma³⁹, A. Maevskaya⁶², M. Mager³³, S.M. Mahmood²⁰, T. Mahmoud⁴², A. Maire¹³⁶, R.D. Majka¹⁴⁶, M. Malaev⁹⁷, Q.W. Malik²⁰, L. Malinina^{75,iv}, D. Mal'Kevich⁹¹, P. Malzacher¹⁰⁶, G. Mandaglio⁵⁵, V. Manko⁸⁷, F. Manso¹³⁴, V. Manzari⁵², Y. Mao⁶, M. Marchisone¹³⁵, J. Mareš⁶⁶, G.V. Margagliotti²⁴, A. Margotti⁵³, J. Margutti⁶³, A. Marín¹⁰⁶, C. Markert¹¹⁹, M. Marquard⁶⁸, N.A. Martin¹⁰³, P. Martinengo³³, J.L. Martinez¹²⁵, M.I. Martínez⁴⁴, G. Martínez García¹¹⁴, M. Martinez Pedreira³³, S. Masciocchi¹⁰⁶, M. Maserà²⁵, A. Masoni⁵⁴, L. Massacrier⁶¹, E. Masson¹¹⁴, A. Mastroserio^{52,138}, A.M. Mathis^{104,117}, O. Matonoha⁸⁰, P.F.T. Matuoka¹²¹, A. Matyja¹¹⁸, C. Mayer¹¹⁸, F. Mazzaschi²⁵, M. Mazzilli⁵², M.A. Mazzoni⁵⁷, A.F. Mechler⁶⁸, F. Meddi²², Y. Melikyan^{62,92}, A. Menchaca-Rocha⁷¹, C. Mengke⁶, E. Meninno^{29,113}, M. Meres¹³, S. Mhlanga¹²⁴, Y. Miake¹³³, L. Micheletti²⁵, D.L. Mihaylov¹⁰⁴, K. Mikhaylov^{75,91}, A.N. Mishra⁶⁹, D. Miśkowiec¹⁰⁶, A. Modak³, N. Mohammadi³³, A.P. Mohanty⁶³, B. Mohanty⁸⁵, M. Mohisin Khan^{16,v}, C. Mordasini¹⁰⁴, D.A. Moreira De Godoy¹⁴⁴, L.A.P. Moreno⁴⁴, I. Morozov⁶², A. Morsch³³, T. Mrnjavac³³, V. Muccifora⁵¹, E. Mudnic³⁴, D. Mühlheim¹⁴⁴, S. Muhuri¹⁴¹, J.D. Mulligan⁷⁹, M.G. Munhoz¹²¹, R.H. Munzer⁶⁸, H. Murakami¹³², S. Murray¹²⁴, L. Musa³³, J. Musinsky⁶⁴, C.J. Myers¹²⁵, J.W. Myrcha¹⁴², B. Naik⁴⁸, R. Nair⁸⁴, B.K. Nandi⁴⁸, R. Nania^{10,53}, E. Nappi⁵², M.U. Naru¹⁴, A.F. Nassirpour⁸⁰, C. Nattrass¹³⁰, R. Nayak⁴⁸, T.K. Nayak⁸⁵, S. Nazarenko¹⁰⁸, A. Neagu²⁰, R.A. Negrao De Oliveira⁶⁸, L. Nellen⁶⁹, S.V. Nesbo³⁵, G. Neskovic³⁸, D. Nesterov¹¹², L.T. Neumann¹⁴², B.S. Nielsen⁸⁸, S. Nikolaev⁸⁷, S. Nikulin⁸⁷, V. Nikulin⁹⁷, F. Noferini^{10,53}, P. Nomokonov⁷⁵, J. Norman^{78,127}, N. Novitzky¹³³, P. Nowakowski¹⁴², A. Nyanin⁸⁷, J. Nystrand²¹, M. Ogino⁸¹, A. Ohlson^{80,103}, J. Olińczak¹⁴², A.C. Oliveira Da Silva^{121,130}, M.H. Oliver¹⁴⁶, C. Oppedisano⁵⁸, R. Orava⁴³, A. Ortiz Velasquez⁶⁹, A. Oskarsson⁸⁰, J. Otwinowski¹¹⁸, K. Oyama⁸¹, Y. Pachmayer¹⁰³, V. Pacik⁸⁸, D. Pagano¹⁴⁰, G. Paic⁶⁹, J. Pan¹⁴³, A.K. Pandey⁴⁸, S. Panebianco¹³⁷, P. Pareek^{49,141}, J. Park⁶⁰, J.E. Parkkila¹²⁶, S. Parmar⁹⁹, S.P. Pathak¹²⁵, R.N. Patra¹⁴¹, B. Paul²³, H. Pei⁶, T. Peitzmann⁶³, X. Peng⁶, L.G. Pereira⁷⁰, H. Pereira Da Costa¹³⁷, D. Peresunko⁸⁷, G.M. Perez⁸, E. Perez Lezama⁶⁸, V. Peskov⁶⁸, Y. Pestov⁴, V. Petráček³⁶, M. Petrovici⁴⁷, R.P. Pezzi⁷⁰, S. Piano⁵⁹, M. Pikna¹³, P. Pillot¹¹⁴, O. Pinazza^{33,53}, L. Pinsky¹²⁵, C. Pinto²⁷, S. Pisano^{10,51}, D. Pistone⁵⁵, M. Płoskoń⁷⁹, M. Planinic⁹⁸, F. Pliquett⁶⁸, S. Pochybova^{145,i}, M.G. Poghosyan⁹⁵, B. Polichtchouk⁹⁰, N. Poljak⁹⁸, A. Pop⁴⁷, H. Poppenborg¹⁴⁴, S. Porteboeuf-Houssais¹³⁴, V. Pozdniakov⁷⁵, S.K. Prasad³, R. Preghenella⁵³, F. Prino⁵⁸, C.A. Pruneau¹⁴³, I. Pshenichnov⁶², M. Puccio^{25,33}, J. Putschke¹⁴³, L. Quaglia²⁵, R.E. Quishpe¹²⁵, S. Ragoni¹¹⁰, S. Raha³, S. Rajput¹⁰⁰, J. Rak¹²⁶, A. Rakotozafindrabe¹³⁷, L. Ramello³¹, F. Rami¹³⁶, R. Raniwala¹⁰¹, S. Raniwala¹⁰¹, S.S. Räsänen⁴³, R. Rath⁴⁹, V. Ratz⁴², I. Ravasenga⁸⁹, K.F. Read^{95,130}, A.R. Redelbach³⁸, K. Redlich^{84,vi}, A. Rehman²¹, P. Reichelt⁶⁸, F. Reidt³³, X. Ren⁶, R. Renfordt⁶⁸, Z. Rescakova³⁷, J.-P. Revol¹⁰, K. Reygers¹⁰³, V. Riabov⁹⁷, T. Richert^{80,88}, M. Richter²⁰, P. Riedler³³, W. Riegler³³, F. Riggi²⁷, C. Ristea⁶⁷, S.P. Rode⁴⁹, M. Rodríguez Cahuantzi⁴⁴, K. Røed²⁰, R. Rogalev⁹⁰, E. Rogochaya⁷⁵, D. Rohr³³, D. Röhrich²¹, P.S. Rokita¹⁴², F. Ronchetti⁵¹, E.D. Rosas⁶⁹, K. Roslon¹⁴², A. Rossi^{28,56}, A. Rotondi¹³⁹, A. Roy⁴⁹, P. Roy¹⁰⁹, O.V. Rueda⁸⁰, R. Rui²⁴,

B. Rumyantsev⁷⁵, A. Rustamov⁸⁶, E. Ryabinkin⁸⁷, Y. Ryabov⁹⁷, A. Rybicki¹¹⁸, H. Rytönen¹²⁶, O.A.M. Saarimäki⁴³, S. Sadhu¹⁴¹, S. Sadovsky⁹⁰, K. Šafařík³⁶, S.K. Saha¹⁴¹, B. Sahoo⁴⁸, P. Sahoo⁴⁸, R. Sahoo⁴⁹, S. Sahoo⁶⁵, P.K. Sahu⁶⁵, J. Saini¹⁴¹, S. Sakai¹³³, S. Sambyal¹⁰⁰, V. Samsonov^{92,97}, D. Sarkar¹⁴³, N. Sarkar¹⁴¹, P. Sarma⁴¹, V.M. Sarti¹⁰⁴, M.H.P. Sas⁶³, E. Scapparone⁵³, B. Schaefer⁹⁵, J. Schambach¹¹⁹, H.S. Scheid⁶⁸, C. Schiaua⁴⁷, R. Schicker¹⁰³, A. Schmah¹⁰³, C. Schmidt¹⁰⁶, H.R. Schmidt¹⁰², M.O. Schmidt¹⁰³, M. Schmidt¹⁰², N.V. Schmidt^{68,95}, A.R. Schmier¹³⁰, J. Schukraft⁸⁸, Y. Schutz^{33,136}, K. Schwarz¹⁰⁶, K. Schweda¹⁰⁶, G. Scioli²⁶, E. Scomparin⁵⁸, M. Šefčík³⁷, J.E. Seger¹⁵, Y. Sekiguchi¹³², D. Sekihata¹³², I. Selyuzhenkov^{92,106}, S. Senyukov¹³⁶, D. Serebryakov⁶², E. Serradilla⁷¹, A. Sevcenco⁶⁷, A. Shabanov⁶², A. Shabetai¹¹⁴, R. Shahoyan³³, W. Shaikh¹⁰⁹, A. Shangaraev⁹⁰, A. Sharma⁹⁹, A. Sharma¹⁰⁰, H. Sharma¹¹⁸, M. Sharma¹⁰⁰, N. Sharma⁹⁹, S. Sharma¹⁰⁰, A.I. Sheikh¹⁴¹, K. Shigaki⁴⁵, M. Shimomura⁸², S. Shirinkin⁹¹, Q. Shou³⁹, Y. Sibiriak⁸⁷, S. Siddhanta⁵⁴, T. Siemiarczuk⁸⁴, D. Silvermyr⁸⁰, G. Simatovic⁸⁹, G. Simonetti^{33,104}, R. Singh⁸⁵, R. Singh¹⁰⁰, R. Singh⁴⁹, V.K. Singh¹⁴¹, V. Singhal¹⁴¹, T. Sinha¹⁰⁹, B. Sitar¹³, M. Sitta³¹, T.B. Skaali²⁰, M. Slupecki¹²⁶, N. Smirnov¹⁴⁶, R.J.M. Snellings⁶³, T.W. Snellman^{43,126}, C. Soncco¹¹¹, J. Song^{60,125}, A. Songmoolnak¹¹⁵, F. Soramel²⁸, S. Sorensen¹³⁰, I. Sputowska¹¹⁸, J. Stachel¹⁰³, I. Stan⁶⁷, P. Stankus⁹⁵, P.J. Steffanic¹³⁰, E. Stenlund⁸⁰, D. Stocco¹¹⁴, M.M. Stortvedt³⁵, L.D. Stritto²⁹, A.A.P. Suaide¹²¹, T. Sugitate⁴⁵, C. Suire⁶¹, M. Suleymanov¹⁴, M. Suljic³³, R. Sultanov⁹¹, M. Šumbera⁹⁴, V. Sumberia¹⁰⁰, S. Sumowidagdo⁵⁰, S. Swain⁶⁵, A. Szabo¹³, I. Szarka¹³, U. Tabassam¹⁴, S.F. Taghavi¹⁰⁴, G. Taillepied¹³⁴, J. Takahashi¹²², G.J. Tambave²¹, S. Tang^{6,134}, M. Tarhini¹¹⁴, M.G. Tarzila⁴⁷, A. Tauro³³, G. Tejada Muñoz⁴⁴, A. Telesca³³, L. Terlizzi²⁵, C. Terrevoli¹²⁵, D. Thakur⁴⁹, S. Thakur¹⁴¹, D. Thomas¹¹⁹, F. Thoresen⁸⁸, R. Tieulent¹³⁵, A. Tikhonov⁶², A.R. Timmins¹²⁵, A. Toia⁶⁸, N. Topilskaya⁶², M. Toppi⁵¹, F. Torres-Acosta¹⁹, S.R. Torres^{36,120}, A. Trifiro⁵⁵, S. Tripathy⁴⁹, T. Tripathy⁴⁸, S. Trogolo²⁸, G. Trombetta³², L. Tropp³⁷, V. Trubnikov², W.H. Trzaska¹²⁶, T.P. Trzcinski¹⁴², B.A. Trzeciak^{36,63}, T. Tsuji¹³², A. Tumkin¹⁰⁸, R. Turrisi⁵⁶, T.S. Tveter²⁰, K. Ullaland²¹, E.N. Umaka¹²⁵, A. Uras¹³⁵, G.L. Usai²³, A. Utrobicic⁹⁸, M. Vala³⁷, N. Valle¹³⁹, S. Vallero⁵⁸, N. van der Kolk⁶³, L.V.R. van Doremalen⁶³, M. van Leeuwen⁶³, P. Vande Vyvre³³, D. Varga¹⁴⁵, Z. Varga¹⁴⁵, M. Varga-Kofarago¹⁴⁵, A. Vargas⁴⁴, M. Vasileiou⁸³, A. Vasiliev⁸⁷, O. Vázquez Doce^{104,117}, V. Vechernin¹¹², A.M. Veen⁶³, E. Vercellin²⁵, S. Vergara Limón⁴⁴, L. Vermunt⁶³, R. Vernet⁷, R. Vértesi¹⁴⁵, L. Vickovic³⁴, Z. Vilakazi¹³¹, O. Villalobos Baillie¹¹⁰, A. Villatoro Tello⁴⁴, G. Vino⁵², A. Vinogradov⁸⁷, T. Virgili²⁹, V. Vislavicius⁸⁸, A. Vodopyanov⁷⁵, B. Volkel³³, M.A. Völkl¹⁰², K. Voloshin⁹¹, S.A. Voloshin¹⁴³, G. Volpe³², B. von Haller³³, I. Vorobyev¹⁰⁴, D. Voscek¹¹⁶, J. Vrláková³⁷, B. Wagner²¹, M. Weber¹¹³, A. Wegrzynek³³, D.F. Weiser¹⁰³, S.C. Wenzel³³, J.P. Wessels¹⁴⁴, J. Wiechula⁶⁸, J. Wikne²⁰, G. Wilk⁸⁴, J. Wilkinson^{10,53}, G.A. Willems¹⁴⁴, E. Willsher¹¹⁰, B. Windelband¹⁰³, M. Winn¹³⁷, W.E. Witt¹³⁰, Y. Wu¹²⁸, R. Xu⁶, S. Yalcin⁷⁷, Y. Yamaguchi⁴⁵, K. Yamakawa⁴⁵, S. Yang²¹, S. Yano¹³⁷, Z. Yin⁶, H. Yokoyama⁶³, I.-K. Yoo¹⁷, J.H. Yoon⁶⁰, S. Yuan²¹, A. Yuncu¹⁰³, V. Yurchenko², V. Zaccolo²⁴, A. Zaman¹⁴, C. Zampolli³³, H.J.C. Zanoli⁶³, N. Zardoshti³³, A. Zarochentsev¹¹², P. Závada⁶⁶, N. Zaviyalov¹⁰⁸, H. Zbroszczyk¹⁴², M. Zhalov⁹⁷, S. Zhang³⁹, X. Zhang⁶, Z. Zhang⁶, V. Zhrebchevskii¹¹², D. Zhou⁶, Y. Zhou⁸⁸, Z. Zhou²¹, J. Zhu^{6,106}, Y. Zhu⁶, A. Zichichi^{10,26}, M.B. Zimmermann³³, G. Zinovjev², N. Zurlo¹⁴⁰,

Affiliation notes

- ⁱ Deceased
- ⁱⁱ Italian National Agency for New Technologies, Energy and Sustainable Economic Development (ENEA), Bologna, Italy
- ⁱⁱⁱ Dipartimento DET del Politecnico di Torino, Turin, Italy
- ^{iv} M.V. Lomonosov Moscow State University, D.V. Skobeltsyn Institute of Nuclear Physics, Moscow, Russia
- ^v Department of Applied Physics, Aligarh Muslim University, Aligarh, India
- ^{vi} Institute of Theoretical Physics, University of Wrocław, Poland

Collaboration Institutes

- ¹ A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia
- ² Bogolyubov Institute for Theoretical Physics, National Academy of Sciences of Ukraine, Kiev, Ukraine
- ³ Bose Institute, Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Kolkata, India
- ⁴ Budker Institute for Nuclear Physics, Novosibirsk, Russia
- ⁵ California Polytechnic State University, San Luis Obispo, California, United States
- ⁶ Central China Normal University, Wuhan, China

- 7 Centre de Calcul de l'IN2P3, Villeurbanne, Lyon, France
- 8 Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba
- 9 Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico
- 10 Centro Fermi - Museo Storico della Fisica e Centro Studi e Ricerche "Enrico Fermi", Rome, Italy
- 11 Chicago State University, Chicago, Illinois, United States
- 12 China Institute of Atomic Energy, Beijing, China
- 13 Comenius University Bratislava, Faculty of Mathematics, Physics and Informatics, Bratislava, Slovakia
- 14 COMSATS University Islamabad, Islamabad, Pakistan
- 15 Creighton University, Omaha, Nebraska, United States
- 16 Department of Physics, Aligarh Muslim University, Aligarh, India
- 17 Department of Physics, Pusan National University, Pusan, Republic of Korea
- 18 Department of Physics, Sejong University, Seoul, Republic of Korea
- 19 Department of Physics, University of California, Berkeley, California, United States
- 20 Department of Physics, University of Oslo, Oslo, Norway
- 21 Department of Physics and Technology, University of Bergen, Bergen, Norway
- 22 Dipartimento di Fisica dell'Università 'La Sapienza' and Sezione INFN, Rome, Italy
- 23 Dipartimento di Fisica dell'Università and Sezione INFN, Cagliari, Italy
- 24 Dipartimento di Fisica dell'Università and Sezione INFN, Trieste, Italy
- 25 Dipartimento di Fisica dell'Università and Sezione INFN, Turin, Italy
- 26 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Bologna, Italy
- 27 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Catania, Italy
- 28 Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Padova, Italy
- 29 Dipartimento di Fisica 'E.R. Caianiello' dell'Università and Gruppo Collegato INFN, Salerno, Italy
- 30 Dipartimento DISAT del Politecnico and Sezione INFN, Turin, Italy
- 31 Dipartimento di Scienze e Innovazione Tecnologica dell'Università del Piemonte Orientale and INFN Sezione di Torino, Alessandria, Italy
- 32 Dipartimento Interateneo di Fisica 'M. Merlin' and Sezione INFN, Bari, Italy
- 33 European Organization for Nuclear Research (CERN), Geneva, Switzerland
- 34 Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, University of Split, Split, Croatia
- 35 Faculty of Engineering and Science, Western Norway University of Applied Sciences, Bergen, Norway
- 36 Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
- 37 Faculty of Science, P.J. Šafárik University, Košice, Slovakia
- 38 Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 39 Fudan University, Shanghai, China
- 40 Gangneung-Wonju National University, Gangneung, Republic of Korea
- 41 Gauhati University, Department of Physics, Guwahati, India
- 42 Helmholtz-Institut für Strahlen- und Kernphysik, Rheinische Friedrich-Wilhelms-Universität Bonn, Bonn, Germany
- 43 Helsinki Institute of Physics (HIP), Helsinki, Finland
- 44 High Energy Physics Group, Universidad Autónoma de Puebla, Puebla, Mexico
- 45 Hiroshima University, Hiroshima, Japan
- 46 Hochschule Worms, Zentrum für Technologietransfer und Telekommunikation (ZTT), Worms, Germany
- 47 Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania
- 48 Indian Institute of Technology Bombay (IIT), Mumbai, India
- 49 Indian Institute of Technology Indore, Indore, India
- 50 Indonesian Institute of Sciences, Jakarta, Indonesia
- 51 INFN, Laboratori Nazionali di Frascati, Frascati, Italy
- 52 INFN, Sezione di Bari, Bari, Italy
- 53 INFN, Sezione di Bologna, Bologna, Italy
- 54 INFN, Sezione di Cagliari, Cagliari, Italy
- 55 INFN, Sezione di Catania, Catania, Italy
- 56 INFN, Sezione di Padova, Padova, Italy
- 57 INFN, Sezione di Roma, Rome, Italy

- 58 INFN, Sezione di Torino, Turin, Italy
- 59 INFN, Sezione di Trieste, Trieste, Italy
- 60 Inha University, Incheon, Republic of Korea
- 61 Institut de Physique Nucléaire d'Orsay (IPNO), Institut National de Physique Nucléaire et de Physique des Particules (IN2P3/CNRS), Université de Paris-Sud, Université Paris-Saclay, Orsay, France
- 62 Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
- 63 Institute for Subatomic Physics, Utrecht University/Nikhef, Utrecht, Netherlands
- 64 Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia
- 65 Institute of Physics, Homi Bhabha National Institute, Bhubaneswar, India
- 66 Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
- 67 Institute of Space Science (ISS), Bucharest, Romania
- 68 Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- 69 Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 70 Instituto de Física, Universidade Federal do Rio Grande do Sul (UFRGS), Porto Alegre, Brazil
- 71 Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico
- 72 iThemba LABS, National Research Foundation, Somerset West, South Africa
- 73 Jeonbuk National University, Jeonju, Republic of Korea
- 74 Johann-Wolfgang-Goethe Universität Frankfurt Institut für Informatik, Fachbereich Informatik und Mathematik, Frankfurt, Germany
- 75 Joint Institute for Nuclear Research (JINR), Dubna, Russia
- 76 Korea Institute of Science and Technology Information, Daejeon, Republic of Korea
- 77 KTO Karatay University, Konya, Turkey
- 78 Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS-IN2P3, Grenoble, France
- 79 Lawrence Berkeley National Laboratory, Berkeley, California, United States
- 80 Lund University Department of Physics, Division of Particle Physics, Lund, Sweden
- 81 Nagasaki Institute of Applied Science, Nagasaki, Japan
- 82 Nara Women's University (NWU), Nara, Japan
- 83 National and Kapodistrian University of Athens, School of Science, Department of Physics, Athens, Greece
- 84 National Centre for Nuclear Research, Warsaw, Poland
- 85 National Institute of Science Education and Research, Homi Bhabha National Institute, Jatni, India
- 86 National Nuclear Research Center, Baku, Azerbaijan
- 87 National Research Centre Kurchatov Institute, Moscow, Russia
- 88 Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
- 89 Nikhef, National institute for subatomic physics, Amsterdam, Netherlands
- 90 NRC Kurchatov Institute IHEP, Protvino, Russia
- 91 NRC Kurchatov Institute - ITEP, Moscow, Russia
- 92 NRNU Moscow Engineering Physics Institute, Moscow, Russia
- 93 Nuclear Physics Group, STFC Daresbury Laboratory, Daresbury, United Kingdom
- 94 Nuclear Physics Institute of the Czech Academy of Sciences, Řež u Prahy, Czech Republic
- 95 Oak Ridge National Laboratory, Oak Ridge, Tennessee, United States
- 96 Ohio State University, Columbus, Ohio, United States
- 97 Petersburg Nuclear Physics Institute, Gatchina, Russia
- 98 Physics department, Faculty of science, University of Zagreb, Zagreb, Croatia
- 99 Physics Department, Panjab University, Chandigarh, India
- 100 Physics Department, University of Jammu, Jammu, India
- 101 Physics Department, University of Rajasthan, Jaipur, India
- 102 Physikalisches Institut, Eberhard-Karls-Universität Tübingen, Tübingen, Germany
- 103 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
- 104 Physik Department, Technische Universität München, Munich, Germany
- 105 Politecnico di Bari, Bari, Italy
- 106 Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, Darmstadt, Germany
- 107 Rudjer Bošković Institute, Zagreb, Croatia
- 108 Russian Federal Nuclear Center (VNIIEF), Sarov, Russia

- 109 Saha Institute of Nuclear Physics, Homi Bhabha National Institute, Kolkata, India
- 110 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
- 111 Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
- 112 St. Petersburg State University, St. Petersburg, Russia
- 113 Stefan Meyer Institut für Subatomare Physik (SMI), Vienna, Austria
- 114 SUBATECH, IMT Atlantique, Université de Nantes, CNRS-IN2P3, Nantes, France
- 115 Suranaree University of Technology, Nakhon Ratchasima, Thailand
- 116 Technical University of Košice, Košice, Slovakia
- 117 Technische Universität München, Excellence Cluster 'Universe', Munich, Germany
- 118 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Cracow, Poland
- 119 The University of Texas at Austin, Austin, Texas, United States
- 120 Universidad Autónoma de Sinaloa, Culiacán, Mexico
- 121 Universidade de São Paulo (USP), São Paulo, Brazil
- 122 Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
- 123 Universidade Federal do ABC, Santo Andre, Brazil
- 124 University of Cape Town, Cape Town, South Africa
- 125 University of Houston, Houston, Texas, United States
- 126 University of Jyväskylä, Jyväskylä, Finland
- 127 University of Liverpool, Liverpool, United Kingdom
- 128 University of Science and Technology of China, Hefei, China
- 129 University of South-Eastern Norway, Tonsberg, Norway
- 130 University of Tennessee, Knoxville, Tennessee, United States
- 131 University of the Witwatersrand, Johannesburg, South Africa
- 132 University of Tokyo, Tokyo, Japan
- 133 University of Tsukuba, Tsukuba, Japan
- 134 Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France
- 135 Université de Lyon, Université Lyon 1, CNRS/IN2P3, IPN-Lyon, Villeurbanne, Lyon, France
- 136 Université de Strasbourg, CNRS, IPHC UMR 7178, F-67000 Strasbourg, France, Strasbourg, France
- 137 Université Paris-Saclay Centre d'Etudes de Saclay (CEA), IRFU, Département de Physique Nucléaire (DPhN), Saclay, France
- 138 Università degli Studi di Foggia, Foggia, Italy
- 139 Università degli Studi di Pavia, Pavia, Italy
- 140 Università di Brescia, Brescia, Italy
- 141 Variable Energy Cyclotron Centre, Homi Bhabha National Institute, Kolkata, India
- 142 Warsaw University of Technology, Warsaw, Poland
- 143 Wayne State University, Detroit, Michigan, United States
- 144 Westfälische Wilhelms-Universität Münster, Institut für Kernphysik, Münster, Germany
- 145 Wigner Research Centre for Physics, Budapest, Hungary
- 146 Yale University, New Haven, Connecticut, United States
- 147 Yonsei University, Seoul, Republic of Korea