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Jet fragmentation transverse momentum distributions in pp and p–Pb collisions at $\sqrt{s}, \sqrt{s_{NN}} = 5.02$ TeV

ALICE Collaboration*

Abstract

Jet fragmentation transverse momentum (j_T) distributions are measured in proton-proton (pp) and proton-lead (p–Pb) collisions at $\sqrt{s_{NN}} = 5.02$ TeV with the ALICE experiment at the LHC. Jets are reconstructed with the ALICE tracking detectors and electromagnetic calorimeter using the anti- k_T algorithm with resolution parameter $R = 0.4$ in the pseudorapidity range $|\eta| < 0.25$. The j_T values are calculated for charged particles inside a fixed cone with a radius $R = 0.4$ around the reconstructed jet axis. The measured j_T distributions are compared with a variety of parton-shower models. Herwig and PYTHIA 8 based models describe the data well for the higher j_T region, while they underestimate the lower j_T region. The j_T distributions are further characterised by fitting them with a function composed of an inverse gamma function for higher j_T values (called the “wide component”), related to the perturbative component of the fragmentation process, and with a Gaussian for lower j_T values (called the “narrow component”), predominantly connected to the hadronisation process. The width of the Gaussian has only a weak dependence on jet transverse momentum, while that of the inverse gamma function increases with increasing jet transverse momentum. For the narrow component, the measured trends are successfully described by all models except for Herwig. For the wide component, Herwig and PYTHIA 8 based models slightly underestimate the data for the higher jet transverse momentum region. These measurements set constraints on models of jet fragmentation and hadronisation.

1 Introduction

Jets are groups of collimated particles mainly resulting from fragmentation of hard scattered partons produced in high-energy particle collisions. Jet production in quantum chromodynamics (QCD) [1–5] can be thought as a two-stage process [6]. After being produced in the hard scattering, partons reduce their virtuality by emitting gluons [7]. Since the momentum transfer scale (Q^2) is large during the showering, perturbative QCD calculations can be applied. When Q^2 becomes of the order of Λ_{QCD} , partons hadronise into final-state particles through processes that cannot be calculated perturbatively [8–14]. Instead, the implementation of specific hadronisation models in Monte Carlo event generators such as PYTHIA [8] and Herwig [10] can be used.

In this work the fragmentation of partons is studied using the jet fragmentation transverse momentum, j_{T} . The j_{T} is defined as the perpendicular component of the momentum of the constituent particle with respect to reconstructed jet momentum, \vec{p}_{jet} . The length of the \vec{j}_{T} vector is

$$j_{\text{T}} = \frac{|\vec{p}_{\text{jet}} \times \vec{p}_{\text{track}}|}{|\vec{p}_{\text{jet}}|}, \quad (1)$$

where \vec{p}_{track} is the momentum of the constituent particles. It is one of many jet shape observables to study the properties of fragmenting particles with respect to the initial hard momentum during the fragmentation process. The j_{T} provides a measurement of the transverse momentum spread of the jet fragments.

Previously, j_{T} has been studied using two-particle correlations where j_{T} is calculated for particles with respect to the highest transverse momentum particle in each event instead of reconstructed jet. The study using the correlation method was done by the CCOR collaboration at ISR in pp collisions at centre-of-mass energies $\sqrt{s} = 31, 45$ and 63 GeV [15] and by the PHENIX collaboration at RHIC in pp collisions at $\sqrt{s} = 200$ GeV [16] and d–Au collisions at a center-of-mass energy per nucleon pair $\sqrt{s_{\text{NN}}} = 200$ GeV [17]. The results showed no clear dependence on the transverse momentum (p_{T}) of the trigger particle. Jet measurements to study j_{T} were done by the CDF collaboration in $p\bar{p}$ collisions at $\sqrt{s} = 1.96$ TeV [18] at Tevatron, by the ATLAS collaboration in pp at $\sqrt{s} = 7$ TeV [19] and by the LHCb collaboration in pp collisions at $\sqrt{s} = 8$ TeV [20] at the LHC. The results show a dependence of the width of j_{T} distributions with respect to the p_{T} of jets at the LHC energies.

Jets are used as an important probe for the study of the deconfined phase of strongly interacting matter, the quark–gluon plasma (QGP) that is formed in high-energy collisions of heavy nuclei. There exists plenty of experimental evidence of jet energy loss, such as the suppression of inclusive hadron spectra at high transverse momentum [21–25], the modification of back-to-back hadron-hadron [26, 27] and direct photon-hadron correlations [28], hadron–jet correlations [29, 30], the modification of reconstructed jet spectra [31, 32] and jet substructure [33–36], as compared to the expectations from elementary proton–proton collisions.

Jet quenching in heavy-ion collisions evolves multi-scale steps from hard to soft processes [37, 38]. Hard scales dominate in the elementary hard scattering. The hard scattering is followed by the subsequent branching process down to non-perturbative scales. Soft scales, of the order of the temperature of the medium, characterise interactions of soft partons produced in the shower with the QGP. Soft scales also govern hadronisation, which is expected to take place in vacuum for sufficiently energetic probes, even though some changes can persist from modifications of colour flow [39–41]. Understanding the contributions from the different processes to the jet shower evolution in medium and their scale dependence is crucial to constrain the dynamics of jet energy loss in the expanding medium [42], and fundamental medium properties like the temperature-dependent transport coefficient [43, 44]. Besides heavy-ion collisions one should study also smaller systems such as p–Pb in order to get an important baseline. Cold nuclear matter effects [45–47] in p–Pb collisions need to be considered to interpret the measurements in heavy-ion collisions.

The results for j_T distributions obtained using two-particle correlations were recently reported by the ALICE Collaboration [48] in pp and p–Pb collisions. In this paper, jet reconstruction provides a better estimate of the initial parton momentum than the leading hadron in two-particle correlations. Additionally, contrary to the correlation studies, the j_T distribution is not smeared by hadrons decaying from a short living resonance.

The j_T distributions are studied by reconstructing jets with the ALICE tracking detectors and electromagnetic calorimeter using the anti- k_T algorithm [49] with resolution parameter $R = 0.4$ in the pseudorapidity range $|\eta| < 0.25$ in pp collisions at $\sqrt{s} = 5.02$ TeV and p–Pb minimum bias collisions at $\sqrt{s_{NN}} = 5.02$ TeV. It is worth noting that there is a shift in the centre-of-mass rapidity of $\Delta y = 0.465$ in the direction of the proton beam because of the asymmetric collision system. The j_T distribution is further analysed by fitting and separating it into two distinct components that are assigned to the parton shower and the hadronisation process. The attempt to separate the two components is presented for the first time using jets in various jet transverse momentum ($p_{T,jet}$) ranges. We also compare the results with those obtained from PYTHIA (PYTHIA 8.3) and Herwig (Herwig 7.2) simulations.

2 Experimental setup and data samples

The data presented here were recorded by the ALICE detector in 2017 for pp collisions at $\sqrt{s} = 5.02$ TeV with 7.6×10^8 minimum-bias events ($\mathcal{L}_{int} = 15.7 \text{ nb}^{-1}$) and in 2013 for p–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV with 1.3×10^8 events ($\mathcal{L}_{int} = 620 \text{ nb}^{-1}$). Detailed information about the ALICE detector during LHC Run 1 and Run 2 can be found in Refs. [50, 51].

The V0 detector [52] provides the information for event triggering. The V0 detector consists of two scintillator hodoscopes that are located on each side of the interaction point along the beam direction. It covers the pseudorapidity region $-3.7 < \eta < -1.7$ (V0C) and $2.8 < \eta < 5.1$ (V0A). To select the minimum-bias trigger signals are required in both the V0A and V0C. This condition is used to reduce the contamination of data from beam-gas events using the timing difference of the signals between the V0A and V0C detectors [51].

The analysis is performed with events that have a primary vertex within $|z_{vtx}| < 10$ cm of the nominal interaction point at $z_{vtx} = 0$ along the beam direction. Charged particles are used for reconstruction of the primary vertex and jets. The charged particles are reconstructed with the Inner Tracking System (ITS) [53] and the Time Projection Chamber (TPC) [54]. These detectors are located inside a large solenoidal magnet that provides a homogeneous magnetic field of 0.5 T. Tracks within a pseudorapidity range $|\eta| < 0.9$ over the full azimuth are accepted. The ITS is made up of the Silicon Pixel Detector (SPD) in the innermost layers, the Silicon Drift Detector (SDD) in the middle layers and the Silicon Strip Detector (SSD) in the outermost layers, each consisting of two layers. The tracks are selected following the hybrid approach [55] which ensures a uniform distribution of tracks as a function of azimuthal angle (ϕ). The hybrid approach combines two different classes of tracks. The first class consists of tracks that have at least one hit in the SPD. The tracks from the second class do not have any SPD associated hit and mainly rely on the position information of the primary vertex when reconstructing the tracks. Combining the information from the ITS and TPC provides a p_T resolution ranging from 1 to 10% for charged particles from 0.15 and 100 GeV/c. For tracks without the ITS information, the momentum resolution is comparable to that of ITS+TPC tracks below transverse momentum $p_T = 10$ GeV/c, but for higher momenta the resolution reaches 20% at $p_T = 50$ GeV/c [51, 56].

The EMCAL covers an area with a range of $|\eta| < 0.7$ in pseudorapidity and 107 degrees in azimuth and is made up of 12288 towers in total. Each tower consists of 76 alternating layers of 1.44 mm lead and 77 layers of 1.76 mm scintillator material. The EMCAL is also used to provide a high-energy photon trigger for a high- $p_{T,jet}$ data sample that is complementary to the minimum bias trigger for a low $p_{T,jet}$ data sample. The EMCAL can be used to trigger on single shower deposits or energy deposits integrated over

a larger area. The latter is used for the high-energy photon trigger. The EMCAL trigger definition for p–Pb collisions in 2013 requires an energy deposit in a group of the towers of either 10 GeV for the low threshold trigger or 20 GeV for the high threshold trigger. A sample of 3×10^6 events ($\mathcal{L}_{\text{int}} = 5 \text{ nb}^{-1}$) with the EMCAL trigger provides increased statistics for $p_{T,\text{jet}} > 60 \text{ GeV}/c$ where the trigger bias disappears in the analysis [57]. The energy of the electromagnetic shower clusters is reconstructed in the EMCAL by searching for a tower with an energy deposit greater than a defined seed energy and merging all towers that share the energy cluster. To avoid double counting, when a cluster is matched with charged particles measured by the ITS and TPC, the sum of the transverse momentum of all the matched tracks are subtracted from the cluster energy.

3 Analysis method

For each collision event, jets are reconstructed with the anti- k_T algorithm [49] and resolution parameter $R = 0.4$ using FastJet [58]. The p_T -recombination scheme is used when reconstructing jets. Jets are selected in $|\eta| < 0.25$ to satisfy the fiducial acceptance of the EMCAL. The jet energy resolution $\text{JER} = \sigma(p_{T,\text{jet}}^{\text{reco}})/p_{T,\text{jet}}^{\text{true}}$ is calculated as 20% (18%) at $p_{T,\text{jet}}^{\text{true}} = 20 \text{ GeV}/c$ and 21% (19%) at 100 GeV/c in pp (p–Pb) collisions. The jet angular resolution is estimated as 29% (28%) and 2% (2%) at $p_{T,\text{jet}} = 20 \text{ GeV}/c$ 20% (19%) and 1.2% (1.2%) at $p_{T,\text{jet}} = 100 \text{ GeV}/c$ in pp (p–Pb) collisions for pseudorapidity and azimuthal angle, respectively. In the jet reconstruction both charged particles with $p_T > 0.15 \text{ GeV}/c$ and EMCAL clusters with $p_T > 0.3 \text{ GeV}/c$ are considered. All charged particles within a fixed cone with a resolution parameter R are taken as jet constituents, instead of using the list of jet constituents provided by the jet algorithm [19, 59]. Results are presented in terms of the jet transverse momentum $p_{T,\text{jet}}$.

The resulting j_T distributions are corrected for detector effects using the unfolding method in Ref. [60]. The response matrix used for the unfolding is obtained from events generated by PYTHIA 8 Monash 2013 (PYTHIA 8.2) [61] for the correction of the data sample in pp collisions and PYTHIA 6 Perugia 2011 (PYTHIA 6.4) [62] for the correction of the one in p–Pb collisions. The events are transported through the ALICE experimental set up described with GEANT 3 [63, 64]. This response matrix $(j_T^{\text{rec}}, p_{T,\text{jet}}^{\text{rec}}; j_T^{\text{true}}, p_{T,\text{jet}}^{\text{true}})$ has 2×2 dimensions to correct the detector inefficiency for jet transverse momentum ($p_{T,\text{jet}}$) and j_T simultaneously, where j_T^{true} and $p_{T,\text{jet}}^{\text{true}}$ are obtained from particle level jets by PYTHIA 6 and 8 and j_T^{rec} and $p_{T,\text{jet}}^{\text{rec}}$ are the corresponding measured values in ALICE, respectively. As a primary method the unfolding is performed with an iterative (Bayesian) algorithm as implemented in the RooUnfold package [60]. The unfolding procedure is tested by dividing the generated data sample into two halves. The first half is used to fill the response matrix. The second half is used to test the closure of the unfolding method. For $40 < p_{T,\text{jet}} < 150 \text{ GeV}/c$, the generated $p_{T,\text{jet}}$ distribution is recovered. For $j_T > 0.1 \text{ GeV}/c$, the j_T distribution is also recovered.

The effect of the underlying event background is estimated by looking at a cone perpendicular to the observed jet axis ($\frac{\pi}{2}$ rotation in ϕ , for details see Refs [65, 66]). The background j_T is calculated for any track that is found within this cone and the rotated jet axis is used as reference for j_T . The background obtained in this manner is subtracted from the unfolded inclusive j_T distribution, which gives the resulting signal distribution as shown in Eq. 2. The probability of events with jets inside the perpendicular cone are estimated as 1–2% of the total number of jets. Jets reconstructed with charged particles only (charged jet) for $R = 0.4$ and $p_{T,\text{jet}}^{\text{ch}} > 5 \text{ GeV}/c$ are used to check other jets inside the perpendicular since charged jets can cover the full azimuthal angle contrary to the case of jets in the EMCAL acceptance. To make sure there is no jet contribution in the background, those events are not used for background estimation. Because of this reason, $N_{\text{perpendicular jets}}$ is less than N_{jets} by about 1–2% in Eq. 2.

$$\frac{1}{N_{\text{jets}}} \frac{dN}{j_T, \text{ch} dj_T, \text{ch}} \Big|_{\text{signal}} = \frac{1}{N_{\text{jets}}} \frac{dN}{j_T, \text{ch} dj_T, \text{ch}} \Big|_{\text{inclusive}} - \frac{1}{N_{\text{perpendicular jets}}} \frac{dN}{j_T, \text{ch} dj_T, \text{ch}} \Big|_{\text{background}} \quad (2)$$

The resulting signal distribution is fitted with the two-component function shown in Eq. 3. A Gaussian

distribution centered at $j_T = 0$ GeV/ c is used for lower j_T and an inverse gamma function is used for j_T above 1 GeV/ c , where B_1 to B_5 are parameters [48].

$$\frac{1}{N_{\text{jets}}} \frac{dN}{j_{T,\text{ch}} dj_{T,\text{ch}}} = \frac{B_2}{B_1 \sqrt{2\pi}} e^{-\frac{j_T^2}{2B_1^2}} + \frac{B_3 B_5^{B_4}}{\Gamma(B_4)} \frac{e^{-\frac{B_5}{j_T}}}{j_T^{B_4+1}} \quad (3)$$

To achieve stable results the fitting is performed in two steps. First, lower and higher parts of the j_T distribution are fitted with a Gaussian and inverse gamma function, respectively. After getting the results from the individual fits, they are combined into a single function with initial values from the individual results and then an additional fit is performed. After getting the fit function, $\sqrt{\langle j_T^2 \rangle}$ (RMS) and yield values are extracted separately from each component. The narrow component RMS from the Gaussian part is determined as

$$\sqrt{\langle j_T^2 \rangle} = \sqrt{2} B_1 \quad (4)$$

and the wide component RMS value from the inverse gamma function is calculated as

$$\sqrt{\langle j_T^2 \rangle} = \frac{B_5}{\sqrt{(B_4 - 2)(B_4 - 3)}} \quad , \quad (5)$$

where it is required that $B_4 > 3$.

4 Systematic uncertainties

The systematic uncertainties in this analysis come from the background estimation, the unfolding procedure and the uncertainties related to track and cluster selection. The effect originating from uncertainty in the tracking efficiency is estimated with a PYTHIA simulation by removing 4% of tracks randomly from each event corresponding to a mismatching probability of tracks between the ITS and TPC. The resulting variations in the RMS values are less than 4% and 5% for the wide and narrow components, respectively. The uncertainty related to the EMCal energy scale was estimated by scaling cluster energies up and down by 2% in the PYTHIA particle level generation in order to reflect a non-linearity correction of the EMCal energy scale ranging from about 7% at 0.5 GeV/ c to a negligible value above 3 GeV/ c . Similarly, the jet momentum was scaled by $\pm 2\%$ when determining $p_{T,\text{jet}}$ to check how the cluster energy affects j_T distributions. The variation of both RMS components is seen to be less than 2%.

The systematic uncertainty on the background estimation was studied using the ‘‘random background’’ method as an alternative to that of the perpendicular-cone. This method assigns new random η and ϕ of the existing tracks in the event using a uniform distribution without changing their p_T values. A random jet cone is also from uniform η and ϕ distributions covering $|\eta| < 0.25$ and $0 < \phi < 2\pi$ and tracks near the jet axis are not used. The resulting uncertainty is below 5% for the wide component RMS and below 9% for the narrow component RMS in p–Pb collisions. To study the effect of background fluctuations in p–Pb collisions, a study based on embedding particles generated with PYTHIA in real events was performed. The embedded particles are simulated by following the multiplicity density information [67] and p_T distribution [68] of charged particles in p–Pb collisions in ALICE. The effect in RMS is negligible for both RMS components.

The systematic uncertainty introduced by the unfolding procedure was determined by repeating the unfolding using the Singular-Value Decomposition (SVD) method as an alternative [69]. Given that the SVD method does not allow for multi-dimensional unfolding, the unfolding is performed separately for different $p_{T,\text{jet}}$ intervals. In a PYTHIA closure test, the true distribution for $j_T > 0.1$ GeV/ c was in general found to be between the unfolded distributions from the iterative and SVD methods within 2%. The difference between the methods when unfolding data is used as an estimate of the unfolding uncertainty.

The iterative unfolding algorithm permits the change of the number of iterations as a regularisation parameter. The stability of the results was verified by using one iteration above and below instead of the default value, where the default value is chosen by checking that unfolded j_T distributions converge. Also, the regularisation parameter k is varied by one unit above and below with respect to the default solution of the SVD method that is determined by following the guideline [69]. The iterative algorithm requires a prior estimate of the shape of the distribution. As a default prior, generated PYTHIA distribution is used. To estimate the effect of the prior, the unfolded j_T distribution is used as a prior instead. The effect of the unfolding for different ranges of $p_{T,\text{jet}}$ is tested by varying the first value of $p_{T,\text{jet}}$ from 5 to 15 GeV/c. These effects are found negligible compared to that for the two different unfolding methods. The resulting uncertainty by the unfolding procedure is below 8% for both wide and narrow component RMS in p–Pb collisions. In pp collisions it is 9% and 12% for the wide and narrow components, respectively.

The model dependence of the unfolding procedure was explored by weighting the response matrix with PYTHIA. The jet yield in the response matrix is varied by $\pm 30\%$ for the angularity $g > 0.1$. The angularity is defined as $g = \sum_i (p_{T,i} \times r_i) / p_{T,\text{jet}}$, where $p_{T,i}$ is the p_T of the i^{th} constituent of the jet and $r_i = \sqrt{\Delta\eta_i^2 + \Delta\phi_i^2}$ is the distance of the i^{th} constituent from the jet axis [32, 70]. The effect is found to be below 2% for the wide component and negligible for the narrow component.

The different sources of systematic uncertainty are considered as uncorrelated and the values are summed in quadrature. The summary table in Table 1 shows an overview of systematic uncertainties for $40 < p_{T,\text{jet}} < 60$ GeV/c in pp and p–Pb collisions.

Table 1: Summary of systematic uncertainties for $40 < p_{T,\text{jet}} < 60$ GeV/c in pp and p–Pb collisions.

source	j_T distribution at $j_T = 0.2\text{--}0.8\text{--}2$ GeV/c		Wide RMS		Narrow RMS	
	pp	p–Pb	pp	p–Pb	pp	p–Pb
Background	2–2–5%	neg.–2–5%	1.1%	5%	2.9%	9%
Unfolding	10–neg.–20%	10–neg.–12%	9%	8%	12%	8%
Tracking	2–2–2%	2–1–neg.%	0.4%	4%	0.2%	5%
EMCal	2–2–5%	2–2–2%	1.8%	1%	0.2%	1%
Model dependence	neg.–2–5%	neg.–neg.–10%	0.5%	2%	neg.	neg.
Total	11–4–22%	10–3–16%	9%	10%	12%	13%

5 Results

The j_T distribution in pp collisions at $\sqrt{s} = 5.02$ TeV is compared with that in p–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV in Fig. 1 for jet transverse momentum in $40 < p_{T,\text{jet}} < 60$ GeV/c. The ratio of the j_T distributions represents the consistence of the result in pp and p–Pb collisions and implies no clear cold nuclear matter effects in p–Pb collisions. For the interval in $100 < p_{T,\text{jet}} < 150$ GeV/c, the comparison is not provided because of the lack of enough statistics in minimum-bias pp collisions and the absence of the data sample with the EMCal trigger in the corresponding pp data taking period.

Figure 2 shows the distributions of j_T for charged particles in different $p_{T,\text{jet}}$ intervals after applying the unfolding correction and background subtraction in p–Pb collisions at $\sqrt{s_{\text{NN}}} = 5.02$ TeV. The yield at low j_T stays constant with increasing $p_{T,\text{jet}}$. At high j_T the yield increases and the distributions become wider with increasing $p_{T,\text{jet}}$ as indicated by the ratios of the j_T distributions shown in the bottom panel. Notably, this is due to kinematical limits. At midrapidity, within a fixed cone the maximum j_T depends on the track momentum by the relation of $j_{T,\text{max}} \approx R \times p_{T,\text{track}}$, resulting in an increase of the possible j_T as $p_{T,\text{jet}}$ increases. Though jets with larger momenta are more collimated, the net effect is an increase of $\langle j_T \rangle$ as $p_{T,\text{jet}}$ increases. These measurements are consistent with the findings by the ATLAS [19] and LHCb collaborations [20].

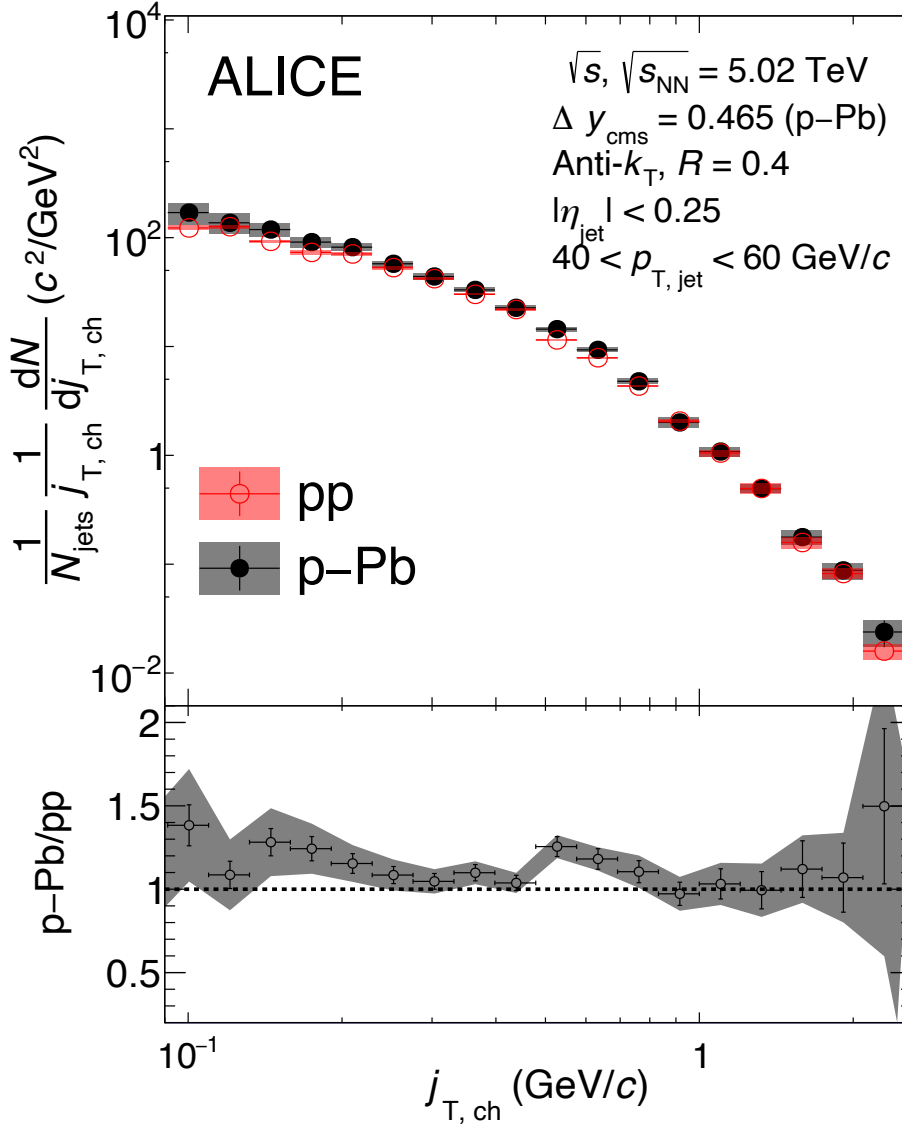


Figure 1: Comparison of the j_T distributions in pp and p–Pb collisions at $\sqrt{s}, \sqrt{s_{NN}} = 5.02$ TeV in $40 < p_{T,\text{jet}} < 60$ GeV/c. The centre-of-mass rapidity in p–Pb collisions is shifted by $\Delta y = 0.465$ in the direction of the proton beam.

Figure 3 shows the j_T distribution in p–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV for jets with $60 < p_{T,\text{jet}} < 80$ GeV/c compared with expectations from various generators in pp collisions at $\sqrt{s} = 5.02$ TeV. PYTHIA 8 based models (PYTHIA 8.3) and Herwig (Herwig 7.2) handle both the showering process and hadronisation differently. PYTHIA 8 uses the Lund string model [71] to perform the hadronisation stage. Herwig uses a cluster model for the hadronisation [9, 10]. PYTHIA 8 has p_T -ordered showers by default while Herwig implements a parton shower using the coherent branching algorithm [72], which has angular ordering as a central feature. The p_T -ordering in a PYTHIA 8 shower is a compromise [73]: ordering in the p_T at splitting ensures the ordering in the hardness and also effectively favours large angles. Herwig

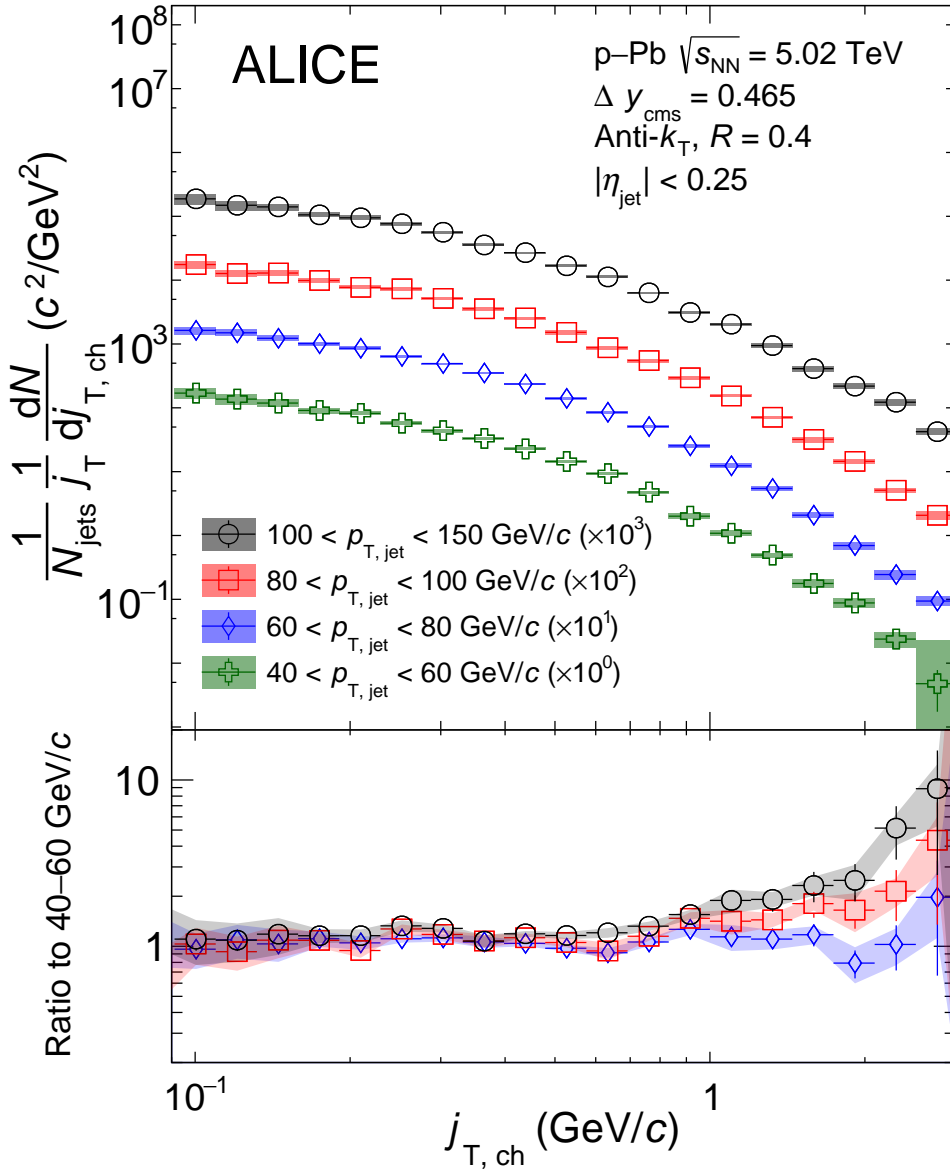


Figure 2: The j_T distributions of charged particles in $R = 0.4$ anti- k_T jets as measured in p–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV for different ranges of $p_{T,jet}$. The centre-of-mass rapidity is shifted by $\Delta y = 0.465$ in the direction of the proton beam. The bottom panel shows ratios of the j_T distributions with respect to that in $40 < p_{T,jet} < 60$ GeV/c.

describes the j_T distribution better than other models for the whole j_T region. Other PYTHIA 8 based models describe the data at high j_T but not in the low j_T region. The results for the other $p_{T,jet}$ intervals are reported in Figs. B1, B2 and B3 that derive the same conclusion. Models describe the data better as $p_{T,jet}$ increases in pp collisions. This is also true at higher j_T , however, models underestimate the data at lower j_T consistently for all $p_{T,jet}$ ranges in p–Pb collisions.

PYTHIA 8 Monash 2013 [61] adopted LHC data to constrain the initial-state radiation and multi-parton interaction parameters based on the default parameters of PYTHIA 8 tune 4C [74]. There is no clear separation of the j_T distributions originating from the different tunes of PYTHIA 8. As of version 8.3 PYTHIA 8 implemented two more shower models as part of the code. Those are VINCIA and Dire Showers that are based on the k_T (transverse momentum of a dipole)-ordered picture of QCD splitting [75, 76]. The j_T distributions generated by the two shower models were obtained by using the default parameters

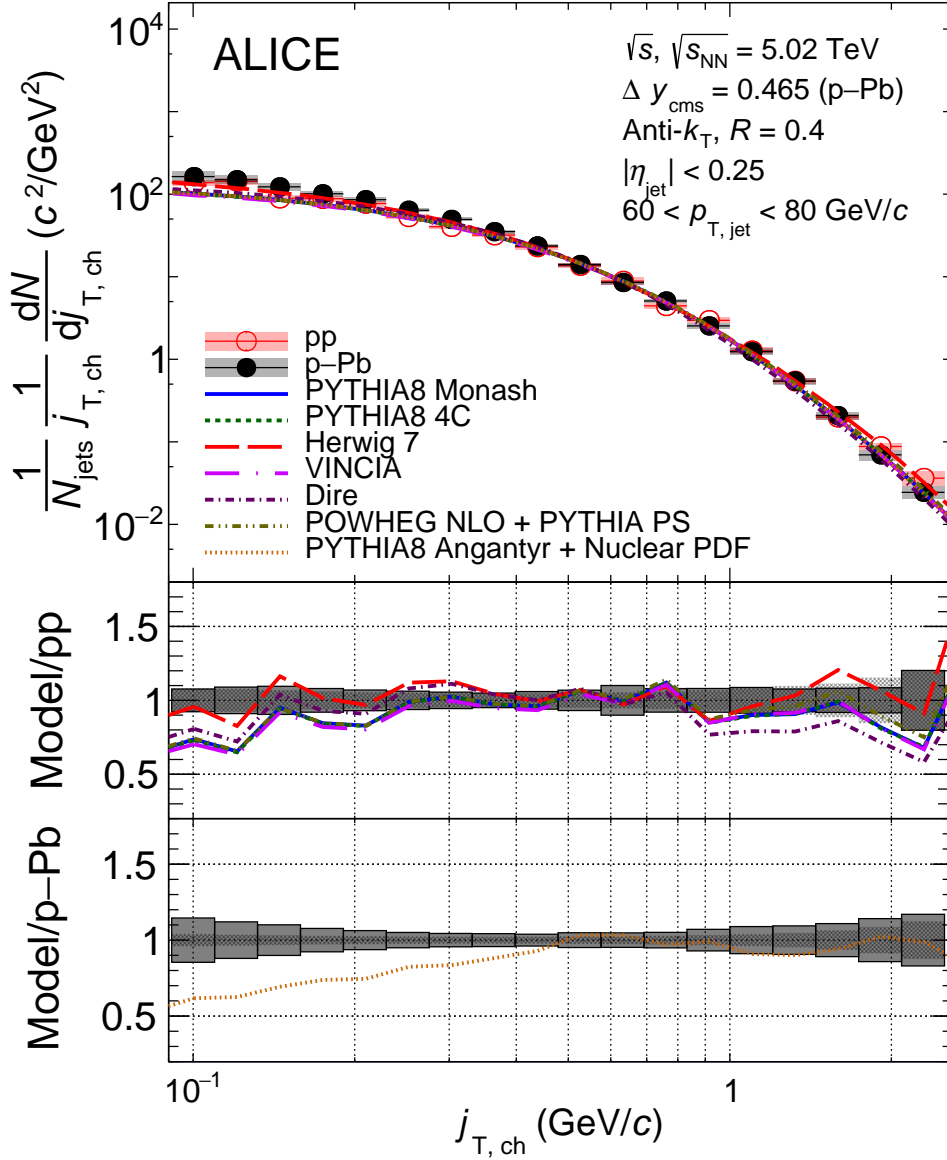


Figure 3: The j_T distribution in p–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV for jets with transverse momentum in $60 < p_{T, \text{jet}} < 80$ GeV/c. The measured data are compared to calculations by theoretical models in pp collisions at $\sqrt{s} = 5.02$ TeV.

of PYTHIA 8 tune 4C. In order to study the effect of the NLO calculation accuracy for the parton showering in PYTHIA 8 (POWHEG NLO + PYTHIA PS), the j_T distribution generated with the combined POWHEG [77] and PYTHIA simulation is also compared to the data. The j_T distributions obtained with the POWHEG NLO calculation and Dire Shower display themselves as upper and lower bounds of the PYTHIA 8 based models for the higher j_T region; however, they are within the systematic uncertainty of the data for the higher j_T region. PYTHIA 8 Angantyr extends pp simulation of PYTHIA 8 to the case of heavy-ion collisions [78]. PYTHIA 8 Angantyr is used to simulate p–Pb collisions with the nuclear parton distribution function (PDF) EPS09LO [47] for the Pb-ion beam. The resulting j_T distribution is almost the same with those by pp simulations with a proton PDF and it does not describe the data for the lower j_T region at all.

The distributions are fitted with the two-component fit motivated by [48]. The function forms are given

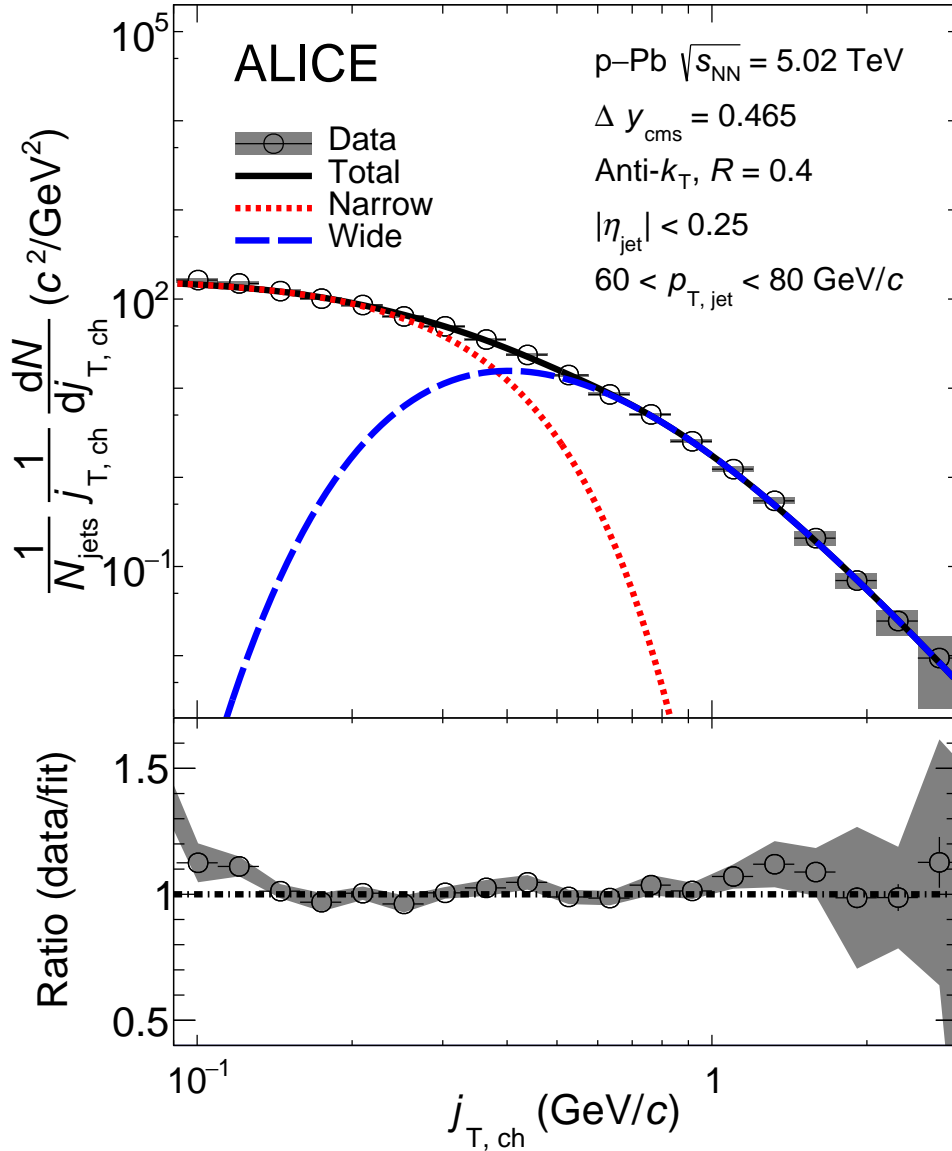


Figure 4: The j_T distribution of charged particles with a two-component fit for $60 < p_{T,\text{jet}} < 80$ GeV/ c . The distribution is fitted with the two-component fit described in Sec. 3.

in Eq. 3. An example of the fitted distribution is shown in Fig. 4 for $60 < p_{T,\text{jet}} < 80$ GeV/ c . The Gaussian term corresponds to the narrow part that can be associated with the hadronisation process, while the inverse gamma corresponds to the wide component characterising the QCD shower. The j_T distributions are described well by the two-component model fit. The corresponding statistical uncertainties are calculated via the general error propagation formulas in Eq. 6

$$\delta \sqrt{\langle j_T^2 \rangle} = \sqrt{2} \delta B_1 \quad \text{and} \quad \sqrt{\left(\frac{(5-2B_4)B_5 \delta B_4}{(2(B_4-2)(B_4-3))^{\frac{3}{2}}} \right)^2 + \left(\frac{\delta B_5}{\sqrt{(B_4-2)(B_4-3)}} \right)^2} \quad (6)$$

for the narrow and wide component RMS values, respectively.

The widths of the j_T distributions are determined as a function of the transverse momentum of jet. The RMS ($\sqrt{\langle j_T^2 \rangle}$) values for the two components are shown in Fig. 5 along with comparisons to Monte

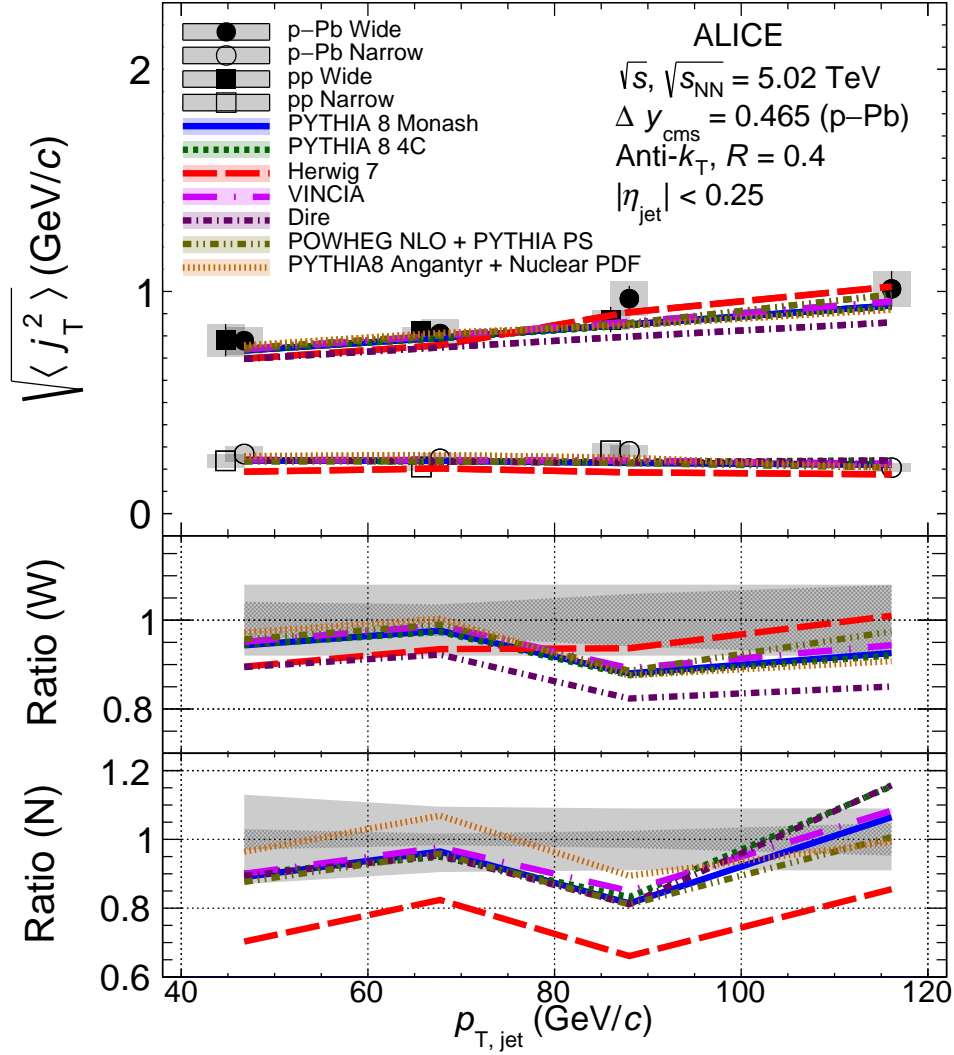


Figure 5: RMS values extracted from the fits for the Gaussian (narrow) and inverse gamma (wide) components. The middle and bottom plots show ratios of models to data for the wide and narrow components, respectively. The grey filled bands with (without) a hatched line in the ratio plots represent the statistical (systematic) uncertainties of the p–Pb data. Note that pp data points are shifted by -2 GeV/c on the horizontal axis to be distinguished from p–Pb data points.

Carlo simulations. There is clear separation in the width of the wide and narrow components of the j_T distributions. The RMS values of the wide component are 3-4 times larger than the narrow component RMS. The wide component RMS shows an increasing trend with increasing $p_{T,jet}$ that is parameterised by a linear function as $\sqrt{\langle j_T^2 \rangle} = 0.005 (\pm 0.004) \times p_{T,jet} + 0.497 (\pm 0.255)$, while the narrow component RMS stays constant with the fitted value of $0.253 (\pm 0.009)$. Both of these trends are qualitatively consistent with the results in the dihadron j_T analysis [48].

All models except for Herwig describe the RMS values relatively well for the narrow RMS component. For the wide RMS component Herwig describes the data best as $p_{T,jet}$ increases. Dire Shower shows clearly lower values compared to data up to 18% for the wide RMS components. Other PYTHIA 8 based models show a good description for the lower j_T region, however, they underestimate the data for the higher $p_{T,jet}$ region.

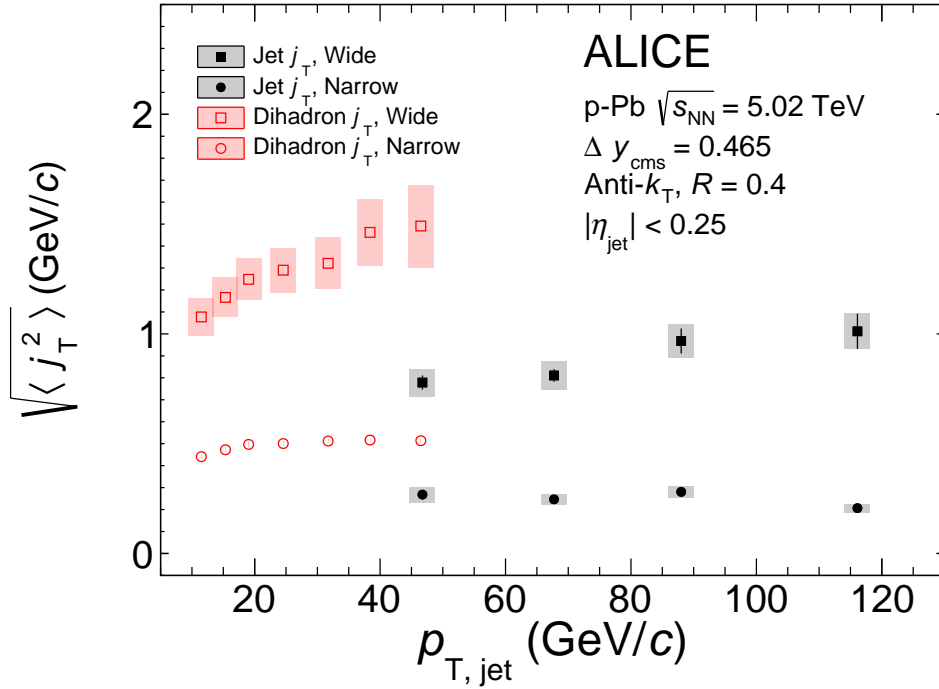


Figure 6: Comparison of results from the jet-based and dihadron-based j_T analyses [48]. Ranges of dihadron trigger p_T ($p_{T, \text{trigger}}$) are converted to corresponding $p_{T, \text{jet}}$ ranges using observed mean $p_{T, \text{jet}}$ values in $p_{T, \text{trigger}}$ bins. Dihadron results are shown for $0.2 < x_{\parallel} < 0.4$, where x_{\parallel} is the longitudinal component fraction of the associated track momentum with respect to the momentum of the trigger track. The difference of the two analyses originates from the different kinematic selections and the choice of the axis used for the j_T calculation. See text for more details.

6 Discussion

The comparison with the results from the dihadron analysis [48] performed for the same collision system and energy is shown in Fig. 6. Different p_T regions of leading particles used in the dihadron analysis are converted to the corresponding average momentum of the jets which contain those leading particles. The wide and narrow components of the dihadron results are for $0.2 < x_{\parallel} < 0.4$, where x_{\parallel} is the projection of the momentum of the associated track to that of the trigger particles. Wide component RMS values tend to increase with increasing $p_{T, \text{trigger}}$ and $p_{T, \text{jet}}$, whereas narrow component RMS values of both results show a weak dependence on $p_{T, \text{jet}}$ above 20 GeV/c. The trends are similar for dihadron and jet j_T results. However, the RMS values of the dihadron analysis are larger than those for the jet analysis both for the narrow and wide components.

The difference in the narrow and wide RMS components can be explained by the following two factors. The first one is due to the different kinematic selections on the charged particles in the same jet from which the j_T values are calculated. The other one is due to the choice of the axis used for the j_T calculation. In the dihadron analysis j_T is calculated for all near-side tracks if the associated tracks satisfy the condition $\vec{p}_{\text{leading}} \times \vec{p}_a > 0$. Here \vec{p}_{leading} and \vec{p}_a are the momentum vectors of the leading and associated tracks, respectively. Thus, the kinematical limit $j_{T, \text{max}}$ can be larger in the dihadron analysis than in the jet analysis in which only particles in a cone with $R = 0.4$ are considered.

The effect of the R parameter choice and $p_{T, \text{jet}}$ dependence on j_T was studied using PYTHIA 8 and the results are shown in Fig. 7a. The usage of a fixed cone sets stringent limits on the possible j_T values. Increasing the cone size loosens these limits and allows for higher j_T values. The effect on the wide and narrow components of the j_T distributions for PYTHIA 8 is shown in Fig. 7b, where the wide component RMS gets larger by about 10% when going from $R = 0.3$ to 0.4 and from 0.4 to 0.5, indicating that the

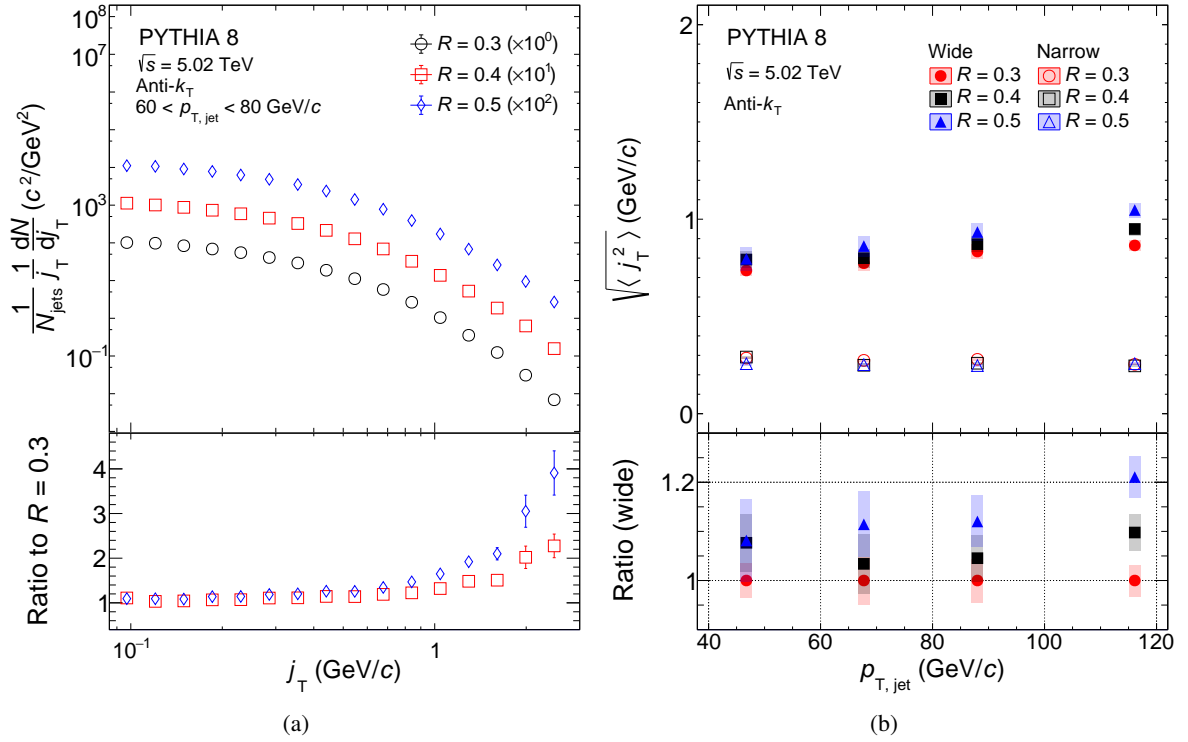


Figure 7: The effect of changing the R parameter in jet finding on j_T distributions obtained with PYTHIA 8 simulations. Comparison of (a) j_T signal distributions for different R parameters and their ratios to that of $R = 0.3$ and (b) RMS values of the wide and narrow components and their ratios to that of $R = 0.3$ for the wide component only.

kinematic limit introduced by increasing R results in a widening of the j_T distribution. For the narrow component the effect is relatively small and they appear independent of the R parameter and $p_{T,\text{jet}}$. There can also be a broadening effect for jets caused by the increasing gluon jet fraction as the kinematical limit increases [70]. Additionally, there is an effect originating from the kinematic cut on x_{\parallel} values in the dihadron analysis that can alter the j_T distributions – but that is not further investigated here.

It is worth noting that the leading-track momentum vector provides an imperfect estimate of the jet axis. Because the leading track in general is at an angle compared to the jet axis, the resulting j_T values based on the leading track are biased from the axis of the jet. Practically, the jet axis found by the jet finding algorithm tends to minimise the j_T of jet constituents. Moreover, in the dihadron correlation analysis the usage of the leading hadron as the trigger particle imposes a trigger bias favouring quark jets resulting in jet narrowing. The impact of the different axes adopted in the two analyses is investigated by measuring j_T with respect to the leading track momentum (leading track reference), instead of the jet axis (jet-axis reference) within the same jet for $R = 0.4$. The results are shown in Fig. 8. The widths of the j_T distributions for the jet-axis reference overall are smaller than those of the leading track reference. The bias of the choice of axis becomes small as $p_{T,\text{jet}}$ increases. As shown in the bottom panels, the ratios of the distributions increase monotonically, implying that the leading track reference makes both the wide and narrow components wider as the ratio distributions show a monotonic increase.

Dihadron j_T distributions [48] are compared to those of jet j_T . Although a direct comparison between jet and dihadron j_T measurements is not possible because of the effects of the different kinematic selection and choice of the axis, RMS values of the wide and narrow components can be quantitatively understood by considering the good agreement between PYTHIA and data.

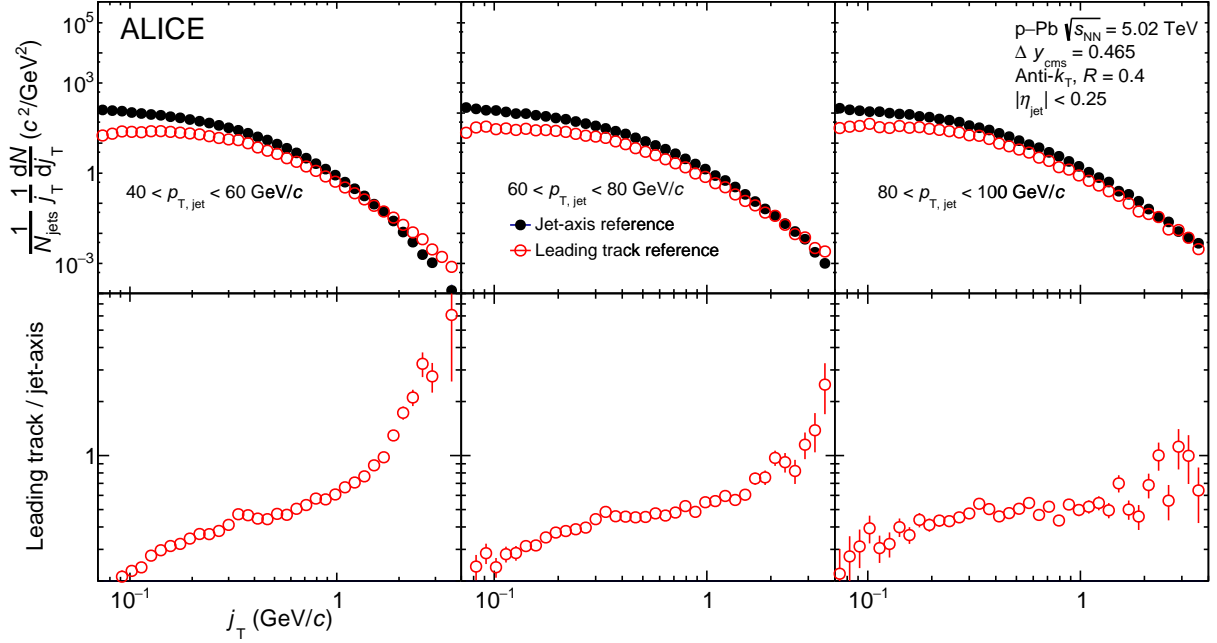


Figure 8: The j_T distributions with respect to the leading track momentum (leading track reference) and the jet axis (jet-axis reference) within the same jet for three different $p_{T,\text{jet}}$ intervals with $R = 0.4$.

7 Conclusion

In this work the jet fragmentation transverse momentum (j_T) distribution of charged particles $\frac{1}{N_{\text{jets}}} \frac{dN}{J_T dJ_T}$ is studied using jet reconstruction in pp and p–Pb collisions at \sqrt{s} , $\sqrt{s_{\text{NN}}} = 5.02$ TeV. The j_T distributions of charged particles in p–Pb collisions become wider as the jet transverse momentum $p_{T,\text{jet}}$ increases. This is understood as an effect of the reduction of the kinematical limit with increasing $p_{T,\text{jet}}$, allowing for higher j_T values. The j_T distribution in p–Pb collision is compared with that in pp collisions for jet transverse momentum in $40 < p_{T,\text{jet}} < 100$ GeV/c, which shows no clear modification of the j_T distribution for the p–Pb collision system. No significant cold nuclear matter effects are observed in the previous and current j_T measurements using dihadron correlations [48] and jet reconstruction. For the jet study, higher statistics in pp collisions for both minimum bias and EMCal trigger is demanded to interpret the effect in lower j_T and higher $p_{T,\text{jet}}$. The j_T distributions in p–Pb collisions are compared with various parton shower and fragmentation models. All models describe the data well for the higher j_T region, while they underestimate the data by about 20% and 40% at lower j_T in pp and p–Pb collisions, respectively.

Two distinct components of the jet fragmentation transverse momentum j_T are extracted for narrow and wide contributions to quantify the j_T distribution further in pp and p–Pb collisions. The width of the narrow component has only a weak dependence on jet transverse momentum, while that of the wide component increases with increasing jet transverse momentum. The results are qualitatively consistent as a function of $p_{T,\text{jet}}$ with the previous j_T study performed with dihadron correlations [48]. We also present a comparison to PYTHIA 8 (PYTHIA 8.3) and Herwig (Herwig 7.2) simulations to figure out if the two distinct components are described well by models or differences are present. For the wide component, Herwig and PYTHIA 8 based models slightly underestimate the data for the higher jet transverse momentum region. For the narrow component, the measured trends are successfully described by all models except for Herwig. This is opposite to the case of the j_T distributions at lower j_T where the narrow component corresponds. This indicates that the shape of the j_T distribution in models is also important to describe the data.

In addition to the result in p–Pb collisions, a high statistics in pp collisions will further constrain predic-

tions in model calculations for jet fragmentation and hadronisation. Future studies of the j_T distribution performed differentially in the longitudinal momentum fraction z can be used to constrain transverse-momentum dependent fragmentation functions [12].

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References

- [1] D. J. Gross and F. Wilczek, “Ultraviolet behavior of non-abelian gauge theories”, *Physical Review Letters* **30** no. 26, (1973) 1343–1346.
- [2] H. D. Politzer, “Reliable perturbative results for strong interactions?”, *Physical Review Letters* **30** no. 26, (1973) 1346–1349.
- [3] D. Gross and F. Wilczek, “Asymptotically Free Gauge Theories - I”, *Phys. Rev. D* **8** (1973) 3633–3652.
- [4] D. J. Gross and F. Wilczek, “Asymptotically Free Gauge Theories - II”, *Phys. Rev. D* **9** (1974) 980–993.
- [5] H. Georgi and H. D. Politzer, “Electroproduction scaling in an asymptotically free theory of strong interactions”, *Physical Review D* **9** no. 2, (1974) 416.
- [6] A. Buckley *et al.*, “General-purpose event generators for LHC physics”, *Phys. Rept.* **504** (2011) 145–233, arXiv:1101.2599 [hep-ph].
- [7] Y. L. Dokshitzer, V. A. Khoze, A. H. Mueller, and S. Troian, *Basics of perturbative QCD*. 1991.
- [8] T. Sjöstrand, S. Mrenna, and P. Z. Skands, “A Brief Introduction to PYTHIA 8.1”, *Comput. Phys. Commun.* **178** (2008) 852–867, arXiv:0710.3820 [hep-ph].
- [9] M. Bahr *et al.*, “Herwig++ Physics and Manual”, *Eur. Phys. J.* **C58** (2008) 639–707, arXiv:0803.0883 [hep-ph].
- [10] J. Bellm *et al.*, “Herwig 7.0/Herwig++ 3.0 release note”, *Eur. Phys. J.* **C76** no. 4, (2016) 196, arXiv:1512.01178 [hep-ph].
- [11] R. Perez-Ramos, F. Arleo, and B. Machet, “Next-to-MLLA corrections to single inclusive k_{\perp} distributions and two-particle correlations in a jet”, *Phys. Rev. D* **78** (2008) 014019, arXiv:0712.2212 [hep-ph].
- [12] Z.-B. Kang, X. Liu, F. Ringer, and H. Xing, “The transverse momentum distribution of hadrons within jets”, *JHEP* **11** (2017) 068, arXiv:1705.08443 [hep-ph].
- [13] Y. I. Azimov, Y. L. Dokshitzer, V. A. Khoze, and S. I. Troyan, “Similarity of Parton and Hadron Spectra in QCD Jets”, *Z. Phys.* **C27** (1985) 65–72.
- [14] D. Gutierrez-Reyes, I. Scimemi, W. J. Waalewijn, and L. Zoppi, “Transverse momentum dependent distributions with jets”, *Phys. Rev. Lett.* **121** no. 16, (2018) 162001, arXiv:1807.07573 [hep-ph].
- [15] **CCOR** Collaboration, A. Angelis *et al.*, “A Measurement of the Transverse Momenta of Partons, and of Jet Fragmentation as a Function of \sqrt{s} in pp Collisions”, *Phys. Lett.* **B97** (1980) 163–168.

- [16] **PHENIX** Collaboration, S. S. Adler *et al.*, “Jet properties from dihadron correlations in p+p collisions at $\sqrt{s} = 200$ GeV”, *Phys. Rev.* **D74** (2006) 072002, arXiv:hep-ex/0605039 [hep-ex].
- [17] **PHENIX** Collaboration, S. S. Adler *et al.*, “Jet structure from dihadron correlations in d+Au collisions at $\sqrt{s(NN)^{1/2}} = 200$ -GeV”, *Phys. Rev.* **C73** (2006) 054903, arXiv:nucl-ex/0510021 [nucl-ex].
- [18] **CDF** Collaboration, T. Aaltonen *et al.*, “Measurement of the k_T Distribution of Particles in Jets Produced in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ -TeV”, *Phys. Rev. Lett.* **102** (2009) 232002, arXiv:0811.2820 [hep-ex].
- [19] **ATLAS** Collaboration, G. Aad *et al.*, “Measurement of the jet fragmentation function and transverse profile in proton-proton collisions at a center-of-mass energy of 7 TeV with the ATLAS detector”, *Eur. Phys. J.* **C71** (2011) 1795, arXiv:1109.5816 [hep-ex].
- [20] **LHCb** Collaboration, R. Aaij *et al.*, “Measurement of charged hadron production in Z-tagged jets in proton-proton collisions at $\sqrt{s} = 8$ TeV”, *Phys. Rev. Lett.* **123** no. 23, (2019) 232001, arXiv:1904.08878 [hep-ex].
- [21] **PHENIX** Collaboration, K. Adcox *et al.*, “Suppression of hadrons with large transverse momentum in central Au+Au collisions at $\sqrt{s_{NN}} = 130$ -GeV”, *Phys. Rev. Lett.* **88** (2002) 022301, arXiv:nucl-ex/0109003 [nucl-ex].
- [22] **STAR** Collaboration, J. Adams *et al.*, “Evidence from d + Au measurements for final state suppression of high p(T) hadrons in Au+Au collisions at RHIC”, *Phys. Rev. Lett.* **91** (2003) 072304, arXiv:nucl-ex/0306024 [nucl-ex].
- [23] **BRAHMS** Collaboration, I. Arsene *et al.*, “Transverse momentum spectra in Au+Au and d+Au collisions at $\sqrt{s_{NN}} = 200$ GeV and the pseudorapidity dependence of high p_T suppression”, *Phys. Rev. Lett.* **91** (2003) 072305, arXiv:nucl-ex/0307003 [nucl-ex].
- [24] **CMS** Collaboration, V. Khachatryan *et al.*, “Charged-particle nuclear modification factors in PbPb and pPb collisions at $\sqrt{s_{NN}} = 5.02$ TeV”, *JHEP* **04** (2017) 039, arXiv:1611.01664 [nucl-ex].
- [25] **ALICE** Collaboration, S. Acharya *et al.*, “Transverse momentum spectra and nuclear modification factors of charged particles in pp, p–Pb and Pb–Pb collisions at the LHC”, *JHEP* **11** (2018) 013, arXiv:1802.09145 [nucl-ex].
- [26] **PHENIX** Collaboration, A. Adare *et al.*, “Transverse momentum and centrality dependence of dihadron correlations in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV: Jet-quenching and the response of partonic matter”, *Phys. Rev.* **C77** (2008) 011901, arXiv:0705.3238 [nucl-ex].
- [27] **ALICE** Collaboration, K. Aamodt *et al.*, “Particle-yield modification in jet-like azimuthal di-hadron correlations in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *Phys. Rev. Lett.* **108** (2012) 092301, arXiv:1110.0121 [nucl-ex].
- [28] **PHENIX** Collaboration, A. Adare *et al.*, “Medium modification of jet fragmentation in Au + Au collisions at $\sqrt{s_{NN}} = 200$ GeV measured in direct photon-hadron correlations”, *Phys. Rev. Lett.* **111** no. 3, (2013) 032301, arXiv:1212.3323 [nucl-ex].
- [29] **ALICE** Collaboration, J. Adam *et al.*, “Measurement of jet quenching with semi-inclusive hadron-jet distributions in central Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *JHEP* **09** (2015) 170, arXiv:1506.03984 [nucl-ex].

- [30] **STAR** Collaboration, L. Adamczyk *et al.*, “Measurements of jet quenching with semi-inclusive hadron+jet distributions in Au+Au collisions at $\sqrt{s_{NN}} = 200$ GeV”, *Phys. Rev. C* **96** no. 2, (2017) 024905, arXiv:1702.01108 [nucl-ex].
- [31] **ALICE** Collaboration, J. Adam *et al.*, “Measurement of jet suppression in central Pb–Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *Phys. Lett.* **B746** (2015) 1–14, arXiv:1502.01689 [nucl-ex].
- [32] **ALICE** Collaboration, S. Acharya *et al.*, “Measurements of inclusive jet spectra in pp and central Pb–Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV”, arXiv:1909.09718 [nucl-ex].
- [33] **CMS** Collaboration, A. M. Sirunyan *et al.*, “Observation of Medium-Induced Modifications of Jet Fragmentation in Pb-Pb Collisions at $\sqrt{s_{NN}} = 5.02$ TeV Using Isolated Photon-Tagged Jets”, *Phys. Rev. Lett.* **121** no. 24, (2018) 242301, arXiv:1801.04895 [hep-ex].
- [34] **CMS** Collaboration, S. Chatrchyan *et al.*, “Measurement of jet fragmentation in PbPb and pp collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *Phys. Rev.* **C90** no. 2, (2014) 024908, arXiv:1406.0932 [nucl-ex].
- [35] **ALICE** Collaboration, S. Acharya *et al.*, “Medium modification of the shape of small-radius jets in central Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *JHEP* **10** (2018) 139, arXiv:1807.06854 [nucl-ex].
- [36] **ALICE** Collaboration, S. Acharya *et al.*, “Exploration of jet substructure using iterative declustering in pp and Pb–Pb collisions at LHC energies”, *Phys. Lett.* **B802** (2020) 135227, arXiv:1905.02512 [nucl-ex].
- [37] A. Kurkela and U. A. Wiedemann, “Picturing perturbative parton cascades in QCD matter”, *Phys. Lett.* **B740** (2015) 172–178, arXiv:1407.0293 [hep-ph].
- [38] **JETSCAPE** Collaboration, Y. Tachibana *et al.*, “Jet substructure modifications in a QGP from multi-scale description of jet evolution with JETSCAPE”, *PoS HardProbes2018* (2018) 099, arXiv:1812.06366 [nucl-th].
- [39] P. Aurenche and B. G. Zakharov, “Jet color chemistry and anomalous baryon production in AA-collisions”, *Eur. Phys. J.* **C71** (2011) 1829, arXiv:1109.6819 [hep-ph].
- [40] A. Beraudo, J. G. Milhano, and U. A. Wiedemann, “Medium-induced color flow softens hadronization”, *Phys. Rev.* **C85** (2012) 031901, arXiv:1109.5025 [hep-ph].
- [41] A. Beraudo, J. G. Milhano, and U. A. Wiedemann, “The Contribution of Medium-Modified Color Flow to Jet Quenching”, *JHEP* **07** (2012) 144, arXiv:1204.4342 [hep-ph].
- [42] J. Casalderrey-Solana, Y. Mehtar-Tani, C. A. Salgado, and K. Tywoniuk, “New picture of jet quenching dictated by color coherence”, *Phys. Lett.* **B725** (2013) 357–360, arXiv:1210.7765 [hep-ph].
- [43] F. D’Eramo, M. Lekaveckas, H. Liu, and K. Rajagopal, “Momentum Broadening in Weakly Coupled Quark-Gluon Plasma (with a view to finding the quasiparticles within liquid quark-gluon plasma)”, *JHEP* **05** (2013) 031, arXiv:1211.1922 [hep-ph].
- [44] A. Ayala, I. Dominguez, J. Jalilian-Marian, and M. E. Tejeda-Yeomans, “Relating \hat{q} , η/s and ΔE in an expanding quark-gluon plasma”, *Phys. Rev.* **C94** no. 2, (2016) 024913, arXiv:1603.09296 [hep-ph].

- [45] R. Baier, Y. L. Dokshitzer, A. H. Mueller, S. Peigne, and D. Schiff, “Radiative energy loss and $p(T)$ broadening of high-energy partons in nuclei”, *Nucl. Phys.* **B484** (1997) 265–282, arXiv:hep-ph/9608322 [hep-ph].
- [46] L. D. McLerran and R. Venugopalan, “Computing quark and gluon distribution functions for very large nuclei”, *Phys. Rev.* **D49** (1994) 2233–2241, arXiv:hep-ph/9309289 [hep-ph].
- [47] K. J. Eskola, H. Paukkunen, and C. A. Salgado, “EPS09: A New Generation of NLO and LO Nuclear Parton Distribution Functions”, *JHEP* **04** (2009) 065, arXiv:0902.4154 [hep-ph].
- [48] ALICE Collaboration, S. Acharya *et al.*, “Jet fragmentation transverse momentum measurements from di-hadron correlations in $\sqrt{s} = 7$ TeV pp and $\sqrt{s_{NN}} = 5.02$ TeV p-Pb collisions”, *JHEP* **03** (2019) 169, arXiv:1811.09742 [nucl-ex].
- [49] M. Cacciari, G. P. Salam, and G. Soyez, “The anti- k_t jet clustering algorithm”, *JHEP* **04** (2008) 063, arXiv:0802.1189 [hep-ph].
- [50] ALICE Collaboration, K. Aamodt *et al.*, “The ALICE experiment at the CERN LHC”, *JINST* **3** (2008) S08002.
- [51] ALICE Collaboration, B. Abelev *et al.*, “Performance of the ALICE Experiment at the CERN LHC”, *Int. J. Mod. Phys.* **A29** (2014) 1430044, arXiv:1402.4476 [nucl-ex].
- [52] ALICE Collaboration, P. Cortese *et al.*, “ALICE technical design report on forward detectors: FMD, T0 and V0”, tech. rep., 9, 2004.
- [53] ALICE Collaboration, K. Aamodt *et al.*, “Alignment of the ALICE Inner Tracking System with cosmic-ray tracks”, *JINST* **5** (2010) P03003, arXiv:1001.0502 [physics.ins-det].
- [54] J. Alme *et al.*, “The ALICE TPC, a large 3-dimensional tracking device with fast readout for ultra-high multiplicity events”, *Nucl. Instrum. Meth.* **A622** (2010) 316–367, arXiv:1001.1950 [physics.ins-det].
- [55] ALICE Collaboration, B. Abelev *et al.*, “Long-range angular correlations on the near and away side in p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV”, *Phys. Lett.* **B719** (2013) 29–41, arXiv:1212.2001 [nucl-ex].
- [56] ALICE Collaboration, B. Abelev *et al.*, “Measurement of Event Background Fluctuations for Charged Particle Jet Reconstruction in Pb-Pb collisions at $\sqrt{s_{NN}} = 2.76$ TeV”, *JHEP* **03** (2012) 053, arXiv:1201.2423 [hep-ex].
- [57] ALICE Collaboration, B. B. Abelev *et al.*, “Performance of the ALICE Experiment at the CERN LHC”, *Int. J. Mod. Phys. A* **29** (2014) 1430044, arXiv:1402.4476 [nucl-ex].
- [58] M. Cacciari, G. P. Salam, and G. Soyez, “FastJet User Manual”, *Eur. Phys. J.* **C72** (2012) 1896, arXiv:1111.6097 [hep-ph].
- [59] CMS Collaboration, S. Chatrchyan *et al.*, “Shape, Transverse Size, and Charged Hadron Multiplicity of Jets in pp Collisions at 7 TeV”, *JHEP* **06** (2012) 160, arXiv:1204.3170 [hep-ex].
- [60] T. Auye, “Unfolding algorithms and tests using RooUnfold”, in *Proceedings, PHYSTAT 2011 Workshop on Statistical Issues Related to Discovery Claims in Search Experiments and Unfolding, CERN, Geneva, Switzerland 17-20 January 2011*, pp. 313–318, CERN. CERN, Geneva, 2011. arXiv:1105.1160 [physics.data-an].

- [61] P. Skands, S. Carrazza, and J. Rojo, “Tuning PYTHIA 8.1: the Monash 2013 Tune”, *Eur. Phys. J.* **C74** no. 8, (2014) 3024, arXiv:1404.5630 [hep-ph].
- [62] P. Z. Skands, “Tuning Monte Carlo Generators: The Perugia Tunes”, *Phys. Rev.* **D82** (2010) 074018, arXiv:1005.3457 [hep-ph].
- [63] **GEANT4** Collaboration, S. Agostinelli *et al.*, “GEANT4: A Simulation toolkit”, *Nucl. Instrum. Meth.* **A506** (2003) 250–303.
- [64] **Geant4** Collaboration, M. Asai, A. Dotti, M. Verderi, and D. H. Wright, “Recent developments in Geant4”, *Annals Nucl. Energy* **82** (2015) 19–28.
- [65] **ALICE** Collaboration, B. Abelev *et al.*, “Charged jet cross sections and properties in proton-proton collisions at $\sqrt{s} = 7$ TeV”, *Phys. Rev.* **D91** no. 11, (2015) 112012, arXiv:1411.4969 [nucl-ex].
- [66] **ALICE** Collaboration, S. Acharya *et al.*, “Charged jet cross section and fragmentation in proton-proton collisions at $\sqrt{s} = 7$ TeV”, *Phys. Rev.* **D99** no. 1, (2019) 012016, arXiv:1809.03232 [nucl-ex].
- [67] **ALICE** Collaboration, B. Abelev *et al.*, “Pseudorapidity density of charged particles in p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV”, *Phys. Rev. Lett.* **110** no. 3, (2013) 032301, arXiv:1210.3615 [nucl-ex].
- [68] **ALICE** Collaboration, B. Abelev *et al.*, “Transverse momentum dependence of inclusive primary charged-particle production in p-Pb collisions at $\sqrt{s_{NN}} = 5.02$ TeV”, *Eur. Phys. J.* **C74** no. 9, (2014) 3054, arXiv:1405.2737 [nucl-ex].
- [69] A. Hocker and V. Kartvelishvili, “SVD approach to data unfolding”, *Nucl. Instrum. Meth. A* **372** (1996) 469–481, arXiv:hep-ph/9509307.
- [70] A. J. Larkoski, J. Thaler, and W. J. Waalewijn, “Gaining (Mutual) Information about Quark/Gluon Discrimination”, *JHEP* **11** (2014) 129, arXiv:1408.3122 [hep-ph].
- [71] B. Andersson, G. Gustafson, G. Ingelman, and T. Sjöstrand, “Parton Fragmentation and String Dynamics”, *Phys. Rept.* **97** (1983) 31–145.
- [72] S. Gieseke, P. Stephens, and B. Webber, “New formalism for QCD parton showers”, *JHEP* **12** (2003) 045, arXiv:hep-ph/0310083 [hep-ph].
- [73] T. Sjostrand and P. Z. Skands, “Transverse-momentum-ordered showers and interleaved multiple interactions”, *Eur. Phys. J.* **C39** (2005) 129–154, arXiv:hep-ph/0408302 [hep-ph].
- [74] R. Corke and T. Sjostrand, “Interleaved Parton Showers and Tuning Prospects”, *JHEP* **03** (2011) 032, arXiv:1011.1759 [hep-ph].
- [75] N. Fischer, S. Prestel, M. Ritzmann, and P. Skands, “Vincia for Hadron Colliders”, *Eur. Phys. J.* **C76** no. 11, (2016) 589, arXiv:1605.06142 [hep-ph].
- [76] S. Höche and S. Prestel, “The midpoint between dipole and parton showers”, *Eur. Phys. J.* **C75** no. 9, (2015) 461, arXiv:1506.05057 [hep-ph].
- [77] C. Oleari, “The POWHEG-BOX”, *Nucl. Phys. Proc. Suppl.* **205-206** (2010) 36–41, arXiv:1007.3893 [hep-ph].
- [78] C. Bierlich, G. Gustafson, L. Lönnblad, and H. Shah, “The Angantyr model for Heavy-Ion Collisions in PYTHIA8”, *JHEP* **10** (2018) 134, arXiv:1806.10820 [hep-ph].

A The ALICE Collaboration

S. Acharya¹⁴², D. Adamová⁹⁷, A. Adler⁷⁵, J. Adolfsson⁸², G. Aglieri Rinella³⁵, M. Agnello³¹, N. Agrawal⁵⁵, Z. Ahammed¹⁴², S. Ahmad¹⁶, S.U. Ahn⁷⁷, Z. Akbar⁵², A. Akindinov⁹⁴, M. Al-Turany¹⁰⁹, D.S.D. Albuquerque¹²⁴, D. Aleksandrov⁹⁰, B. Alessandro⁶⁰, H.M. Alfanda⁷, R. Alfaro Molina⁷², B. Ali¹⁶, Y. Ali¹⁴, A. Alici²⁶, N. Alizadehvandchali¹²⁷, A. Alkin³⁵, J. Alme²¹, T. Alt⁶⁹, L. Altenkamper²¹, I. Altsybeev¹¹⁵, M.N. Anaam⁷, C. Andrei⁴⁹, D. Andreou⁹², A. Andronic¹⁴⁵, V. Anguelov¹⁰⁶, T. Antičić¹¹⁰, F. Antinori⁵⁸, P. Antonioli⁵⁵, N. Apadula⁸¹, L. Aphecetche¹¹⁷, H. Appelshäuser⁶⁹, S. Arcelli²⁶, R. Arnaldi⁶⁰, M. Arratia⁸¹, I.C. Arsene²⁰, M. Arslanodok^{147,106}, A. Augustinus³⁵, R. Averbeck¹⁰⁹, S. Aziz⁷⁹, M.D. Azmi¹⁶, A. Badalà⁵⁷, Y.W. Baek⁴², X. Bai¹⁰⁹, R. Bailhache⁶⁹, R. Bala¹⁰³, A. Balbino³¹, A. Baldisseri¹³⁹, M. Ball⁴⁴, D. Banerjee⁴, R. Barbera²⁷, L. Barioglio²⁵, M. Barlou⁸⁶, G.G. Barnaföldi¹⁴⁶, L.S. Barnby⁹⁶, V. Barret¹³⁶, C. Bartels¹²⁹, K. Barth³⁵, E. Bartsch⁶⁹, F. Baruffaldi²⁸, N. Bastid¹³⁶, S. Basu^{82,144}, G. Batigne¹¹⁷, B. Batyunya⁷⁶, D. Bauri⁵⁰, J.L. Bazo Alba¹¹⁴, I.G. Bearden⁹¹, C. Beattie¹⁴⁷, I. Belikov¹³⁸, A.D.C. Bell Hechavarria¹⁴⁵, F. Bellini³⁵, R. Bellwied¹²⁷, S. Belokurova¹¹⁵, V. Belyaev⁹⁵, G. Bencedi^{70,146}, S. Beole²⁵, A. Bercuci⁴⁹, Y. Berdnikov¹⁰⁰, A. Berdnikova¹⁰⁶, D. Berenyi¹⁴⁶, L. Bergmann¹⁰⁶, M.G. Besoiu⁶⁸, L. Betev³⁵, P.P. Bhaduri¹⁴², A. Bhasin¹⁰³, I.R. Bhat¹⁰³, M.A. Bhat⁴, B. Bhattacharjee⁴³, P. Bhattacharya²³, A. Bianchi²⁵, L. Bianchi²⁵, N. Bianchi⁵³, J. Bielčák³⁸, J. Bielčíková⁹⁷, A. Bilandzic¹⁰⁷, G. Biro¹⁴⁶, S. Biswas⁴, J.T. Blair¹²¹, D. Blau⁹⁰, M.B. Blidaru¹⁰⁹, C. Blume⁶⁹, G. Boca²⁹, F. Bock⁹⁸, A. Bogdanov⁹⁵, S. Boi²³, J. Bok⁶², L. Boldizsár¹⁴⁶, A. Bolozdynya⁹⁵, M. Bombara³⁹, G. Bonomi¹⁴¹, H. Borel¹³⁹, A. Borissov^{83,95}, H. Bossi¹⁴⁷, E. Botta²⁵, L. Bratrud⁶⁹, P. Braun-Munzinger¹⁰⁹, M. Bregant¹²³, M. Broz³⁸, G.E. Bruno^{108,34}, M.D. Buckland¹²⁹, D. Budnikov¹¹¹, H. Buesching⁶⁹, S. Bufalino³¹, O. Bugnon¹¹⁷, P. Buhler¹¹⁶, P. Buncic³⁵, Z. Buthelezi^{73,133}, J.B. Butt¹⁴, S.A. Bysiak¹²⁰, D. Caffarri⁹², A. Caliva¹⁰⁹, E. Calvo Villar¹¹⁴, J.M.M. Camacho¹²², R.S. Camacho⁴⁶, P. Camerini²⁴, F.D.M. Canedo¹²³, A.A. Capon¹¹⁶, F. Carnesecchi²⁶, R. Caron¹³⁹, J. Castillo Castellanos¹³⁹, E.A.R. Casula⁵⁶, F. Catalano³¹, C. Ceballos Sanchez⁷⁶, P. Chakraborty⁵⁰, S. Chandra¹⁴², W. Chang⁷, S. Chapeland³⁵, M. Chartier¹²⁹, S. Chattopadhyay¹⁴², S. Chattopadhyay¹¹², A. Chauvin²³, C. Cheshkov¹³⁷, B. Cheynis¹³⁷, V. Chibante Barroso³⁵, D.D. Chinellato¹²⁴, S. Cho⁶², P. Chochula³⁵, P. Christakoglou⁹², C.H. Christensen⁹¹, P. Christiansen⁸², T. Chujo¹³⁵, C. Cicalo⁵⁶, L. Cifarelli²⁶, F. Cindolo⁵⁵, M.R. Ciupek¹⁰⁹, G. Clai^{II,55}, J. Cleymans¹²⁶, F. Colamaria⁵⁴, J.S. Colburn¹¹³, D. Colella⁵⁴, A. Collu⁸¹, M. Colocci^{35,26}, M. Concas^{III,60}, G. Conesa Balbastre⁸⁰, Z. Conesa del Valle⁷⁹, G. Contin²⁴, J.G. Contreras³⁸, T.M. Cormier⁹⁸, P. Cortese³², M.R. Cosentino¹²⁵, F. Costa³⁵, S. Costanza²⁹, P. Crochet¹³⁶, E. Cuautle⁷⁰, P. Cui⁷, L. Cunqueiro⁹⁸, A. Dainese⁵⁸, F.P.A. Damas^{117,139}, M.C. Danisch¹⁰⁶, A. Danu⁶⁸, D. Das¹¹², I. Das¹¹², P. Das⁸⁸, P. Das⁴, S. Das⁴, S. Dash⁵⁰, S. De⁸⁸, A. De Caro³⁰, G. de Cataldo⁵⁴, L. De Cilladi²⁵, J. de Cuveland⁴⁰, A. De Falco²³, D. De Gruttola³⁰, N. De Marco⁶⁰, C. De Martin²⁴, S. De Pasquale³⁰, S. Deb⁵¹, H.F. Degenhardt¹²³, K.R. Deja¹⁴³, S. Delsanto²⁵, W. Deng⁷, P. Dhankher¹⁹, D. Di Bari³⁴, A. Di Mauro³⁵, R.A. Diaz⁸, T. Dietel¹²⁶, P. Dillenseger⁶⁹, Y. Ding⁷, R. Divià³⁵, D.U. Dixit¹⁹, Ø. Djuvsland²¹, U. Dmitrieva⁶⁴, J. Do⁶², A. Dobrin⁶⁸, B. Dönigus⁶⁹, O. Dordic²⁰, A.K. Dubey¹⁴², A. Dubla^{109,92}, S. Dudi¹⁰², M. Dukhishyam⁸⁸, P. Dupieux¹³⁶, T.M. Eder¹⁴⁵, R.J. Ehlers⁹⁸, V.N. Eikeland²¹, D. Elia⁵⁴, B. Erasmus¹¹⁷, F. Ercolessi²⁶, F. Erhardt¹⁰¹, A. Erokhin¹¹⁵, M.R. Ersdal²¹, B. Espagnon⁷⁹, G. Eulisse³⁵, D. Evans¹¹³, S. Evdokimov⁹³, L. Fabbietti¹⁰⁷, M. Faggin²⁸, J. Faivre⁸⁰, F. Fan⁷, A. Fantoni⁵³, M. Fasel⁹⁸, P. Fedichio³¹, A. Feliciello⁶⁰, G. Feofilov¹¹⁵, A. Fernández Téllez⁴⁶, A. Ferrero¹³⁹, A. Ferretti²⁵, A. Festanti³⁵, V.J.G. Feuillard¹⁰⁶, J. Figiel¹²⁰, S. Filchagin¹¹¹, D. Finogeev⁶⁴, F.M. Fionda²¹, G. Fiorenza⁵⁴, F. Flor¹²⁷, A.N. Flores¹²¹, S. Foertsch⁷³, P. Foka¹⁰⁹, S. Fokin⁹⁰, E. Fragiaco⁶¹, U. Fuchs³⁵, C. Furget⁸⁰, A. Furs⁶⁴, M. Fusco Girard³⁰, J.J. Gaardhøje⁹¹, M. Gagliardi²⁵, A.M. Gago¹¹⁴, A. Gal¹³⁸, C.D. Galvan¹²², P. Ganoti⁸⁶, C. Garabatos¹⁰⁹, J.R.A. Garcia⁴⁶, E. Garcia-Solis¹⁰, K. Garg¹¹⁷, C. Gargiulo³⁵, A. Garibli⁸⁹, K. Garner¹⁴⁵, P. Gasik¹⁰⁷, E.F. Gauger¹²¹, M.B. Gay Ducati⁷¹, M. Germain¹¹⁷, J. Ghosh¹¹², P. Ghosh¹⁴², S.K. Ghosh⁴, M. Giacalone²⁶, P. Gianotti⁵³, P. Giubellino^{109,60}, P. Giubilato²⁸, A.M.C. Glaenger¹³⁹, P. Glässel¹⁰⁶, V. Gonzalez¹⁴⁴, L.H. González-Trueba⁷², S. Gorbunov⁴⁰,

L. Görlich¹²⁰, S. Gotovac³⁶, V. Grabski⁷², L.K. Graczykowski¹⁴³, K.L. Graham¹¹³, L. Greiner⁸¹,
 A. Grelli⁶³, C. Grigoras³⁵, V. Grigoriev⁹⁵, A. Grigoryan^{1,1}, S. Grigoryan⁷⁶, O.S. Groettvik²¹,
 F. Grosa⁶⁰, J.F. Grosse-Oetringhaus³⁵, R. Grosso¹⁰⁹, R. Guernane⁸⁰, M. Guilbaud¹¹⁷, M. Guittiere¹¹⁷,
 K. Gulbrandsen⁹¹, T. Gunji¹³⁴, A. Gupta¹⁰³, R. Gupta¹⁰³, I.B. Guzman⁴⁶, R. Haake¹⁴⁷, M.K. Habib¹⁰⁹,
 C. Hadjidakis⁷⁹, H. Hamagaki⁸⁴, G. Hamar¹⁴⁶, M. Hamid⁷, R. Hannigan¹²¹, M.R. Haque^{143,88},
 A. Harlenderova¹⁰⁹, J.W. Harris¹⁴⁷, A. Harton¹⁰, J.A. Hasenbichler³⁵, H. Hassan⁹⁸, D. Hatzifotiadou⁵⁵,
 P. Hauer⁴⁴, L.B. Havener¹⁴⁷, S. Hayashi¹³⁴, S.T. Heckel¹⁰⁷, E. Hellbär⁶⁹, H. Helstrup³⁷, T. Herman³⁸,
 E.G. Hernandez⁴⁶, G. Herrera Corral⁹, F. Herrmann¹⁴⁵, K.F. Hetland³⁷, H. Hillemanns³⁵, C. Hills¹²⁹,
 B. Hippolyte¹³⁸, B. Hohlweger¹⁰⁷, J. Honermann¹⁴⁵, G.H. Hong¹⁴⁸, D. Horak³⁸, S. Hornung¹⁰⁹,
 R. Hosokawa¹⁵, P. Hristov³⁵, C. Huang⁷⁹, C. Hughes¹³², P. Huhn⁶⁹, T.J. Humanic⁹⁹, H. Hushnud¹¹²,
 L.A. Husova¹⁴⁵, N. Hussain⁴³, D. Hutter⁴⁰, J.P. Iddon^{35,129}, R. Ilkaev¹¹¹, H. Ilyas¹⁴, M. Inaba¹³⁵,
 G.M. Innocenti³⁵, M. Ippolitov⁹⁰, A. Isakov^{38,97}, M.S. Islam¹¹², M. Ivanov¹⁰⁹, V. Ivanov¹⁰⁰,
 V. Izucheev⁹³, B. Jacak⁸¹, N. Jacazio^{35,55}, P.M. Jacobs⁸¹, S. Jadlovská¹¹⁹, J. Jadlovsky¹¹⁹, S. Jaelani⁶³,
 C. Jahnke¹²³, M.J. Jakubowska¹⁴³, M.A. Janik¹⁴³, T. Janson⁷⁵, M. Jercic¹⁰¹, O. Jevons¹¹³, M. Jin¹²⁷,
 F. Jonas^{98,145}, P.G. Jones¹¹³, J. Jung⁶⁹, M. Jung⁶⁹, A. Jusko¹¹³, P. Kalinak⁶⁵, A. Kalweit³⁵, V. Kaplin⁹⁵,
 S. Kar⁷, A. Karasu Uysal⁷⁸, D. Karatovic¹⁰¹, O. Karavichev⁶⁴, T. Karavicheva⁶⁴, P. Karczmarczyk¹⁴³,
 E. Karpechev⁶⁴, A. Kazantsev⁹⁰, U. Kebschull⁷⁵, R. Keidel⁴⁸, M. Keil³⁵, B. Ketzer⁴⁴, Z. Khabanova⁹²,
 A.M. Khan⁷, S. Khan¹⁶, A. Khanzadeev¹⁰⁰, Y. Kharlov⁹³, A. Khatun¹⁶, A. Khuntia¹²⁰, B. Kileng³⁷,
 B. Kim^{17,62}, D. Kim¹⁴⁸, D.J. Kim¹²⁸, E.J. Kim⁷⁴, H. Kim¹⁷, J. Kim¹⁴⁸, J.S. Kim⁴², J. Kim¹⁰⁶, J. Kim¹⁴⁸,
 J. Kim⁷⁴, M. Kim¹⁰⁶, S. Kim¹⁸, T. Kim¹⁴⁸, T. Kim¹⁴⁸, S. Kirsch⁶⁹, I. Kisel⁴⁰, S. Kiselev⁹⁴, A. Kisiel¹⁴³,
 J.L. Klay⁶, J. Klein^{35,60}, S. Klein⁸¹, C. Klein-Bösing¹⁴⁵, M. Kleiner⁶⁹, T. Klemenz¹⁰⁷, A. Kluge³⁵,
 A.G. Knospe¹²⁷, C. Kobdaj¹¹⁸, M.K. Köhler¹⁰⁶, T. Kollegger¹⁰⁹, A. Kondratyev⁷⁶, N. Kondratyeva⁹⁵,
 E. Kondratyuk⁹³, J. König⁶⁹, S.A. Königstorfer¹⁰⁷, P.J. Konopka^{2,35}, G. Kornakov¹⁴³, S.D. Koryciak²,
 L. Koska¹¹⁹, O. Kovalenko⁸⁷, V. Kovalenko¹¹⁵, M. Kowalski¹²⁰, I. Králik⁶⁵, A. Kravčáková³⁹,
 L. Kreis¹⁰⁹, M. Krivda^{113,65}, F. Krizek⁹⁷, K. Krizkova Gajdosova³⁸, M. Kroesen¹⁰⁶, M. Krüger⁶⁹,
 E. Kryshen¹⁰⁰, M. Krzewicki⁴⁰, V. Kučera³⁵, C. Kuhn¹³⁸, P.G. Kuijper⁹², T. Kumaoka¹³⁵, L. Kumar¹⁰²,
 S. Kundu⁸⁸, P. Kurashvili⁸⁷, A. Kurepin⁶⁴, A.B. Kurepin⁶⁴, A. Kuryakin¹¹¹, S. Kushpil⁹⁷, J. Kvapil¹¹³,
 M.J. Kweon⁶², J.Y. Kwon⁶², Y. Kwon¹⁴⁸, S.L. La Pointe⁴⁰, P. La Rocca²⁷, Y.S. Lai⁸¹, A. Lakrathok¹¹⁸,
 M. Lamanna³⁵, R. Langoy¹³¹, K. Lapidus³⁵, P. Larionov⁵³, E. Laudi³⁵, L. Lautner³⁵, R. Lavicka³⁸,
 T. Lazareva¹¹⁵, R. Lea²⁴, J. Lee¹³⁵, S. Lee¹⁴⁸, J. Lehrbach⁴⁰, R.C. Lemmon⁹⁶, I. León Monzón¹²²,
 E.D. Lesser¹⁹, M. Lettrich³⁵, P. Lévai¹⁴⁶, X. Li¹¹, X.L. Li⁷, J. Lien¹³¹, R. Lietava¹¹³, B. Lim¹⁷,
 S.H. Lim¹⁷, V. Lindenstruth⁴⁰, A. Lindner⁴⁹, C. Lippmann¹⁰⁹, A. Liu¹⁹, J. Liu¹²⁹, I.M. Lofnes²¹,
 V. Loginov⁹⁵, C. Loizides⁹⁸, P. Loncar³⁶, J.A. Lopez¹⁰⁶, X. Lopez¹³⁶, E. López Torres⁸, J.R. Luhder¹⁴⁵,
 M. Lunardon²⁸, G. Luparello⁶¹, Y.G. Ma⁴¹, A. Maevskaya⁶⁴, M. Mager³⁵, S.M. Mahmood²⁰,
 T. Mahmoud⁴⁴, A. Maire¹³⁸, R.D. Majka^{1,147}, M. Malaev¹⁰⁰, Q.W. Malik²⁰, L. Malinina^{IV,76},
 D. Mal'Kevich⁹⁴, N. Mallick⁵¹, P. Malzacher¹⁰⁹, G. Mandaglio^{33,57}, V. Manko⁹⁰, F. Manso¹³⁶,
 V. Manzari⁵⁴, Y. Mao⁷, M. Marchisone¹³⁷, J. Mareš⁶⁷, G.V. Margagliotti²⁴, A. Margotti⁵⁵, A. Marín¹⁰⁹,
 C. Markert¹²¹, M. Marquard⁶⁹, N.A. Martin¹⁰⁶, P. Martinengo³⁵, J.L. Martinez¹²⁷, M.I. Martínez⁴⁶,
 G. Martínez García¹¹⁷, S. Masciocchi¹⁰⁹, M. Maserà²⁵, A. Masoni⁵⁶, L. Massacrier⁷⁹,
 A. Mastroserio^{140,54}, A.M. Mathis¹⁰⁷, O. Matonoha⁸², P.F.T. Matuoka¹²³, A. Matyja¹²⁰, C. Mayer¹²⁰,
 F. Mazzaschi²⁵, M. Mazzilli^{35,54}, M.A. Mazzoni⁵⁹, A.F. Mechler⁶⁹, F. Meddi²², Y. Melikyan⁶⁴,
 A. Menchaca-Rocha⁷², C. Mengke⁷, E. Meninno^{116,30}, A.S. Menon¹²⁷, M. Meres¹³, S. Mhlanga¹²⁶,
 Y. Miake¹³⁵, L. Micheletti²⁵, L.C. Migliorin¹³⁷, D.L. Mihaylov¹⁰⁷, K. Mikhaylov^{76,94},
 A.N. Mishra^{146,70}, D. Miśkowiec¹⁰⁹, A. Modak⁴, N. Mohammadi³⁵, A.P. Mohanty⁶³, B. Mohanty⁸⁸,
 M. Mohisin Khan¹⁶, Z. Moravcova⁹¹, C. Mordasini¹⁰⁷, D.A. Moreira De Godoy¹⁴⁵, L.A.P. Moreno⁴⁶,
 I. Morozov⁶⁴, A. Morsch³⁵, T. Mrnjavac³⁵, V. Muccifora⁵³, E. Mudnic³⁶, D. Mühlheim¹⁴⁵,
 S. Muhuri¹⁴², J.D. Mulligan⁸¹, A. Mulliri^{23,56}, M.G. Munhoz¹²³, R.H. Munzer⁶⁹, H. Murakami¹³⁴,
 S. Murray¹²⁶, L. Musa³⁵, J. Musinsky⁶⁵, C.J. Myers¹²⁷, J.W. Myrcha¹⁴³, B. Naik⁵⁰, R. Nair⁸⁷,
 B.K. Nandi⁵⁰, R. Nania⁵⁵, E. Nappi⁵⁴, M.U. Naru¹⁴, A.F. Nassirpour⁸², C. Nattrass¹³², S. Nazarenko¹¹¹,
 A. Neagu²⁰, L. Nellen⁷⁰, S.V. Nesbo³⁷, G. Neskovic⁴⁰, D. Nesterov¹¹⁵, B.S. Nielsen⁹¹, S. Nikolaev⁹⁰,

S. Nikulin⁹⁰, V. Nikulin¹⁰⁰, F. Noferini⁵⁵, S. Noh¹², P. Nomokonov⁷⁶, J. Norman¹²⁹, N. Novitzky¹³⁵,
 P. Nowakowski¹⁴³, A. Nyanin⁹⁰, J. Nystrand²¹, M. Ogino⁸⁴, A. Ohlson⁸², J. Oleniacz¹⁴³, A.C. Oliveira
 Da Silva¹³², M.H. Oliver¹⁴⁷, B.S. Onnerstad¹²⁸, C. Oppedisano⁶⁰, A. Ortiz Velasquez⁷⁰, T. Osako⁴⁷,
 A. Oskarsson⁸², J. Otwinowski¹²⁰, K. Oyama⁸⁴, Y. Pachmayer¹⁰⁶, S. Padhan⁵⁰, D. Pagano¹⁴¹, G. Paic⁷⁰,
 J. Pan¹⁴⁴, S. Panebianco¹³⁹, P. Pareek¹⁴², J. Park⁶², J.E. Parkkila¹²⁸, S. Parmar¹⁰², S.P. Pathak¹²⁷,
 B. Paul²³, J. Pazzini¹⁴¹, H. Pei⁷, T. Peitzmann⁶³, X. Peng⁷, L.G. Pereira⁷¹, H. Pereira Da Costa¹³⁹,
 D. Peresunko⁹⁰, G.M. Perez⁸, S. Perrin¹³⁹, Y. Pestov⁵, V. Petráček³⁸, M. Petrovici⁴⁹, R.P. Pezzi⁷¹,
 S. Piano⁶¹, M. Pikna¹³, P. Pillot¹¹⁷, O. Pinazza^{55,35}, L. Pinsky¹²⁷, C. Pinto²⁷, S. Pisano⁵³, M. Płoskoń⁸¹,
 M. Planinic¹⁰¹, F. Pliquet⁶⁹, M.G. Poghosyan⁹⁸, B. Polichtchouk⁹³, N. Poljak¹⁰¹, A. Pop⁴⁹,
 S. Porteboeuf-Houssais¹³⁶, J. Porter⁸¹, V. Pozdniakov⁷⁶, S.K. Prasad⁴, R. Preghenella⁵⁵, F. Prino⁶⁰,
 C.A. Pruneau¹⁴⁴, I. Pshenichnov⁶⁴, M. Puccio³⁵, S. Qiu⁹², L. Quaglia²⁵, R.E. Quishpe¹²⁷, S. Ragoni¹¹³,
 J. Rak¹²⁸, A. Rakotozafindrabe¹³⁹, L. Ramello³², F. Rami¹³⁸, S.A.R. Ramirez⁴⁶, A.G.T. Ramos³⁴,
 R. Raniwala¹⁰⁴, S. Raniwala¹⁰⁴, S.S. Räsänen⁴⁵, R. Rath⁵¹, I. Ravasenga⁹², K.F. Read^{98,132},
 A.R. Redelbach⁴⁰, K. Redlich⁸⁷, A. Rehman²¹, P. Reichelt⁶⁹, F. Reidt³⁵, R. Renfordt⁶⁹,
 Z. Rescakova³⁹, K. Reygers¹⁰⁶, A. Riabov¹⁰⁰, V. Riabov¹⁰⁰, T. Richert^{82,91}, M. Richter²⁰, P. Riedler³⁵,
 W. Riegler³⁵, F. Riggi²⁷, C. Ristea⁶⁸, S.P. Rode⁵¹, M. Rodríguez Cahuantzi⁴⁶, K. Røed²⁰, R. Rogalev⁹³,
 E. Rogochaya⁷⁶, T.S. Rogoschinski⁶⁹, D. Rohr³⁵, D. Röhrich²¹, P.F. Rojas⁴⁶, P.S. Rokita¹⁴³,
 F. Ronchetti⁵³, A. Rosano^{33,57}, E.D. Rosas⁷⁰, A. Rossi⁵⁸, A. Rotondi²⁹, A. Roy⁵¹, P. Roy¹¹²,
 N. Rubini²⁶, O.V. Rueda⁸², R. Rui²⁴, B. Rumyantsev⁷⁶, A. Rustamov⁸⁹, E. Ryabinkin⁹⁰, Y. Ryabov¹⁰⁰,
 A. Rybicki¹²⁰, H. Rytönen¹²⁸, O.A.M. Saarimäki⁴⁵, R. Sadek¹¹⁷, S. Sadovsky⁹³, J. Saetre²¹,
 K. Šafařík³⁸, S.K. Saha¹⁴², S. Saha⁸⁸, B. Sahoo⁵⁰, P. Sahoo⁵⁰, R. Sahoo⁵¹, S. Sahoo⁶⁶, D. Sahu⁵¹,
 P.K. Sahu⁶⁶, J. Saini¹⁴², S. Sakai¹³⁵, S. Sambyal¹⁰³, V. Samsonov^{100,95}, D. Sarkar¹⁴⁴, N. Sarkar¹⁴²,
 P. Sarma⁴³, V.M. Sarti¹⁰⁷, M.H.P. Sas^{147,63}, J. Schambach^{98,121}, H.S. Scheid⁶⁹, C. Schiaua⁴⁹,
 R. Schicker¹⁰⁶, A. Schmah¹⁰⁶, C. Schmidt¹⁰⁹, H.R. Schmidt¹⁰⁵, M.O. Schmidt¹⁰⁶, M. Schmidt¹⁰⁵,
 N.V. Schmidt^{98,69}, A.R. Schmier¹³², R. Schotter¹³⁸, J. Schukraft³⁵, Y. Schutz¹³⁸, K. Schwarz¹⁰⁹,
 K. Schweda¹⁰⁹, G. Scioli²⁶, E. Scomparin⁶⁰, J.E. Seger¹⁵, Y. Sekiguchi¹³⁴, D. Sekihata¹³⁴,
 I. Selyuzhenkov^{109,95}, S. Senyukov¹³⁸, J.J. Seo⁶², D. Serebryakov⁶⁴, L. Šerkšnytė¹⁰⁷, A. Sevcenco⁶⁸,
 A. Shabanov⁶⁴, A. Shabetai¹¹⁷, R. Shahoyan³⁵, W. Shaikh¹¹², A. Shangaraev⁹³, A. Sharma¹⁰²,
 H. Sharma¹²⁰, M. Sharma¹⁰³, N. Sharma¹⁰², S. Sharma¹⁰³, O. Sheibani¹²⁷, A.I. Sheikh¹⁴², K. Shigaki⁴⁷,
 M. Shimomura⁸⁵, S. Shirinkin⁹⁴, Q. Shou⁴¹, Y. Sibiriak⁹⁰, S. Siddhanta⁵⁶, T. Siemiarzczuk⁸⁷,
 D. Silvermyr⁸², G. Simatovic⁹², G. Simonetti³⁵, B. Singh¹⁰⁷, R. Singh⁸⁸, R. Singh¹⁰³, R. Singh⁵¹,
 V.K. Singh¹⁴², V. Singhal¹⁴², T. Sinha¹¹², B. Sitar¹³, M. Sitta³², T.B. Skaali²⁰, M. Slupecki⁴⁵,
 N. Smirnov¹⁴⁷, R.J.M. Snellings⁶³, T.W. Snellman¹²⁸, C. Soncco¹¹⁴, J. Song¹²⁷, A. Songmoolnak¹¹⁸,
 F. Soramel²⁸, S. Sorensen¹³², I. Sputowska¹²⁰, J. Stachel¹⁰⁶, I. Stan⁶⁸, P.J. Steffanic¹³²,
 S.F. Stiefelmaier¹⁰⁶, D. Stocco¹¹⁷, M.M. Stortvedt³⁷, L.D. Stritto³⁰, C.P. Stylianidis⁹²,
 A.A.P. Suaide¹²³, T. Sugitate⁴⁷, C. Suire⁷⁹, M. Suljic³⁵, R. Sultanov⁹⁴, M. Šumbera⁹⁷, V. Sumberia¹⁰³,
 S. Sumowidagdo⁵², S. Swain⁶⁶, A. Szabo¹³, I. Szarka¹³, U. Tabassam¹⁴, S.F. Taghavi¹⁰⁷,
 G. TAILLEPIED¹³⁶, J. Takahashi¹²⁴, G.J. Tambave²¹, S. Tang^{136,7}, Z. Tang¹³⁰, M. Tardini¹¹⁷,
 M.G. Tarzila⁴⁹, A. Tauro³⁵, G. Tejada Muñoz⁴⁶, A. Telesca³⁵, L. Terlizzi²⁵, C. Terrevoli¹²⁷,
 G. Tersimonov³, S. Thakur¹⁴², D. Thomas¹²¹, R. Tieulent¹³⁷, A. Tikhonov⁶⁴, A.R. Timmins¹²⁷,
 M. Tkacik¹¹⁹, A. Toia⁶⁹, N. Topilskaya⁶⁴, M. Toppi⁵³, F. Torales-Acosta¹⁹, S.R. Torres^{38,9},
 A. Trifiro^{33,57}, S. Tripathy⁷⁰, T. Tripathy⁵⁰, S. Trogolo²⁸, G. Trombetta³⁴, L. Tropp³⁹, V. Trubnikov³,
 W.H. Trzaska¹²⁸, T.P. Trzcinski¹⁴³, B.A. Trzeciak³⁸, A. Tumkin¹¹¹, R. Turrisi⁵⁸, T.S. Tveter²⁰,
 K. Ullaland²¹, E.N. Umaka¹²⁷, A. Uras¹³⁷, G.L. Usai²³, M. Vala³⁹, N. Valle²⁹, S. Vallero⁶⁰, N. van der
 Kolk⁶³, L.V.R. van Doremalen⁶³, M. van Leeuwen⁹², P. Vande Vyvre³⁵, D. Varga¹⁴⁶, Z. Varga¹⁴⁶,
 M. Varga-Kofarago¹⁴⁶, A. Vargas⁴⁶, M. Vasileiou⁸⁶, A. Vasiliev⁹⁰, O. Vázquez Doce¹⁰⁷,
 V. Vechernin¹¹⁵, E. Vercellin²⁵, S. Vergara Limón⁴⁶, L. Vermunt⁶³, R. Vértesi¹⁴⁶, M. Verweij⁶³,
 L. Vickovic³⁶, Z. Vilakazi¹³³, O. Villalobos Baillie¹¹³, G. VINO⁵⁴, A. Vinogradov⁹⁰, T. Virgili³⁰,
 V. Vislavicius⁹¹, A. Vodopyanov⁷⁶, B. Volkel³⁵, M.A. Völkl¹⁰⁵, K. Voloshin⁹⁴, S.A. Voloshin¹⁴⁴,
 G. Volpe³⁴, B. von Haller³⁵, I. Vorobyev¹⁰⁷, D. Voscek¹¹⁹, J. Vrláková³⁹, B. Wagner²¹, M. Weber¹¹⁶,

A. Wegrzynek³⁵, S.C. Wenzel³⁵, J.P. Wessels¹⁴⁵, J. Wiechula⁶⁹, J. Wikne²⁰, G. Wilk⁸⁷, J. Wilkinson¹⁰⁹, G.A. Willems¹⁴⁵, E. Willsher¹¹³, B. Windelband¹⁰⁶, M. Winn¹³⁹, W.E. Witt¹³², J.R. Wright¹²¹, Y. Wu¹³⁰, R. Xu⁷, S. Yalcin⁷⁸, Y. Yamaguchi⁴⁷, K. Yamakawa⁴⁷, S. Yang²¹, S. Yano^{47,139}, Z. Yin⁷, H. Yokoyama⁶³, I.-K. Yoo¹⁷, J.H. Yoon⁶², S. Yuan²¹, A. Yuncu¹⁰⁶, V. Yurchenko³, V. Zaccolo²⁴, A. Zaman¹⁴, C. Zampolli³⁵, H.J.C. Zanoli⁶³, N. Zardoshti³⁵, A. Zarochentsev¹¹⁵, P. Závada⁶⁷, N. Zaviyalov¹¹¹, H. Zbroszczyk¹⁴³, M. Zhalov¹⁰⁰, S. Zhang⁴¹, X. Zhang⁷, Y. Zhang¹³⁰, V. Zherebchevskii¹¹⁵, Y. Zhi¹¹, D. Zhou⁷, Y. Zhou⁹¹, J. Zhu^{7,109}, Y. Zhu⁷, A. Zichichi²⁶, G. Zinovjev³, N. Zurlo¹⁴¹

Affiliation Notes

^I Deceased

^{II} Also at: Italian National Agency for New Technologies, Energy and Sustainable Economic Development (ENEA), Bologna, Italy

^{III} Also at: Dipartimento DET del Politecnico di Torino, Turin, Italy

^{IV} Also at: M.V. Lomonosov Moscow State University, D.V. Skobeltsyn Institute of Nuclear, Physics, Moscow, Russia

^V Also at: Institute of Theoretical Physics, University of Wroclaw, Poland

Collaboration Institutes

¹ A.I. Alikhanyan National Science Laboratory (Yerevan Physics Institute) Foundation, Yerevan, Armenia

² AGH University of Science and Technology, Cracow, Poland

³ Bogolyubov Institute for Theoretical Physics, National Academy of Sciences of Ukraine, Kiev, Ukraine

⁴ Bose Institute, Department of Physics and Centre for Astroparticle Physics and Space Science (CAPSS), Kolkata, India

⁵ Budker Institute for Nuclear Physics, Novosibirsk, Russia

⁶ California Polytechnic State University, San Luis Obispo, California, United States

⁷ Central China Normal University, Wuhan, China

⁸ Centro de Aplicaciones Tecnológicas y Desarrollo Nuclear (CEADEN), Havana, Cuba

⁹ Centro de Investigación y de Estudios Avanzados (CINVESTAV), Mexico City and Mérida, Mexico

¹⁰ Chicago State University, Chicago, Illinois, United States

¹¹ China Institute of Atomic Energy, Beijing, China

¹² Chungbuk National University, Cheongju, Republic of Korea

¹³ Comenius University Bratislava, Faculty of Mathematics, Physics and Informatics, Bratislava, Slovakia

¹⁴ COMSATS University Islamabad, Islamabad, Pakistan

¹⁵ Creighton University, Omaha, Nebraska, United States

¹⁶ Department of Physics, Aligarh Muslim University, Aligarh, India

¹⁷ Department of Physics, Pusan National University, Pusan, Republic of Korea

¹⁸ Department of Physics, Sejong University, Seoul, Republic of Korea

¹⁹ Department of Physics, University of California, Berkeley, California, United States

²⁰ Department of Physics, University of Oslo, Oslo, Norway

²¹ Department of Physics and Technology, University of Bergen, Bergen, Norway

²² Dipartimento di Fisica dell'Università 'La Sapienza' and Sezione INFN, Rome, Italy

²³ Dipartimento di Fisica dell'Università and Sezione INFN, Cagliari, Italy

- ²⁴ Dipartimento di Fisica dell'Università and Sezione INFN, Trieste, Italy
- ²⁵ Dipartimento di Fisica dell'Università and Sezione INFN, Turin, Italy
- ²⁶ Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Bologna, Italy
- ²⁷ Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Catania, Italy
- ²⁸ Dipartimento di Fisica e Astronomia dell'Università and Sezione INFN, Padova, Italy
- ²⁹ Dipartimento di Fisica e Nucleare e Teorica, Università di Pavia and Sezione INFN, Pavia, Italy
- ³⁰ Dipartimento di Fisica 'E.R. Caianiello' dell'Università and Gruppo Collegato INFN, Salerno, Italy
- ³¹ Dipartimento DISAT del Politecnico and Sezione INFN, Turin, Italy
- ³² Dipartimento di Scienze e Innovazione Tecnologica dell'Università del Piemonte Orientale and INFN Sezione di Torino, Alessandria, Italy
- ³³ Dipartimento di Scienze MIFT, Università di Messina, Messina, Italy
- ³⁴ Dipartimento Interateneo di Fisica 'M. Merlin' and Sezione INFN, Bari, Italy
- ³⁵ European Organization for Nuclear Research (CERN), Geneva, Switzerland
- ³⁶ Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, University of Split, Split, Croatia
- ³⁷ Faculty of Engineering and Science, Western Norway University of Applied Sciences, Bergen, Norway
- ³⁸ Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Prague, Czech Republic
- ³⁹ Faculty of Science, P.J. Šafárik University, Košice, Slovakia
- ⁴⁰ Frankfurt Institute for Advanced Studies, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
- ⁴¹ Fudan University, Shanghai, China
- ⁴² Gangneung-Wonju National University, Gangneung, Republic of Korea
- ⁴³ Gauhati University, Department of Physics, Guwahati, India
- ⁴⁴ Helmholtz-Institut für Strahlen- und Kernphysik, Rheinische Friedrich-Wilhelms-Universität Bonn, Bonn, Germany
- ⁴⁵ Helsinki Institute of Physics (HIP), Helsinki, Finland
- ⁴⁶ High Energy Physics Group, Universidad Autónoma de Puebla, Puebla, Mexico
- ⁴⁷ Hiroshima University, Hiroshima, Japan
- ⁴⁸ Hochschule Worms, Zentrum für Technologietransfer und Telekommunikation (ZTT), Worms, Germany
- ⁴⁹ Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest, Romania
- ⁵⁰ Indian Institute of Technology Bombay (IIT), Mumbai, India
- ⁵¹ Indian Institute of Technology Indore, Indore, India
- ⁵² Indonesian Institute of Sciences, Jakarta, Indonesia
- ⁵³ INFN, Laboratori Nazionali di Frascati, Frascati, Italy
- ⁵⁴ INFN, Sezione di Bari, Bari, Italy
- ⁵⁵ INFN, Sezione di Bologna, Bologna, Italy
- ⁵⁶ INFN, Sezione di Cagliari, Cagliari, Italy
- ⁵⁷ INFN, Sezione di Catania, Catania, Italy
- ⁵⁸ INFN, Sezione di Padova, Padova, Italy
- ⁵⁹ INFN, Sezione di Roma, Rome, Italy
- ⁶⁰ INFN, Sezione di Torino, Turin, Italy
- ⁶¹ INFN, Sezione di Trieste, Trieste, Italy
- ⁶² Inha University, Incheon, Republic of Korea
- ⁶³ Institute for Gravitational and Subatomic Physics (GRASP), Utrecht University/Nikhef, Utrecht, Netherlands
- ⁶⁴ Institute for Nuclear Research, Academy of Sciences, Moscow, Russia
- ⁶⁵ Institute of Experimental Physics, Slovak Academy of Sciences, Košice, Slovakia

- ⁶⁶ Institute of Physics, Homi Bhabha National Institute, Bhubaneswar, India
⁶⁷ Institute of Physics of the Czech Academy of Sciences, Prague, Czech Republic
⁶⁸ Institute of Space Science (ISS), Bucharest, Romania
⁶⁹ Institut für Kernphysik, Johann Wolfgang Goethe-Universität Frankfurt, Frankfurt, Germany
⁷⁰ Instituto de Ciencias Nucleares, Universidad Nacional Autónoma de México, Mexico City, Mexico
⁷¹ Instituto de Física, Universidade Federal do Rio Grande do Sul (UFRGS), Porto Alegre, Brazil
⁷² Instituto de Física, Universidad Nacional Autónoma de México, Mexico City, Mexico
⁷³ iThemba LABS, National Research Foundation, Somerset West, South Africa
⁷⁴ Jeonbuk National University, Jeonju, Republic of Korea
⁷⁵ Johann-Wolfgang-Goethe Universität Frankfurt Institut für Informatik, Fachbereich Informatik und Mathematik, Frankfurt, Germany
⁷⁶ Joint Institute for Nuclear Research (JINR), Dubna, Russia
⁷⁷ Korea Institute of Science and Technology Information, Daejeon, Republic of Korea
⁷⁸ KTO Karatay University, Konya, Turkey
⁷⁹ Laboratoire de Physique des 2 Infinis, Irène Joliot-Curie, Orsay, France
⁸⁰ Laboratoire de Physique Subatomique et de Cosmologie, Université Grenoble-Alpes, CNRS-IN2P3, Grenoble, France
⁸¹ Lawrence Berkeley National Laboratory, Berkeley, California, United States
⁸² Lund University Department of Physics, Division of Particle Physics, Lund, Sweden
⁸³ Moscow Institute for Physics and Technology, Moscow, Russia
⁸⁴ Nagasaki Institute of Applied Science, Nagasaki, Japan
⁸⁵ Nara Women's University (NWU), Nara, Japan
⁸⁶ National and Kapodistrian University of Athens, School of Science, Department of Physics, Athens, Greece
⁸⁷ National Centre for Nuclear Research, Warsaw, Poland
⁸⁸ National Institute of Science Education and Research, Homi Bhabha National Institute, Jatni, India
⁸⁹ National Nuclear Research Center, Baku, Azerbaijan
⁹⁰ National Research Centre Kurchatov Institute, Moscow, Russia
⁹¹ Niels Bohr Institute, University of Copenhagen, Copenhagen, Denmark
⁹² Nikhef, National institute for subatomic physics, Amsterdam, Netherlands
⁹³ NRC Kurchatov Institute IHEP, Protvino, Russia
⁹⁴ NRC «Kurchatov» Institute - ITEP, Moscow, Russia
⁹⁵ NRNU Moscow Engineering Physics Institute, Moscow, Russia
⁹⁶ Nuclear Physics Group, STFC Daresbury Laboratory, Daresbury, United Kingdom
⁹⁷ Nuclear Physics Institute of the Czech Academy of Sciences, Řež u Prahy, Czech Republic
⁹⁸ Oak Ridge National Laboratory, Oak Ridge, Tennessee, United States
⁹⁹ Ohio State University, Columbus, Ohio, United States
¹⁰⁰ Petersburg Nuclear Physics Institute, Gatchina, Russia
¹⁰¹ Physics department, Faculty of science, University of Zagreb, Zagreb, Croatia
¹⁰² Physics Department, Panjab University, Chandigarh, India
¹⁰³ Physics Department, University of Jammu, Jammu, India
¹⁰⁴ Physics Department, University of Rajasthan, Jaipur, India
¹⁰⁵ Physikalisches Institut, Eberhard-Karls-Universität Tübingen, Tübingen, Germany
¹⁰⁶ Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
¹⁰⁷ Physik Department, Technische Universität München, Munich, Germany
¹⁰⁸ Politecnico di Bari and Sezione INFN, Bari, Italy
¹⁰⁹ Research Division and ExtreMe Matter Institute EMMI, GSI Helmholtzzentrum für Schwerionenforschung GmbH, Darmstadt, Germany
¹¹⁰ Rudjer Bošković Institute, Zagreb, Croatia
¹¹¹ Russian Federal Nuclear Center (VNIIEF), Sarov, Russia

- 112 Saha Institute of Nuclear Physics, Homi Bhabha National Institute, Kolkata, India
113 School of Physics and Astronomy, University of Birmingham, Birmingham, United Kingdom
114 Sección Física, Departamento de Ciencias, Pontificia Universidad Católica del Perú, Lima, Peru
115 St. Petersburg State University, St. Petersburg, Russia
116 Stefan Meyer Institut für Subatomare Physik (SMI), Vienna, Austria
117 SUBATECH, IMT Atlantique, Université de Nantes, CNRS-IN2P3, Nantes, France
118 Suranaree University of Technology, Nakhon Ratchasima, Thailand
119 Technical University of Košice, Košice, Slovakia
120 The Henryk Niewodniczanski Institute of Nuclear Physics, Polish Academy of Sciences, Cracow, Poland
121 The University of Texas at Austin, Austin, Texas, United States
122 Universidad Autónoma de Sinaloa, Culiacán, Mexico
123 Universidade de São Paulo (USP), São Paulo, Brazil
124 Universidade Estadual de Campinas (UNICAMP), Campinas, Brazil
125 Universidade Federal do ABC, Santo Andre, Brazil
126 University of Cape Town, Cape Town, South Africa
127 University of Houston, Houston, Texas, United States
128 University of Jyväskylä, Jyväskylä, Finland
129 University of Liverpool, Liverpool, United Kingdom
130 University of Science and Technology of China, Hefei, China
131 University of South-Eastern Norway, Tonsberg, Norway
132 University of Tennessee, Knoxville, Tennessee, United States
133 University of the Witwatersrand, Johannesburg, South Africa
134 University of Tokyo, Tokyo, Japan
135 University of Tsukuba, Tsukuba, Japan
136 Université Clermont Auvergne, CNRS/IN2P3, LPC, Clermont-Ferrand, France
137 Université de Lyon, CNRS/IN2P3, Institut de Physique des 2 Infinis de Lyon, Lyon, France
138 Université de Strasbourg, CNRS, IPHC UMR 7178, F-67000 Strasbourg, France, Strasbourg, France
139 Université Paris-Saclay Centre d'Etudes de Saclay (CEA), IRFU, Département de Physique Nucléaire (DPhN), Saclay, France
140 Università degli Studi di Foggia, Foggia, Italy
141 Università di Brescia and Sezione INFN, Brescia, Italy
142 Variable Energy Cyclotron Centre, Homi Bhabha National Institute, Kolkata, India
143 Warsaw University of Technology, Warsaw, Poland
144 Wayne State University, Detroit, Michigan, United States
145 Westfälische Wilhelms-Universität Münster, Institut für Kernphysik, Münster, Germany
146 Wigner Research Centre for Physics, Budapest, Hungary
147 Yale University, New Haven, Connecticut, United States
148 Yonsei University, Seoul, Republic of Korea

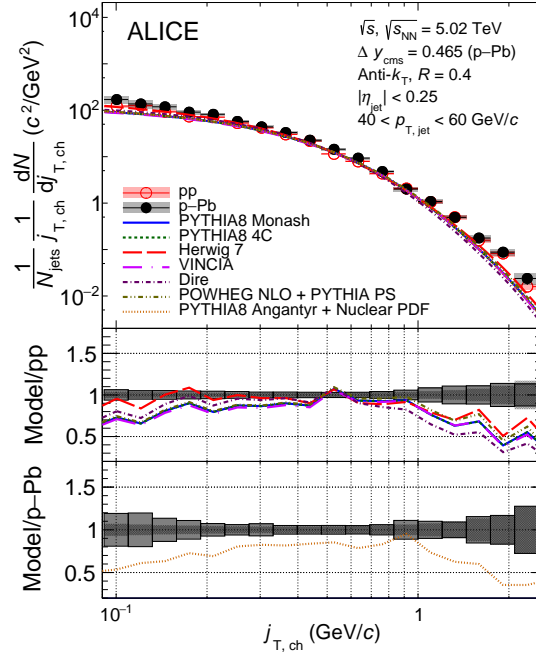
B Comparison of the j_T distributions with models for other $p_{T,\text{jet}}$ regions


Figure B.1: The j_T distribution in pp and p–Pb collisions at $\sqrt{s}, \sqrt{s_{\text{NN}}} = 5.02$ TeV for $40 < p_{T,\text{jet}} < 60$ GeV/c comparing to theoretical models in pp and p–Pb collisions.

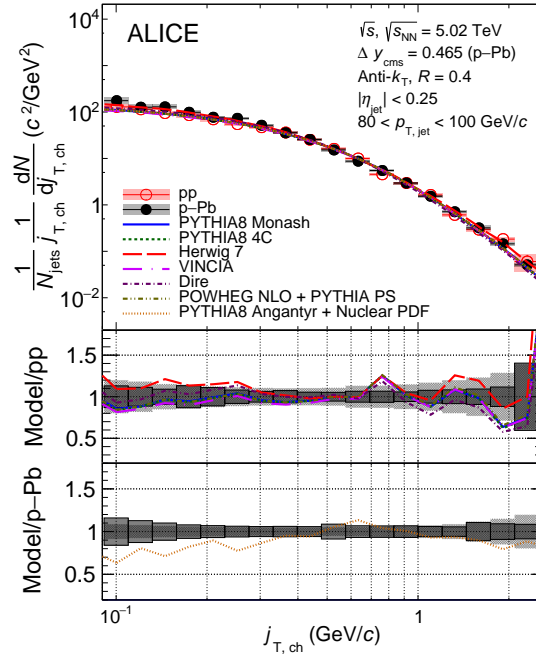


Figure B.2: The j_T distribution in p–Pb collisions at $\sqrt{s}, \sqrt{s_{\text{NN}}} = 5.02$ TeV for $80 < p_{T,\text{jet}} < 100$ GeV/c comparing to theoretical models in pp and p–Pb collisions.

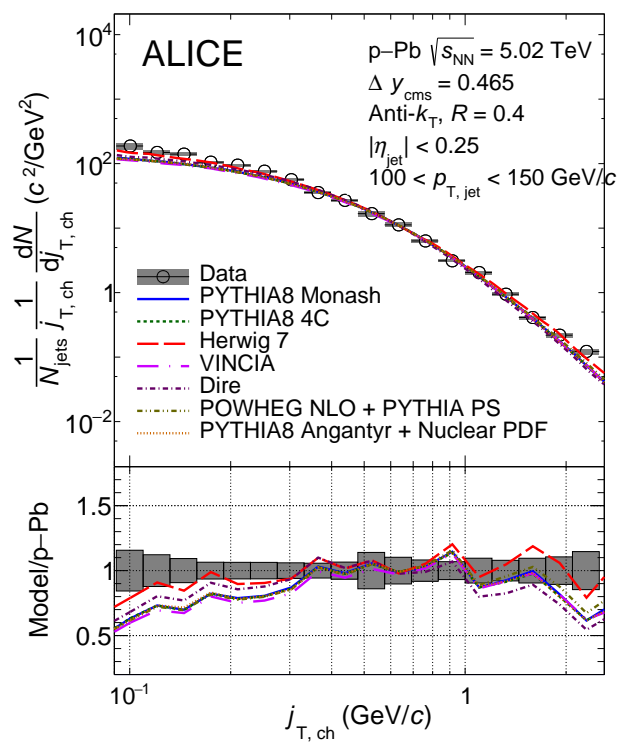


Figure B.3: The j_T distribution in $p\text{-Pb}$ collisions at $\sqrt{s_{\text{NN}}} = 5.02 \text{ TeV}$ for $100 < p_{T,\text{jet}} < 150 \text{ GeV}/c$ comparing to theoretical models in pp and $p\text{-Pb}$ collisions.